

## Twelfth exercise sheet on Relativity and Cosmology I

Winter term 2018/19

**Release:** Mon, Jan. 21<sup>st</sup>

**Submit:** Mon, Jan. 28<sup>th</sup> in lecture

**Discuss:** Thu, Jan. 31<sup>st</sup>

In the following exercises, consider the coordinates  $(t, r, \theta, \phi)$ . The 2-dimensional equatorial spatial slice  $\Sigma$  is defined by  $t = \text{const.}$ ,  $\theta = \pi/2$ , and its induced metric reads  $d\sigma_\Sigma^2$ . Moreover, the Schwarzschild metric in such coordinates reads

$$ds^2 = -f(r) dt^2 + \frac{dr^2}{f(r)} + r^2 d\Omega^2, \quad f(r) := 1 - \frac{R_S}{r}, \quad R_S := 2GM, \quad d\Omega^2 := d\theta^2 + \sin^2 \theta d\phi^2.$$

Furthermore, denote the Euclidean metric on  $\mathbb{R}^n$  by  $d\bar{x}_{\mathbb{R}^n}^2$ .

**Exercise 38** (10 points): *Schwarzschild metric in isotropic coordinates*

Consider a coordinate transformation

$$r = \left(1 + \frac{R_S}{4\bar{r}}\right)^2 \bar{r}, \quad t = \bar{t}, \quad \theta = \bar{\theta}, \quad \phi = \bar{\phi}.$$

**38.1** Express the Schwarzschild metric in the new coordinates  $(\bar{t}, \bar{r}, \bar{\theta}, \bar{\phi})$ , which are called *isotropic*.

**38.2** Compare briefly the metric components in old and new coordinates at the horizon.

**38.3** In the *standard* coordinates, consider a radial range  $R_S < r < R$ .

Use the *isotropic* coordinates to calculate

- the surface area on the equatorial spatial slice between these radii, and
- the volume of a spherical shell with  $t = \text{const.}$  within the range.

**38.4** Compare your results in **38.3** to the corresponding quantities in the Euclidean space.

**Exercise 39** (4 points): *Isometric embedding I: the Schwarzschild space*

**39.1** Consider the cylindrical coordinates  $(r, \phi, z)$  of  $\mathbb{R}^3$ . Set

$$d\bar{x}_{\mathbb{R}^3}^2 \Big|_{z=z(r)} \equiv d\sigma_\Sigma^2$$

and integrate the resulting equation (*Flamm's paraboloid*).

**39.2** In the current case, if  $d\sigma_\Sigma^2 \rightarrow d\bar{x}_{\mathbb{R}^2}^2$  as  $r \rightarrow +\infty$ ,  $\Sigma$  is called *asymptotically flat*.

Analytically extend the embedding in **39.1**, and show that there can be *two* distinct regions, which are both asymptotically flat (*Einstein and Rosen*).

*Hint:* Sketching the embedding may help.

See overleaf.

**Exercise 40** (6 points): *Isometric embedding II: a wormhole*

Consider a spacetime  $\mathcal{M}_W$  with the metric

$$ds^2 = -dt^2 + dr^2 + (b^2 + r^2)d\Omega^2$$

where  $b$  is a constant of dimension length.

**40.1** Argue briefly that the equatorial spatial slice  $\Sigma$  of  $\mathcal{M}_W$  is representative for the latter.

**40.2** Consider the cylindrical coordinates  $(\bar{r}, \phi, z)$  of  $\mathbb{R}^3$ . Set

$$d\bar{x}_{\mathbb{R}^3}^2 \Big|_{\bar{r}=\bar{r}(r), z=z(r)} \equiv d\sigma_{\Sigma}^2$$

and integrate the resulting equation for  $z$  and  $\bar{r}$ .

**40.3** Argue briefly that there is a hole-like structure in  $\mathcal{M}_W$ .

*Hint:* Sketching the embedding may help.

*Remark.* The Einstein tensor of the metric reads

$$G_{\mu\nu} dx^\mu dx^\nu = \frac{b^2}{(b^2 + r^2)^2} (-dt^2 - dr^2 + (b^2 + r^2)d\Omega^2).$$

If the matter were modelled by an ideal fluid, its energy density would turn out to be negative,

$$\rho = T_{\mu\nu} u^\mu u^\nu = -\frac{1}{\kappa} \frac{b^2}{(b^2 + r^2)^2} < 0.$$

In other words, matter with negative energy density is needed to source this spacetime.

Inversely, one may also ask, if it is possible to have a viable wormhole spacetime sourced by matter with positive energy density. Under precise and additional assumptions, this has been excluded by the so-called *topological censorship theorem*.