

Projected entangled-pair states for chiral topological phases

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Projected entangled-pair state (PEPS)

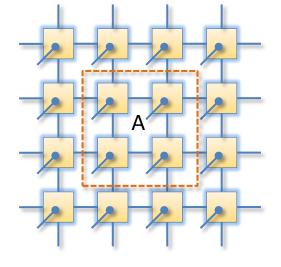
$$\begin{split} |P\rangle_{i} &= \sum_{s_{i}=1}^{d} \sum_{\alpha_{i},\beta_{i},\gamma_{i},\delta_{i}=1}^{D} A_{\alpha_{i}\beta_{i}\gamma_{i}\delta_{i}}^{s_{i}} |s_{i}\rangle |\alpha_{i},\beta_{i},\gamma_{i},\delta_{i}\rangle \\ &= \sum_{\alpha=1}^{D} |\alpha_{i},\alpha_{j}\rangle \\ &= \sum_{\alpha=1}^{D} |\alpha_{i},\alpha_{j}\rangle \\ D: \text{ bond dimension} \\ |\psi\rangle &= \left(\bigotimes_{\langle ij\rangle} \langle I|_{i,j}\right) \left(\bigotimes_{i=1}^{N} |P\rangle_{i}\right) \\ &= \sum_{\{s\}} \left(\sum_{\{\alpha\beta\gamma\delta\}} \cdots A_{\alpha_{i}\beta_{i}\gamma_{i}\delta_{i}}^{s_{i}} \cdots\right) |\ldots s_{i}\ldots\rangle \\ Vertrate \& \text{ Given cond-mat/0407066} \end{split}$$

Verstraete & Cirac, cond-mat/0407066

Parent Hamiltonian

Existence of null space if $d^{N_A} > D^{N_{\partial A}}$

$$H = \sum_{i} h_{i} \qquad \begin{array}{c} h_{i} \geq 0 \\ h_{i} |\psi\rangle = 0 \end{array}$$



 $\operatorname{Rank}(\rho_A) \le D^{N_{\partial A}}$

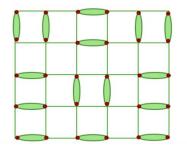
$$S(A) \sim N_{\partial A}$$

Conjecture: area law holds for all gapped ground states of local Hamiltonians

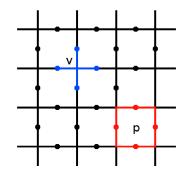
Topological PEPS

Resonating valence bonds

Anderson, Mater. Res. Bull. (1973)

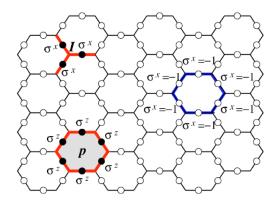


Toric code Kitaev, Ann. Phys. (2003)



Levin-Wen string nets

Levin & Wen, PRB (2005)



Q: What about chiral topological states?

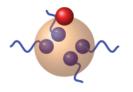
Verstraete, Wolf, Perez-Garcia & Cirac, PRL (2006); Buerschaper & Aguado, PRB (2009); Gu, Levin, Swingle & Wen, PRB (2009)

PEPS for free fermionic chiral topological states

T.B. Wahl, HHT, N. Schuch & J.I. Cirac, PRL (2013); T.B. Wahl, S.T. Hassler, HHT, N. Schuch & J.I. Cirac, arXiv:1405.0447. See also J. Dubail & N. Read, arXiv: 1307.7726.

Fermionic Gaussian state

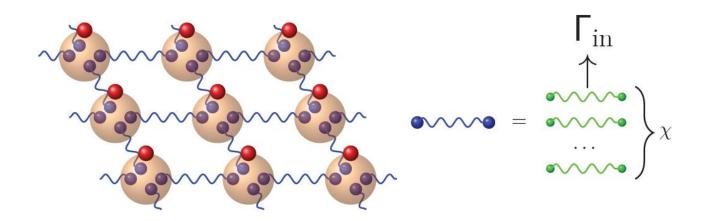
- PEPS projector $|P\rangle$ is a Gaussian state (ground/thermal states of quadratic Hamiltonians)



• Gaussian states are characterized by the convariance matrix

$$\begin{split} \Gamma_{jl} &= \frac{i}{2} \mathrm{tr}(\rho[c_j,c_l]) & \Gamma \leftrightarrow \rho \\ \end{split}$$
Pure state: $\Gamma\Gamma^{\top} = \mathbb{I}$ Majorana
Mixed state: $\Gamma\Gamma^{\top} < \mathbb{I}$

Gaussian Fermionic PEPS (GFPEPS)



• GFPEPS projector characterized by a convariance matrix:

$$M = \begin{pmatrix} A & B \\ -B^{\top} & D \end{pmatrix}$$
 Pure to pure: $MM^{\top} = \mathbb{I}$
Pure to mix: $MM^{\top} < \mathbb{I}$

