# Majorana Fermions in 1D Nanowires

Majorana Fermions in solid state systems feature nonabelian statistics, i.e braiding not only changes the phase but leads to different quantum state. Therefor an application as qbits seems possible.

## Kitaev chain

A toy tight binding Hamiltonian for a p-wave superconductor reads

$$H = \sum_{j=1}^{L} \left[ -\mu a_j^{\dagger} a_j - t \left( a_{j+1}^{\dagger} a_j + \text{h.c} \right) + \left( \Delta a_{j+1}^{\dagger} a_j^{\dagger} + \text{h.c} \right) \right]$$

Introduce Majorana operators  $\gamma_{2j-1} = a_j + ia_j^{\dagger}$  and  $\gamma_{2j} = -ia_j + ia_j^{\dagger}$  with  $\{\gamma_l, \gamma_m\} = 2\delta_{lm}$  and  $\gamma_j = \gamma_j^{\dagger}$  Easy to solve in cases

- 1.  $\Delta=t=0,\,\mu<0$ 
  - $\rightarrow H = -\frac{i}{2} \sum \gamma_{2j-1} \gamma_{2j}$
  - $\rightarrow$  Majoranas on same site pair
  - $\rightarrow$  Unique unoccupied groundstate

2. 
$$\Delta = t \neq 0, \ \mu = 0$$

 $\rightarrow H = it \sum_{j=1}^{L-1} \gamma_{2j} \gamma_{2j+1}$ 

- $\rightarrow$  Majoranas on neighbouring sites pair
- $\rightarrow \gamma_1$  and  $\gamma_{2L}$  are unpaired. Combine to highly nonlocal fermion  $d_M = 1/2 (\gamma_1 + i\gamma_{2L})$
- $\rightarrow$  groundstate two-fold degenerate:  $|0\rangle$  and  $|1\rangle=d_{M}^{\dagger}|0\rangle$

Associated with the transition is a topological invariant  $\nu = s_0 s_{\pi}$  where  $s_{0/\pi}$  is the sign of the kinetic energy at momentum  $0/\pi$ , i.e  $\nu$  is -1 for an odd number of band crossings in half the Brioullin zone and can only changes when the gap closes. For parameters in the range  $|\mu| < 2t$ , the topological phase is realized with MFs no longer sharply localized to the ends, but decaying exponentially and overlapping, giving rise to an interaction

$$H_{int} = \frac{i}{2} t \gamma' \gamma''$$

with  $t \propto \exp(-L/\xi)$  Therfor states split, but only exponentially  $\rightarrow$  neglectable. Realisation is difficult, because no pwave-SC is available and fermions must appear spin-less.

### **Physical Realization**

Solution: Semiconductor/s-wave-SC heterostructure can mimic p-wave-SC through Spin-Orbit coupling.

$$H = \int \mathrm{dx} \,\psi_{\alpha}^{\dagger}(x) \left\{ -\frac{\partial_x^2}{2m} - \mu - i\alpha \partial_x \sigma_y + V \sigma_x \right\} \psi_{\beta}(x) + H_{SC}$$
with

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$$H_{SC} = \int \mathrm{d}\mathbf{x} \left\{ \Delta \psi_{\uparrow} \psi_{\downarrow} + \mathrm{h.c} \right\}$$

Spin-Orbit- and Zeemann-coupling split bands and give rise to an effective spinless regime for  $|\mu| < V$ . The Hamiltonian then maps to Kitaev's model and the topological phase occurs at  $V > \sqrt{\Delta^2 + \mu^2}$ .



### Experiments

Two main effects should be visible. Measuring the DoS via tunnel spectroscopy should reveal a zero-energy-peak. A junction between two superconductors via a topological SC exhibits a fractional Josephson current.

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