

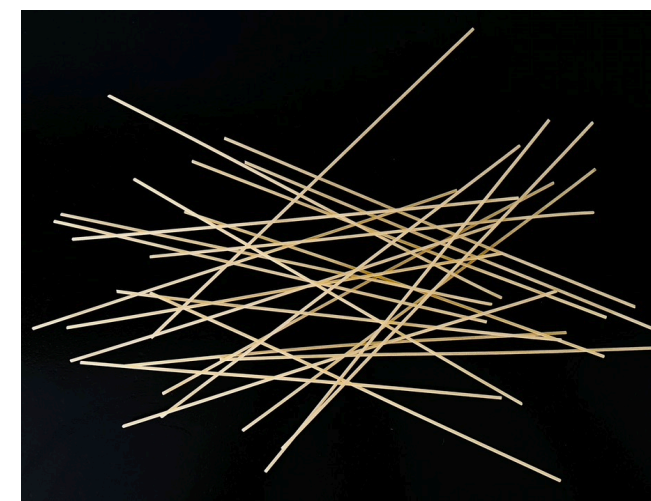
Universal principles of moiré band structures

Simon Trebst
University of Cologne

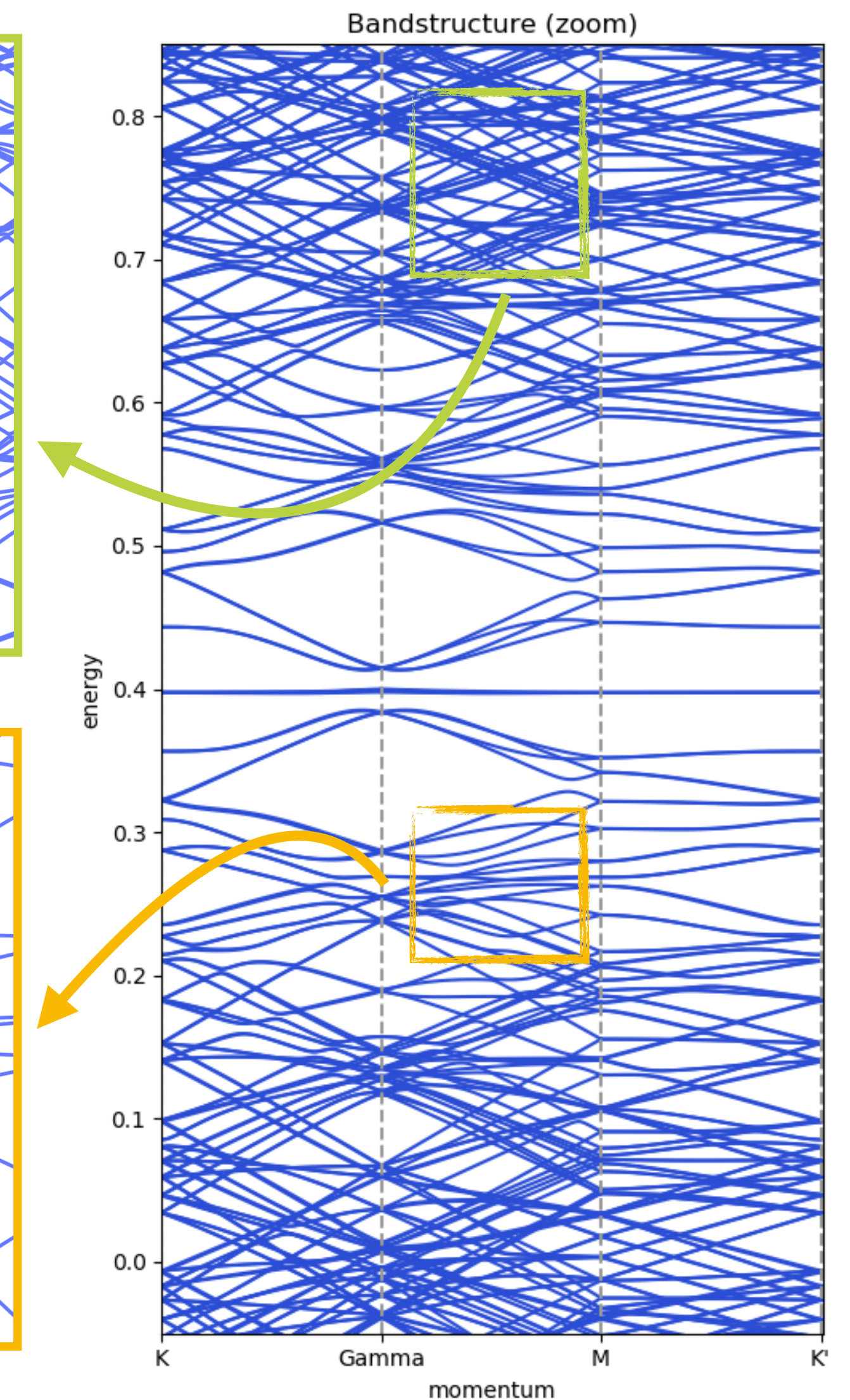
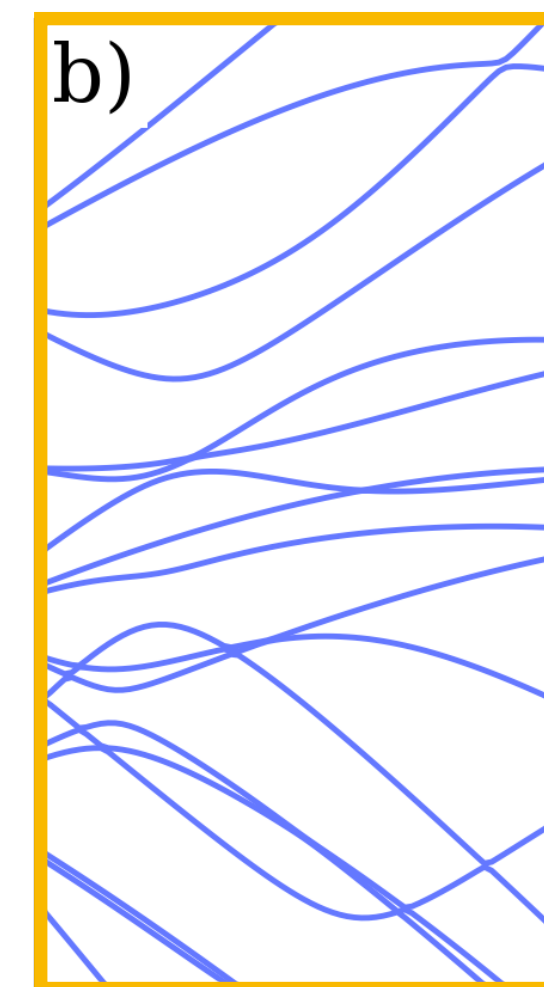
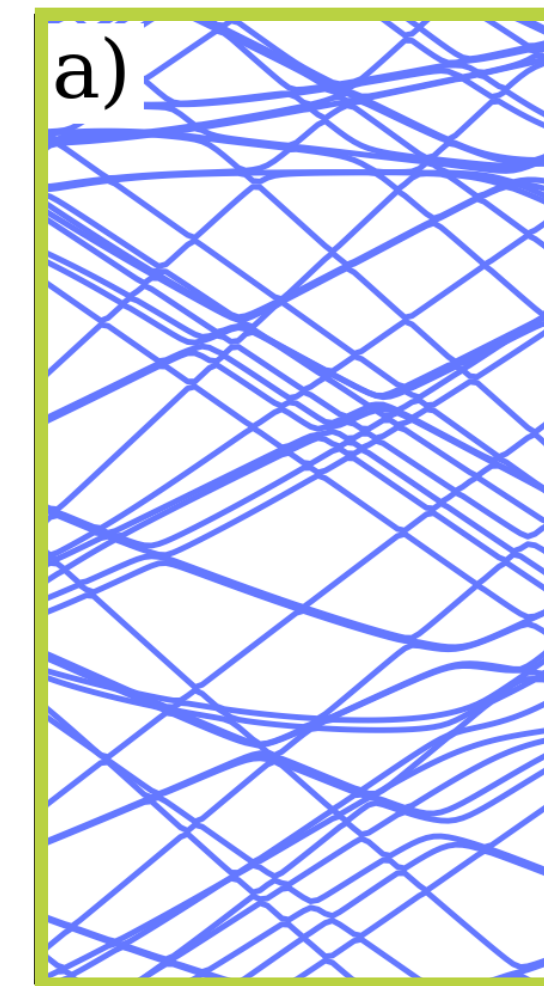
why most moiré bands are **not** flat

This talk

- the physics of moiré systems with **giant unit cells**
- **statistical** analysis
- **quantum chaos** versus Anderson **localization**
- **twisted bilayer graphene** and beyond



