

## Exercise 2: Metropolis Algorithm for the Ising Model

In this second exercise we will study the thermal phase transition that occurs in the two-dimensional Ising model, the first example of a dynamic system. The Ising model is defined by the Hamiltonian

$$\mathcal{H} = -J \sum_{\langle i,j \rangle} \sigma_i \sigma_j, \quad (1)$$

where the spins  $\sigma_i = \pm 1$  correspond to “up” and “down” spins and the sum runs over all pairs of nearest neighbor spins. We will consider the ferromagnetic case with  $J > 0$ .

For the underlying lattice we will consider a square lattice of linear extent  $L$  and  $N = L \times L$  sites. In order to minimize finite-size effects we will use *periodic* boundary conditions, which results in a lattice with a total of  $2N$  bonds.

To identify and characterize the thermal phase transition we will investigate thermal averages for a number of observables, which are generally defined as

$$A(T) = \frac{1}{Z} \sum_i A_i \exp(-\beta E_i),$$

where  $Z$  is the partition function of the system, the sum runs over all possible spin configurations  $i$ ,  $A_i$  and  $E_i$  are the values which the observable  $A$  and energy  $E$  have for a given configuration  $i$ , and  $\beta$  is the inverse temperature  $\beta = 1/(k_B T)$ .

We will calculate these thermal averages via Monte Carlo sampling for a range of temperatures  $T = 0.1, 0.2, \dots, 4$  (where we fix units by setting  $J = 1$ ) and system sizes  $L = 8, 16, 32$ .

1. Implement a **single spin-flip Metropolis algorithm** for this 2D Ising model. (You can build on your code for the percolation problem.)
2. Plot Monte Carlo averages of the **magnetization**  $M = \sum_i \sigma_i$  for the full temperature range. Perform measurements only after an initial set of – say – 10,000 thermalization sweeps where one sweep corresponds to  $N$  attempted spin flips. After this thermalization phase, perform one measurement for every sweep.
3. Plot Monte Carlo averages for the **Binder cumulant** of the magnetization  $U = 1 - \frac{\langle M^4 \rangle}{3 \langle M^2 \rangle^2}$  for the full temperature range.
4. Plot Monte Carlo averages for the **energy**  $E$  for the full temperature range.

5. Plot Monte Carlo averages for the **specific heat**  $C_v$ , which you can estimate either by (numerically) calculating the derivative  $dE/dT$  of the energy curve above, or more directly via  $\langle C_v \rangle = \beta^2/N(\langle E^2 \rangle - \langle E \rangle^2)$  by measuring Monte Carlo estimates for  $E^2$ .
6. From the two results plotted above can you reproduce the estimate of the thermal phase transition  $T_c = 2/\ln(1 + \sqrt{2}) \approx 2.269186$  for the square lattice?
7. Implement a **cluster update**, such as the Wolf algorithm explained in the lecture.
8. *Optional exercise:* For a fixed system size  $L = 32$  perform  $2^{16} = 65536$  measurements of the energy at the thermal transition temperature  $T = 2.269186$  for the single spin-flip algorithm and the cluster update algorithm.  
For both sequences perform a **binning analysis** of the sampled energies – for which algorithm does the error converge?