

Outline

- 1 Prelims
 - Physics
 - Math
- 2 Platonic QHE
- 3 Virtual work & Hall conductance
- 4 Chern=Kubo

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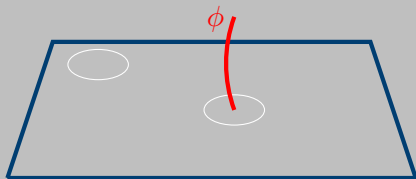
- 4 Chern=Kubo

Aharonov-Bohm flux tubes

Quantum flux

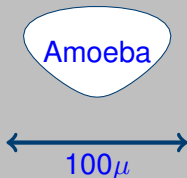


$$\oint A \cdot dx = \begin{cases} \phi & \text{winds origin} \\ 0 & \text{otherwise} \end{cases}$$



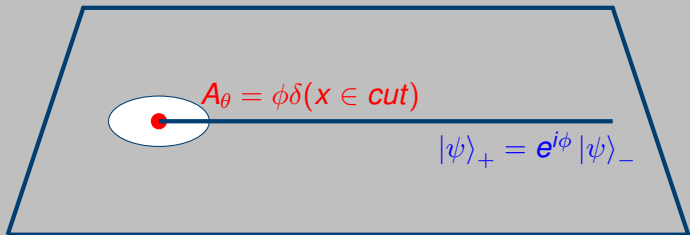
$$\Phi_0 = 2\pi \underbrace{\frac{\hbar}{e}}_{\text{fundamental}} = 2\pi$$

- Flux through a micro-organism



AB Periodicity

Flux tube modifies boundary condition:



b.c. 2π periodic in ϕ :

$$|\psi\rangle_+ = e^{i\phi} |\psi\rangle_-$$

AB periodicity

$$H(\phi + 2\pi) = UH(\phi)U^*$$

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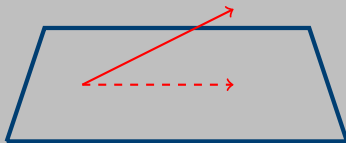
4 Chern=Kubo

Projections: P

- Orthogonal projections

$$\underbrace{P^2 = P}_{\text{projection}}, \quad \underbrace{P = P^*}_{\text{orthogonal}}$$

$$P = \sum_1^d |\psi_j\rangle \langle \psi_j|, \quad \langle \psi_i | \psi_j \rangle = \delta_{ij}$$



$$\underbrace{P_{\perp} = \mathbb{1} - P}_{\text{complementary}}$$

$$PP_{\perp} = P_{\perp}P = 0$$

Family of projections

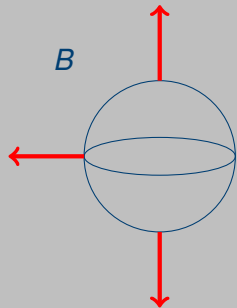
Paradigm

Berry: Spin in magnetic field B

$$P(B) = \frac{\mathbb{1} + H(\hat{B})}{2}, \quad \hat{B} = \frac{B}{|B|}$$

$$H(B) = B \cdot \sigma = \frac{1}{2} \begin{pmatrix} B_3 & B_1 - iB_2 \\ B_1 + iB_2 & -B_3 \end{pmatrix}$$

$P(B)$ sick at $B = 0 \Leftrightarrow H(0)$ degenerates



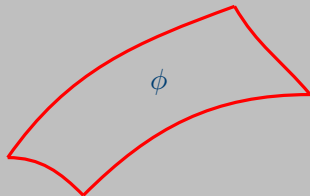
Family of projections

Parameter=control space

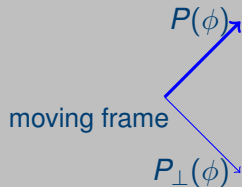
$\phi \in$ parameter space=control space