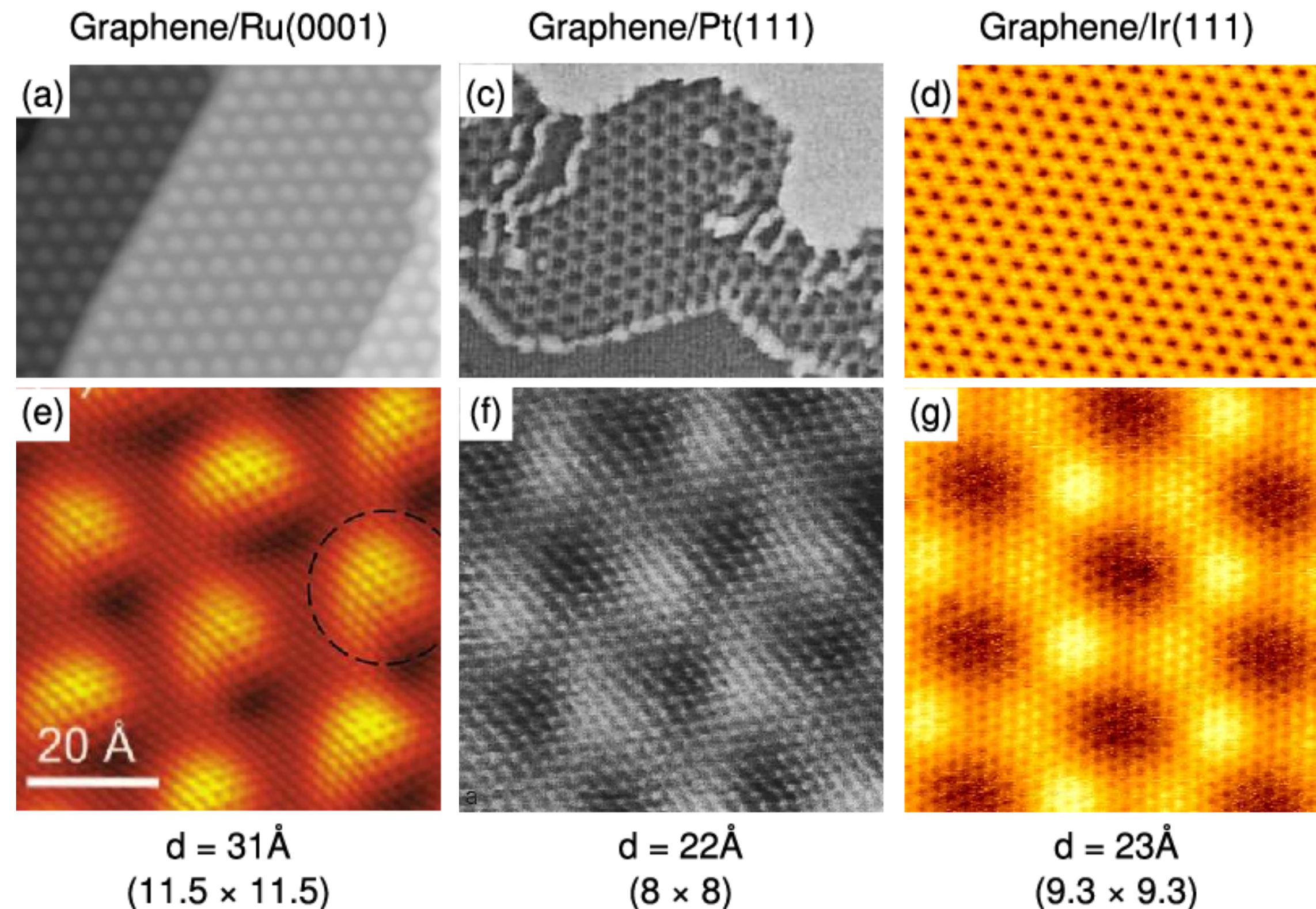
- What happens with the **spaghetti**?