• Covariance matrix for GFPEPS

$$\Gamma_{\rm out}(\mathbf{k}) = B[D - \Gamma_{\rm in}(\mathbf{k})]^{-1}B^{\top} + A$$

Kraus, Schuch, Verstraete & Cirac, PRA (2010)

Example of a chiral PEPS: Chern insulator

$$A_{\lambda} = (-1+2\lambda) \begin{pmatrix} \omega & 0\\ 0 & -\omega \end{pmatrix}$$

$$B_{\lambda} = \sqrt{\frac{\lambda-\lambda^2}{2}} \begin{pmatrix} I-\omega & I+\omega & -\sqrt{2}\,\omega & \sqrt{2}\,I\\ I-\omega & -I-\omega & \sqrt{2}\,I & -\sqrt{2}\,\omega \end{pmatrix}$$

$$D_{\lambda} = \begin{pmatrix} 0 & (-1+\lambda)\,I & -\frac{\lambda}{\sqrt{2}}\,I & \frac{\lambda}{\sqrt{2}}\,I\\ (1-\lambda)\,I & 0 & -\frac{\lambda}{\sqrt{2}}\,I & -\frac{\lambda}{\sqrt{2}}\,I\\ \frac{\lambda}{\sqrt{2}}\,I & \frac{\lambda}{\sqrt{2}}\,I & 0 & (-1+\lambda)\,I\\ -\frac{\lambda}{\sqrt{2}}\,I & \frac{\lambda}{\sqrt{2}}\,I & (1-\lambda)\,I & 0 \end{pmatrix}$$

$$I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$
$$\omega = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$
$$0 < \lambda < 1$$

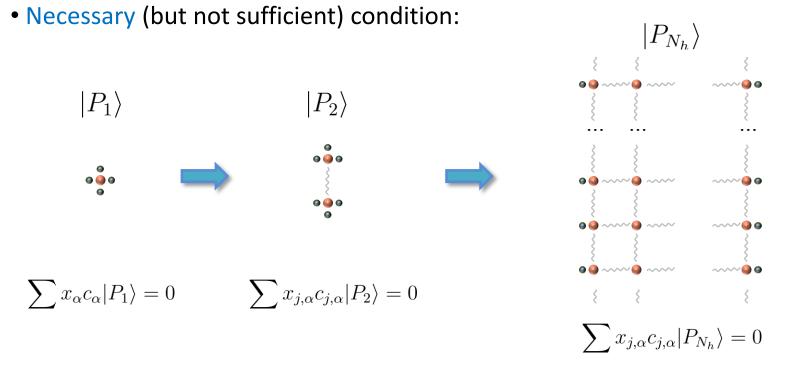
• Gapless Hamiltonian with short-range hoppings

$$H_1 = -i \sum_{\mathbf{k},s,t} \varepsilon(\mathbf{k}) [\Gamma_{\text{out}}(\mathbf{k})]_{st} d_{\mathbf{k},s} d_{-\mathbf{k},t}$$

• Gapped Hamiltonian with powerlaw decaying hoppings $(1/r^3)$, C = -1

$$H_2 = -i \sum_{\mathbf{k},s,t} [\Gamma_{\text{out}}(\mathbf{k})]_{st} d_{\mathbf{k},s} d_{-\mathbf{k},t}$$

Chirality of GFPEPS



• The existence of $\sum x_{\alpha}c_{\alpha}$ is related to the existence of chiral edge modes

 The GFPEPS is non-injective (otherwise adiabatically connected to a trivial state)

Approximating a Chern insulator

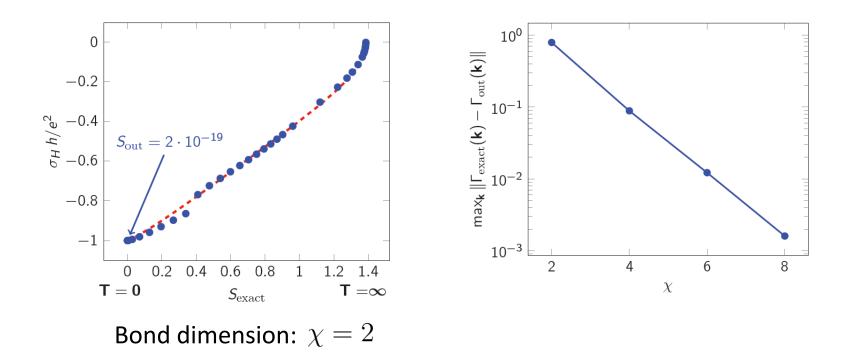
Do PEPS provide a good approximation to the ground/thermal state of a Chern insulator?

$$H = \sum_{\mathbf{k}} (a_{\mathbf{k},\uparrow}^{\dagger}, a_{\mathbf{k},\downarrow}^{\dagger}) \left(\boldsymbol{\sigma} \cdot \mathbf{d}(\mathbf{k}) \right) (a_{\mathbf{k},\uparrow}, a_{\mathbf{k},\downarrow})^{\top}$$
$$\mathbf{d}(\mathbf{k}) = (\sin k_y, -\sin k_x, 2 - \cos k_x - \cos k_y - e_S)$$
$$e_S = 1 \quad C = -1$$

X.-L. Qi, Y.-S. Wu & S.-C. Zhang, PRB (2006)

Approximating a Chern insulator

Do PEPS provide a good approximation to the ground/thermal state of a Chern insulator?