$P(\phi) : (\text{parameter space}) \mapsto \text{smooth projections}$

Parameter space



Hilbert space



dP

Motion of projections

$$\underbrace{P^2 = P}_{\text{matrices}}$$

$$P dP + dP P = dP$$

Corrolary

$$dP P = (\mathbb{1} - P)dP = P_{\perp} dP$$

$$P dP P = \underbrace{P P_{\perp}}_{=0} dP = 0$$

Kato

$$P dP P = 0$$

Kato evolution

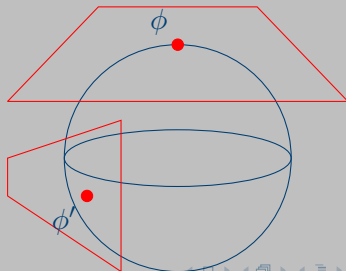
Unitary evolution in evolving subspaces

Kato's unitary evolution

$$\underbrace{P = U P_0 U^*}_{\text{Notion of parallel transport}}, \quad U = U(\phi), \quad P = P(\phi)$$

Who generates U ?

$$i dU = \underbrace{A}_{\text{generator}} U$$



Kato's evolution

Commutator equation

Generator satisfies commutator equation

$$dP = i[\mathcal{A}, P]$$

Proof:

$$P_0 = U^* P U \implies 0 = (dU^*) P U + U^* dP U + U^* P dU$$

$$0 = \underbrace{U(dU^*)}_{-(dU)U^*} P + dP + P \underbrace{(dU)U^*}_{-i\mathcal{A}}$$

\mathcal{A} : Not unique!

Ambiguity: commutant (P)

Kato's evolution

Generator

Commutator equation for \mathcal{A} :

$$dP = i[\mathcal{A}, P]$$

A Generator

$$\mathcal{A} = \underbrace{i(dU)U^*}_{\text{Definition}} = \underbrace{-i[dP, P]}_{\text{Generator}}$$

Verify:

$$i[\mathcal{A}, P] = [[dP, P], P] = (dP)P - \underbrace{2P(dP)P}_{=0} + P dP = dP$$

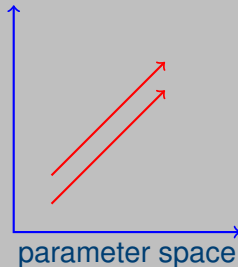
Parallel transport

Connection

Parallel transport: No motion in P

$$\underbrace{|\psi\rangle = P|\psi\rangle}_{\text{vector} \in P}, \quad \underbrace{0 = Pd|\psi\rangle}_{\text{no-motion}}$$

$$\begin{aligned} d|\psi\rangle &= d(P|\psi\rangle) = (dP)|\psi\rangle + \underbrace{Pd|\psi\rangle}_{=0} \\ &= (dP)P|\psi\rangle = \underbrace{[dP, P]}_{i\mathcal{A}}|\psi\rangle \end{aligned}$$



Covariant derivative:

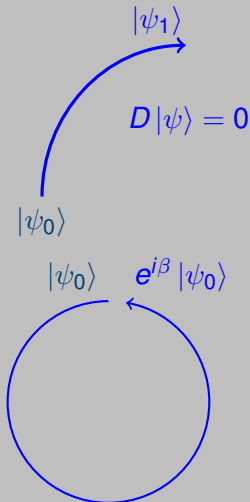
$$D = (d - i\mathcal{A}), \quad D|\psi\rangle = 0 \Leftrightarrow Pd|\psi\rangle = 0$$

Parallel transport

Berry's phase

- 1-D projection: $P = |\psi\rangle \langle\psi|$
- Parallel transport:
 $0 = P d|\psi\rangle = |\psi\rangle \langle\psi| d\psi$
- Parallel transport
 \implies **No local Berry's phase**

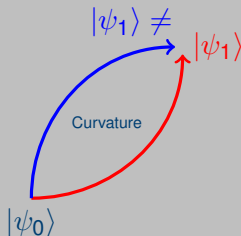
$$\begin{aligned}
 0 &= \langle\psi|d\psi\rangle - \frac{1}{2}d(\underbrace{\langle\psi|\psi\rangle}_{=1}) \\
 &= \frac{\langle\psi|d\psi\rangle - \langle d\psi|\psi\rangle}{2} \\
 &= i \underbrace{\text{Im} \langle\psi|d\psi\rangle}_{\text{Berry's phase}}
 \end{aligned}$$