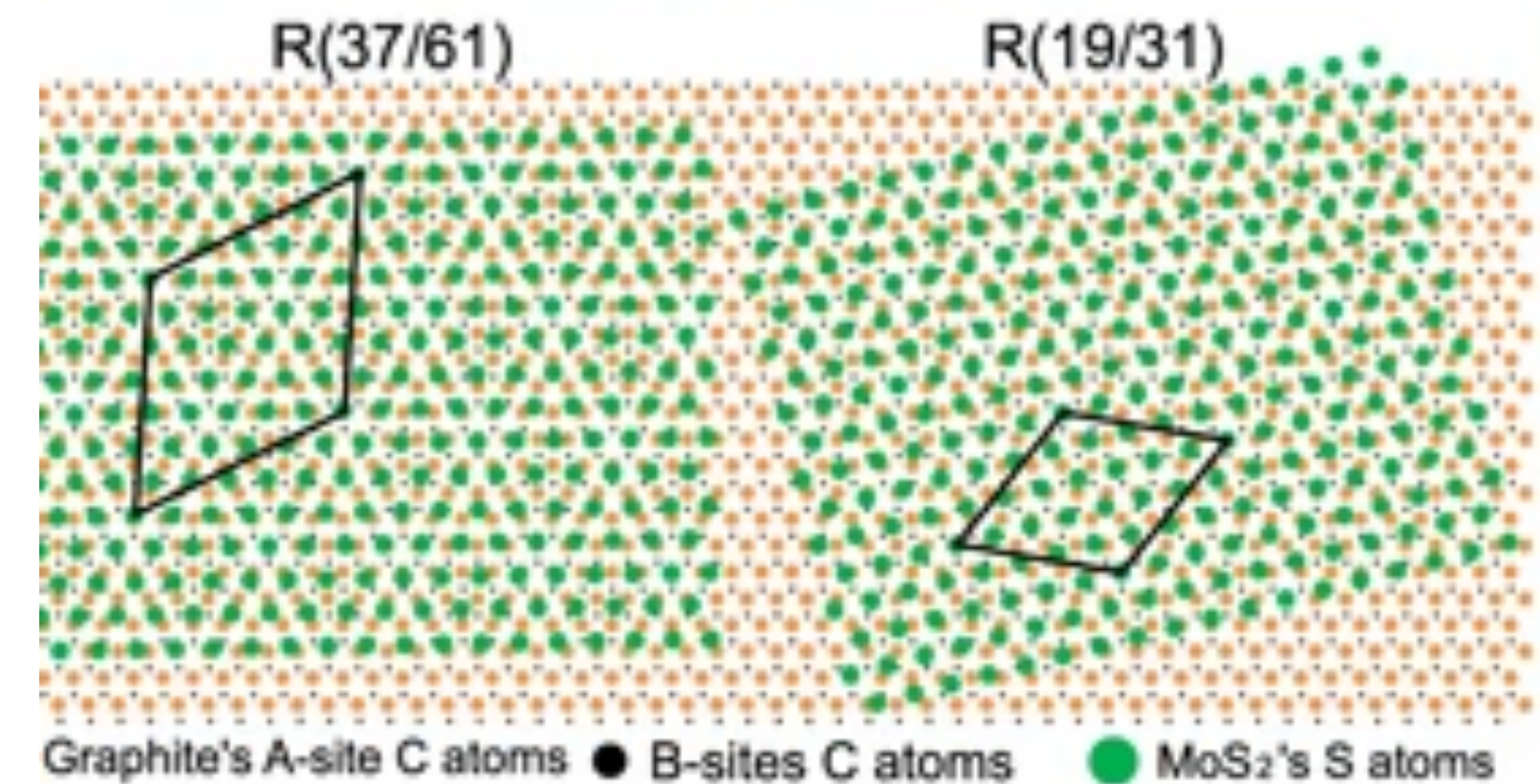
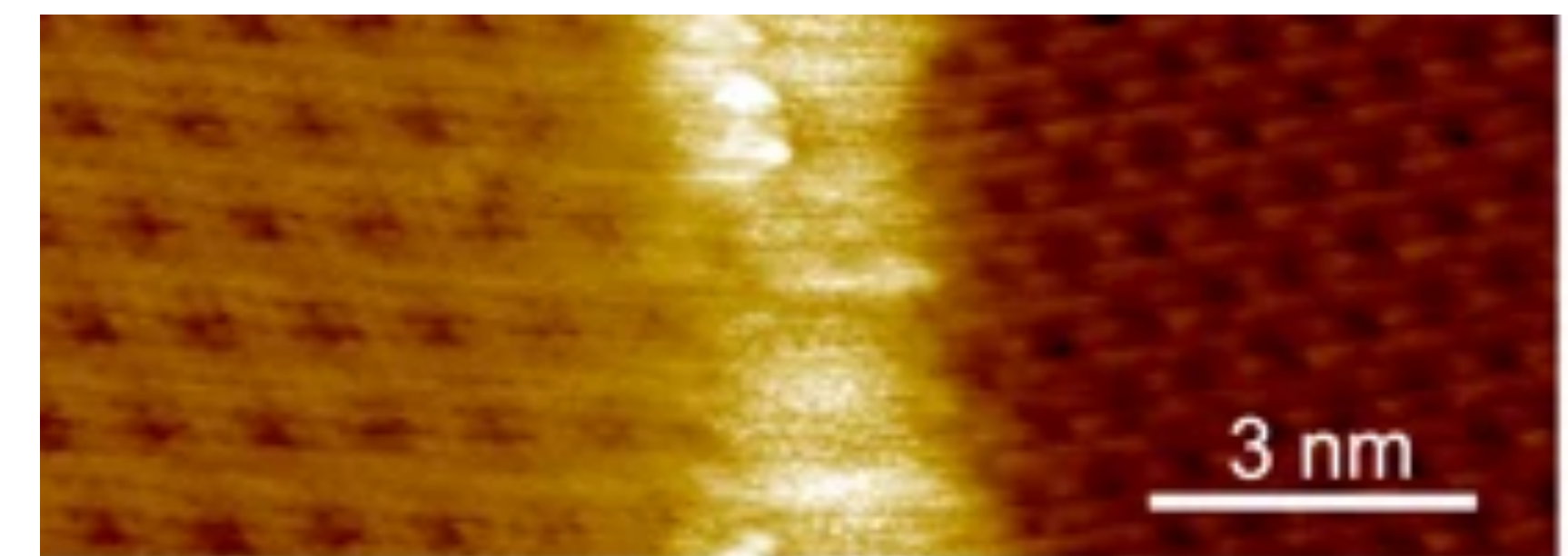
moiré materials

