Homotopy Theory of Topological Insulators

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SFB | TR12



• Single particle Hilbert space $l^2(\mathbb{Z}^d)\otimes \mathbb{C}^n$

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· Insulator ground state given by assignment

 $\mathbf{k} \mapsto V(\mathbf{k}) =$ subspace spanned by *m* occupied states



Classifying map

 $T^d o \operatorname{Gr}_m(\mathbb{C}^n)$ $\mathbf{k} \mapsto V(\mathbf{k})$



Classifying map $T^d \rightarrow \operatorname{Gr}_m(\mathbb{C}^n)$ $\mathbf{k} \mapsto V(\mathbf{k})$ Vector subbundle of $T^d \times \mathbb{C}^n$

$$p: \mathcal{V} o T^d$$
 $p^{-1}(\mathbf{k}) = V(\mathbf{k})$



Classifying mapVector subbundle of $T^d \times \mathbb{C}^n$ $T^d \to \operatorname{Gr}_m(\mathbb{C}^n)$ $p: \mathcal{V} \to T^d$ $\mathbf{k} \mapsto V(\mathbf{k})$ $p^{-1}(\mathbf{k}) = V(\mathbf{k})$

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- Example: d = 1, n = 2 and m = 1





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Other equivalence relations:

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- Isomorphic as vector bundles: $\operatorname{Vect}_m^{\mathbb{C}}(T^d) = [T^d, \operatorname{Gr}_m(\mathbb{C}^{\infty})]$
- Stably equivalent: $\tilde{K}(T^d) = [T^d, \operatorname{Gr}_{\infty}(\mathbb{C}^{\infty})]$



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- 1. More topological phases
- 2. More restrictive definition of "strong" topological insulators



Symmetry $T \circ I$ in d = 1 with m = 1 of n = 2 bands occupied:

 $\mathcal{S}^1 o \operatorname{Gr}_1(\mathbb{R}^2)$ $(x, y) \mapsto \text{line through } (x, y)$



2 in
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class	1	2	3
А	0	Z	0
AIII	\mathbb{Z}	0	\mathbb{Z}
D	\mathbb{Z}_2	\mathbb{Z}	0
DIII	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}
All	0	\mathbb{Z}_2	\mathbb{Z}_2
CII	\mathbb{Z}	0	\mathbb{Z}_2
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[Moore,Ran,Wen 2008]



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Usual definition^[Kitaev 2009; RK,Guggenheim 2014] for large *n*, *m*:

$$[T^d, \operatorname{Gr}_m(\mathbb{C}^n)] = \prod_{p=1}^d [S^p, \operatorname{Gr}_m(\mathbb{C}^n)]^{\binom{d}{p}}$$



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Example $d = 3, m, n \gg 1$: $[S^1, \operatorname{Gr}_m(\mathbb{C}^n)] = 0$ $[S^2, \operatorname{Gr}_m(\mathbb{C}^n)] = \mathbb{Z}$ (Chern number) $[S^3, \operatorname{Gr}_m(\mathbb{C}^n)] = 0$



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Interpretation: Chern insulators stacked orthogonal to vector $(n_1, n_2, n_3) \in \mathbb{Z} \times \mathbb{Z} \times \mathbb{Z}$



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Generalises to all d and \mathbb{Z}_2 -equivariant classes^[RK,Guggenheim 2014]



Example with symmetry $T \circ I$ for d = 2, m = 1, n = 3:

$$[S^1, \operatorname{Gr}_1(\mathbb{R}^3)] = \mathbb{Z}_2$$
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Breakdown of strong invariant $\mathbb{N} \to \mathbb{Z}_2$





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Stacked realisation



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- Stability typically reached quickly, but some exceptions exist (Hopf insulator/superconductor)
- Only non-trivial maps from S^d are strong in general



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