Curvature

Failure of parallel transport

Parallel transport is path dependent:



Curvature=Failure of parallel transport

$$\Omega_{jk} = \underbrace{i[D_j, D_k]}_{\text{definition}} = \underbrace{(\partial_j \mathcal{A}_k - \partial_k \mathcal{A}_j) - i[\mathcal{A}_j, \mathcal{A}_k]}_{\text{Non-abelian magnetic fields}}$$

Curvature for projections

$$P(dP)(dP)P$$

$$\text{Curvature} = iP(dp)(dP)P$$

$$\Omega_{jk} = \underbrace{i[D_j, D_k]}_{\text{definition}} = i[\partial_j P, \partial_k P]$$

Proof:

$$\begin{aligned} [D_j, D_k] P &= [Pd_j, Pd_k]P \\ &= P(\partial_j P)(\partial_k P) - P(\partial_k P)(\partial_j P) \\ &= P[\partial_j P, \partial_k P] \end{aligned}$$

Curvature

1-D projection

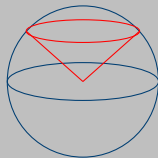
$$1\text{-D: } P = |\psi\rangle \langle\psi|$$

$$\Omega_{jk} P = i [\partial_j P, \partial_k P] P = \overbrace{i(\langle\partial_j\psi|\partial_k\psi\rangle - \langle\partial_k\psi|\partial_j\psi\rangle)}^{\text{Berry's curvature}} P$$

Example: Spin 1/2

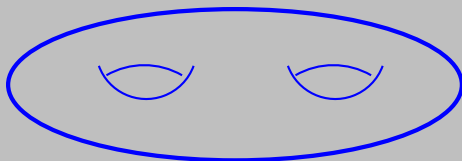
$$H = \Phi \cdot \sigma, \quad P = \frac{\mathbb{1} + \hat{H}}{2}$$

$$\Omega_{jk} P = \underbrace{\varepsilon_{jkl} \frac{\Phi_\ell d\Phi_j d\Phi_k}{4|\Phi|^3}}_{1/2 \text{ spherical angle}} P$$



Gauss Bonnet

Geometry meets topology



Gauss-Bonnet: Gaussian curvature & genus

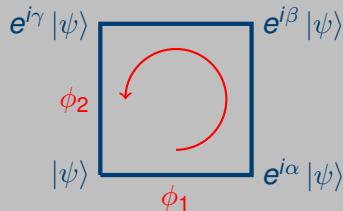
$$\frac{1}{2\pi} \int \underbrace{\Omega}_{\text{Curvature}} dS = 2(1 - \text{genus})$$

Chern numbers

Proof for torus (TKNN)

- $P(\phi_1, \phi_2)$ periodic
- $|\psi(\phi_1, \phi_2)\rangle$ periodic up to phase:

$$\begin{aligned} |\psi(0, 0)\rangle &= e^{-i\alpha} |\psi(2\pi, 0)\rangle \\ &= e^{-i\gamma} |\psi(0, 2\pi)\rangle \\ &= e^{-i\beta} |\psi(2\pi, 2\pi)\rangle \end{aligned}$$



- Angle counted mod 2π

$$(\alpha - 0)_{\text{mod } 2\pi} + (\beta - \alpha)_{\text{mod } 2\pi} + (\gamma - \beta)_{\text{mod } 2\pi} + (0 - \gamma)_{\text{mod } 2\pi} = 0$$

Chern numbers

$$i \int_T \langle d\psi | d\psi \rangle = i \oint_{\partial T} \langle \psi | d\psi \rangle \in 2\pi\mathbb{Z}$$

Chern numbers

Projections

Chern numbers

$$\text{Chern}(P, \mathcal{M}) = \frac{i}{2\pi} \int_{\mathcal{M}} \text{Tr} P[\partial_j P, \partial_k P] d\Phi_j d\Phi_k \in \mathbb{Z},$$