grow material 1 on material 2
if weakly coupled & **mismatch** of lattice constants

use **twisting**: “stamp” material onto
other/same material with rotation angle



Michely group, Cologne

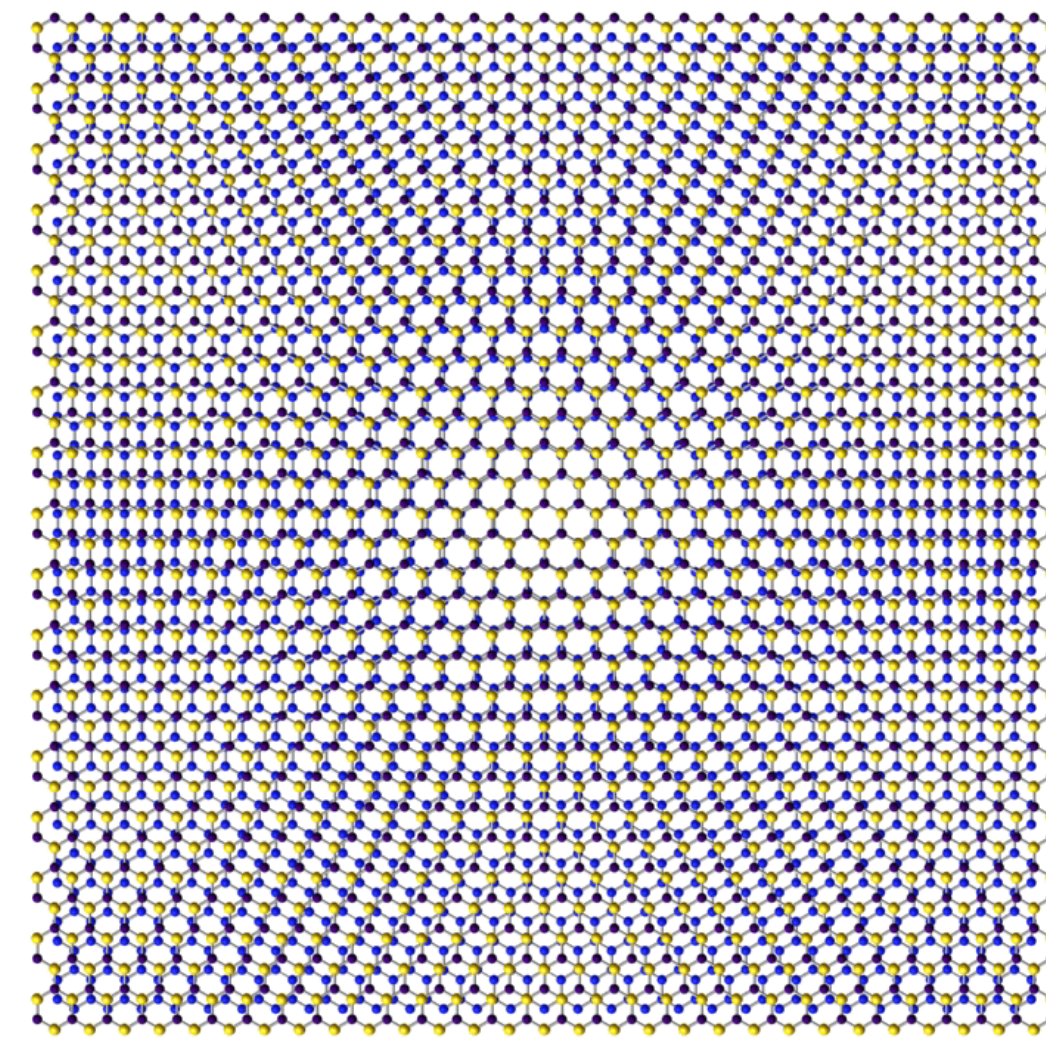
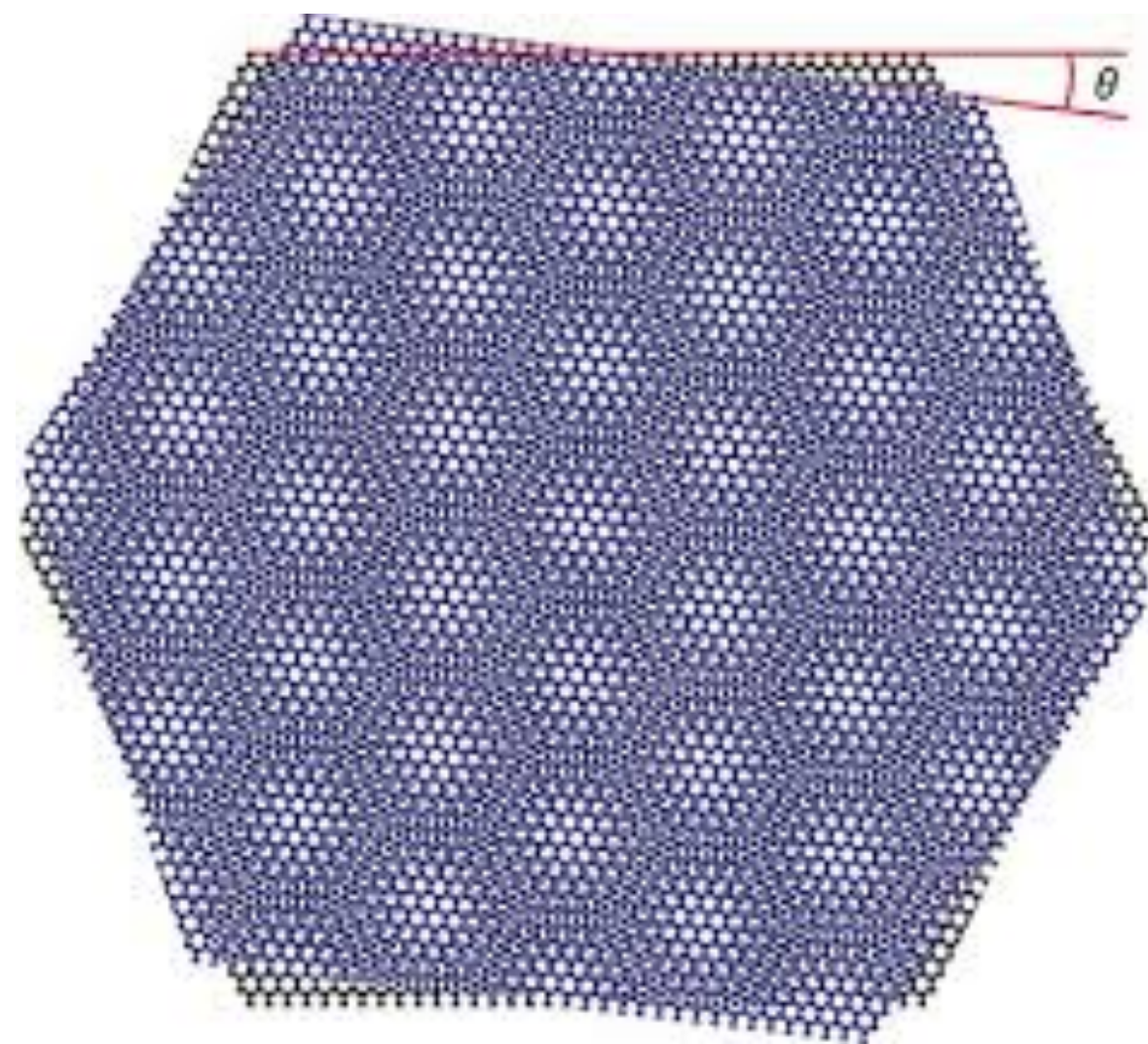