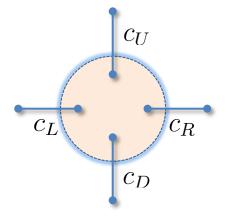


PEPS for interacting chiral topological states

In preparation...

Chiral PEPS example from projective construction

$$|\text{GPEPS}\rangle = \left(\bigotimes_{\langle ij \rangle} \langle \omega |_{i,j}\right) \left(\bigotimes_{i=1}^{N} |P\rangle_{i}\right)$$

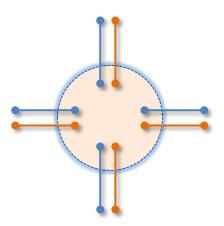


topological superconductor with C = 1 (class D)

Projective construction:

$$|\Psi\rangle = P_{\rm G} |\rm{GPEPS}\rangle_1 |\rm{GPEPS}\rangle_2$$

Gutzwiller projector -- only single occupancy allowed!



Projective construction of $SO(n)_1$ state

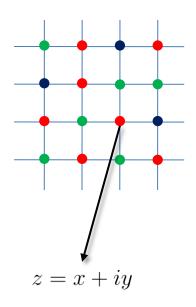
$$|\Psi\rangle = P_{\rm G} \prod_{\alpha=1}^{n} \exp\left(\sum_{i < j} \frac{1}{z_i - z_j} a_{i,\alpha}^{\dagger} a_{j,\alpha}^{\dagger}\right) |0\rangle$$

n even

$$\Psi(\alpha_1,\ldots,\alpha_N) = \langle \chi^{\alpha_1}(z_1)\chi^{\alpha_2}(z_2)\cdots\chi^{\alpha_N}(z_N)\rangle$$

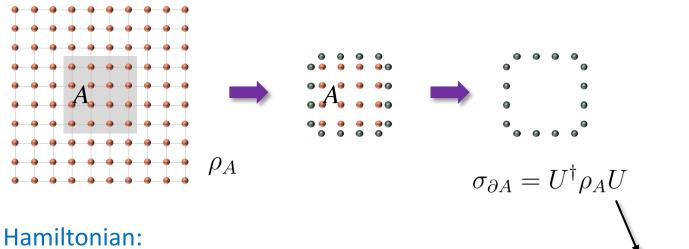
Bulk-edge correspondence (Moore-Read):

- Edge CFT: SO(n)₁ with central charge c = n/2٠
- I, s, \bar{s}, v Anyonic quasiparticles ٠ I, s, vn odd $h_I = 0$ $h_s = h_{\bar{s}} = \frac{n}{16}$ $h_v = \frac{1}{2}$



HHT, Phys. Rev. B 87, 041103 (2013)

Boundary theory of PEPS



Boundary Hamiltonian:

$$\sigma_{\partial A} = e^{-H_{\partial A}}$$

isometry

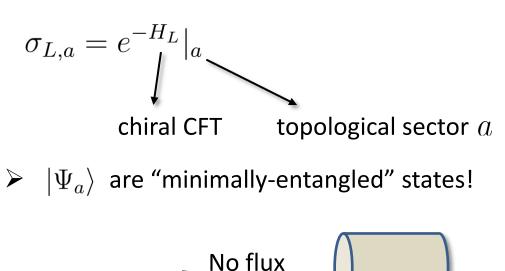
> ... gives entanglement spectrum

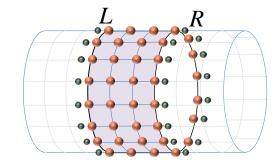
> ... can be easily determined (exactly or approximiately)

Cirac, Poilblanc, Schuch & Verstraete, PRB (2011)

Boundary theory of chiral PEPS

Entanglement spectrum for chiral states $|\Psi_a\rangle$:





Each contains two sectors!

Li & Haldane, PRL (2008); Qi, Katsura & Ludwig, PRL (2012); Zhang, Grover, Turner, Oshikawa & Vishwanath, PRB (2012)

With flux

PEPS

Outlook

• Chiral PEPS with exponentially decaying correlations and gapped shortrange parent Hamiltonian?

> Approach different from projective construction and discretization of conformal blocks?

• Gauge symmetry of PEPS local tensor as a unified description of both chiral and non-chiral topological states?

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Thank you for your attention!