- \mathcal{M} : 2-D compact manifold (no bdry= $\partial\mathcal{M} = 0$)
- $\text{Chern}(P, \mathcal{M})$ invariant under smooth deformations of P
- P singular at eigenvalue crossing—dim P jumps

Chern numbers

Facts

- $\text{Chern}(0, \mathcal{M}) = \text{Chern}(\mathbb{1}, \mathcal{M}) = 0$
- $\text{Chern}(P_1 \oplus P_2, \mathcal{M}) = \text{Chern}(P_1, \mathcal{M}) + \text{Chern}(P_2, \mathcal{M})$

Chern numbers

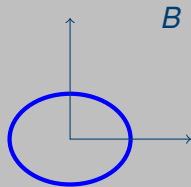
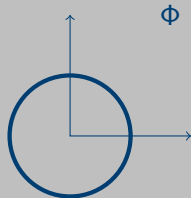
From sphere to ball

- $H = B \cdot \sigma$, $B = \underbrace{(B_x, B_y, B_z)}_{3-D \text{ space}}$
- Linear map of parameter space:
 $B = g \Phi$, $\Phi = (\Phi_x, \Phi_y, \Phi_z)$
- $\det g \neq 0$

Chern

$$H(\Phi) = \sum_{j,k=1}^3 g_{jk} \Phi^k \sigma^j, \quad P(\Phi) = \frac{\mathbb{1} + \hat{H}}{2}$$

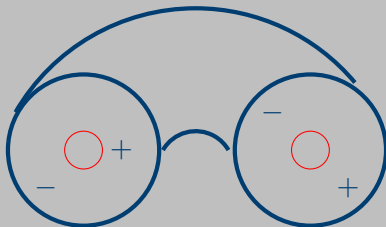
$$\text{Chern}(P) = \text{sgn } \det g$$



Chern numbers

What is counted?

Contracting into the solid torus

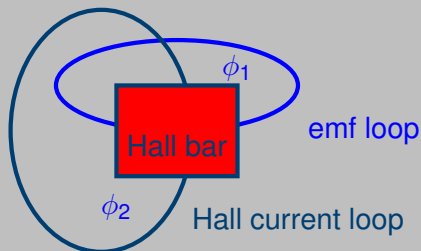


Simon

$$Chern(P, T) = \underbrace{\sum \text{sgn det } g(\Phi_d)}_{\text{degeneracies}}$$

QHE

Driving and response

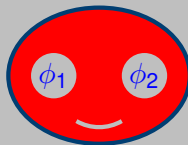
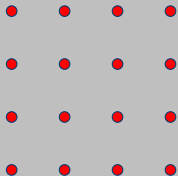


Platonic

- Driving: $\text{emf} = \dot{\phi}_1$
- Response: Hall current = $\frac{\partial H}{\partial \phi_2}$
- $H(\phi_1, \phi_2)$: Periodic, nondegenerate, hermitian matrix

Variations on a theme

Bloch momenta & controls



- Periodic
- (k_1, k_2) conserved
- Bloch momenta
- Brillouin Zone
- ∞ noninteracting (gapped) fermions
- Multiply connected
- (ϕ_1, ϕ_2) controls
- Fluxes
- Aharonov-Bohm period
- **Interacting** (finite)

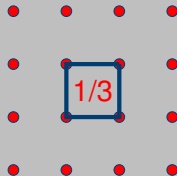
Example: 3×3 matrix function

Hofstadter Butterfly with flux $1/3$

$$H(\phi) = e^{i\phi_1} \underbrace{T}_{\text{translation}} + e^{i\phi_2} \underbrace{S}_{\text{shift}} + h.c.$$

$$T = \underbrace{\begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}}_{\text{lattice translation}}, \quad S = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & \omega & 0 \\ 0 & 0 & \bar{\omega} \end{pmatrix}}_{S=FTF^*}$$

$$\omega = e^{2\pi i/3}$$



Hofstadter Model $B = 1/3$

Virtual work

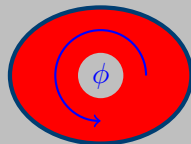
Q-observable

- $H : (\text{parameter space } \phi) \mapsto \text{Hamiltonian}$
- Virtual work

$$\delta H = \underbrace{\frac{dH(\phi)}{d\phi}}_{\text{observable}} \delta\phi$$