twisted MoS₂/graphite moiré, Chun-I Lu et al. 2017

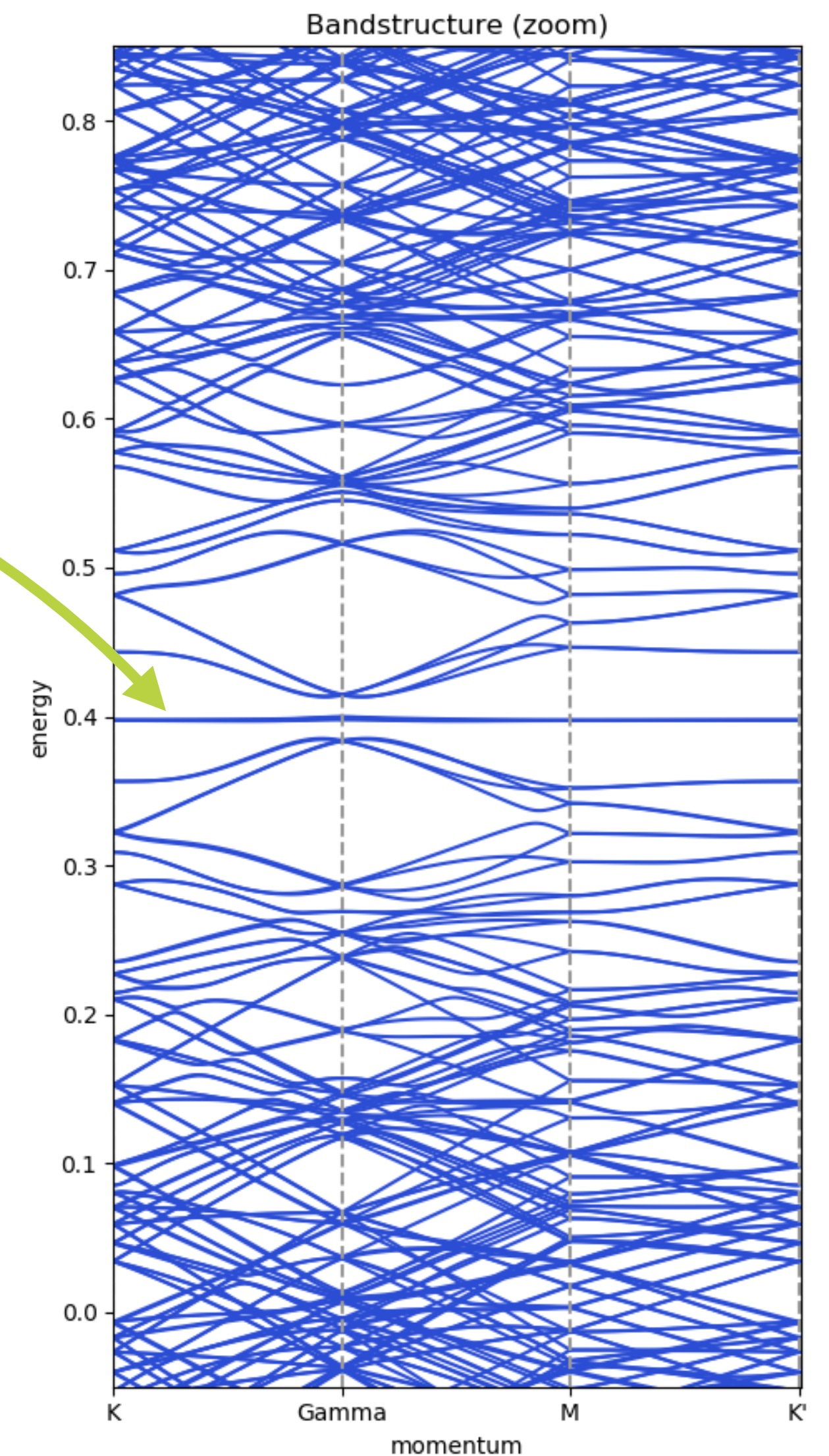
moiré materials

twisted bilayer graphene

- flat bands at **magic twist angle of 1.2°**
due to interference effect
- giant unit cell of 10^4 atoms



www.condmat.physics.manchester.ac.uk/imagelibrary/



moiré materials

twisted bilayer graphene

- flat bands at **magic twist angle of 1.2°**
due to interference effect
- tunable by gate voltage