- Loop current:
Virtual work of Aharonov-Bohm flux

$$I = \frac{dH}{d\phi}$$



Charge transport

Time dependent Feynman-Hellman

Virtual work=Rate of Berry's phase

$$\underbrace{\langle \psi | \partial_\phi H | \psi \rangle}_{\text{Virtual work}} = \partial_t \underbrace{(i \langle \psi | \partial_\phi \psi \rangle)}_{\text{Berry's phase}}$$

- Schrödinger

$$i \partial_t |\psi\rangle = H(\phi) |\psi\rangle$$

- Pf:

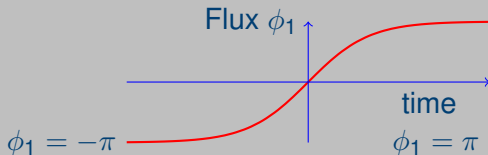
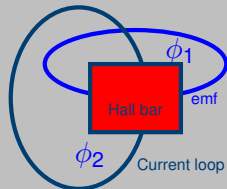
$$\begin{aligned} \langle \psi | \partial_\phi H | \psi \rangle &= \partial_\phi \langle \psi | H | \psi \rangle - \overbrace{\langle \partial_\phi \psi | H | \psi \rangle - \langle \psi | H | \partial_\phi \psi \rangle}^{\text{time-independent}=0} \\ &= \partial_\phi \langle \psi | H | \psi \rangle - \underbrace{\langle \partial_\phi \psi | H | \psi \rangle}_{i \langle \partial_t \psi |} - \underbrace{\langle \psi | H | \partial_\phi \psi \rangle}_{-i \langle \partial_t \psi |} \\ &= \partial_t (i \langle \psi | \partial_\phi \psi \rangle) \end{aligned}$$

Hall conductance

Definition

Def: Flux averaged Hall conductance:

$$2\pi\sigma = \underbrace{\frac{1}{2\pi} \int_{\pi}^{\pi} d\phi_2 \int_{-\infty}^{\infty} dt \langle \psi_t | \partial_2 H | \psi_t \rangle dt}_{\text{flux average charge transport}}$$



Control averaging \iff Filled band.
Averaging: Gets rid of persistent currents

Chern=Kubo

Geometry of transport

In adiabatic limit: flux average transport=Chern

$$2\pi\sigma \xrightarrow{\text{adiabatic}} \underbrace{\frac{i}{2\pi} \int_{\pi}^{\pi} d\phi_1 d\phi_2 \text{Tr} P[\partial_1 P, \partial_2 P]}_{\text{Chern}}$$

- Loop currents: $\langle \psi | \partial_{\phi} H | \psi \rangle = \partial_t (i \langle \psi | \partial_{\phi} \psi \rangle)$
- Adiabatic limit: $|\psi\rangle \mapsto |\psi_A\rangle$, $H \mapsto \mathcal{A}$
- Adiabatic loop currents: $\partial_t (i \langle \psi_A | \partial_{\phi} \psi_A \rangle) = \langle \psi_A | \partial_{\phi} \mathcal{A} | \psi_A \rangle$

$$i \langle \psi_A | \partial_{\phi} [\dot{P}, P] | \psi_A \rangle = i \text{Tr} P \partial_{\phi} [\partial_t P, P] = i \text{Tr} P [\partial_{\phi} P, \partial_t P]$$

Critique

Why Platonic?

- $\phi = \text{control}$: Too general
- $\phi = \text{Bloch momenta}$: Too special
- Gap condition: Too strong—localization
- Where is 2-D?
- Where is thermodynamic limit
- Why average?
- Sample appear to be less important than connecting circuit
- What about fractions?