Mott insulators

superconductors

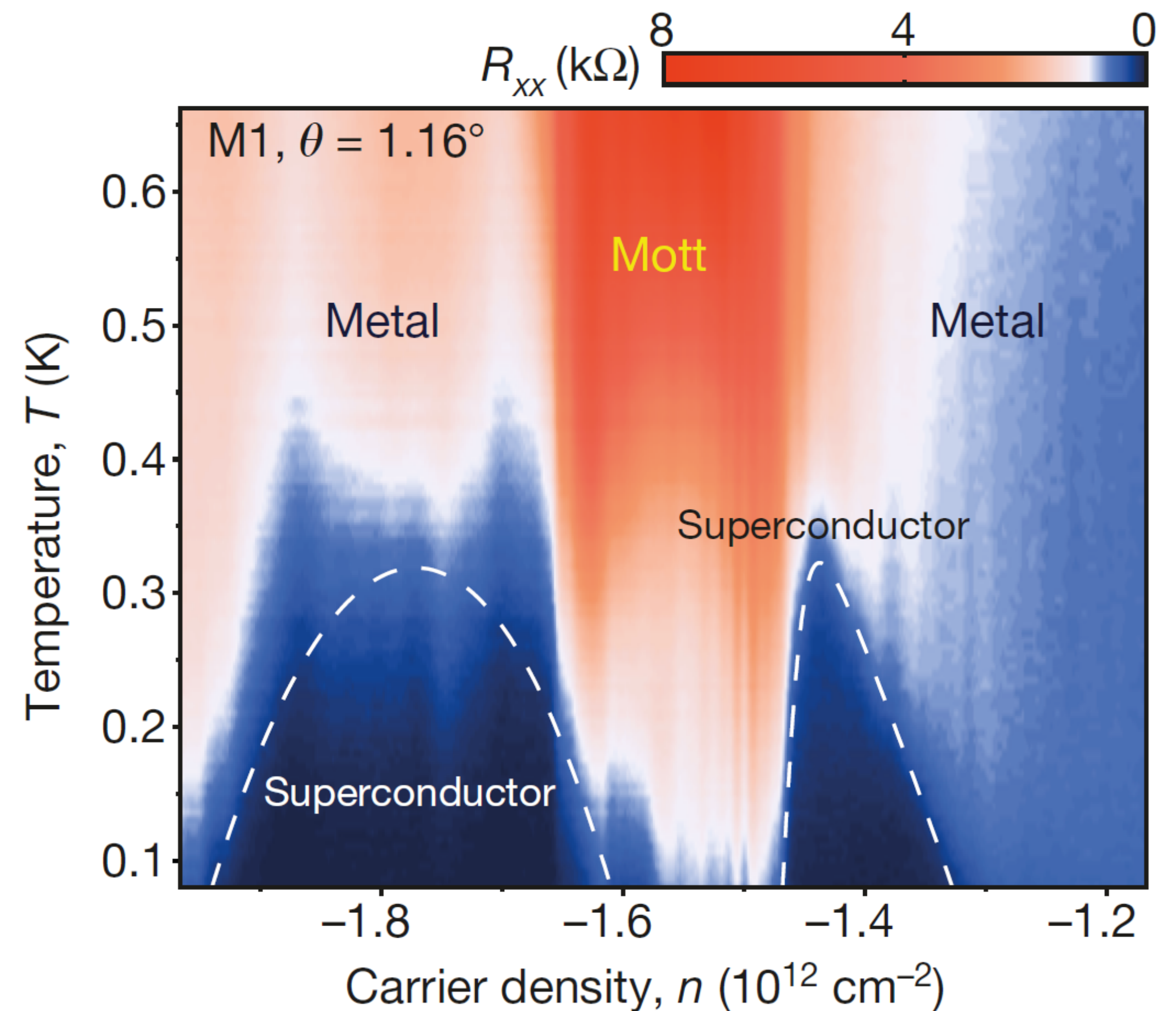
topological bands & anomalous QHE

magnetic phases

nematic order

...

**correlations & topology
in a *single* highly-tunable system**



Cao et al., Nature (2018)

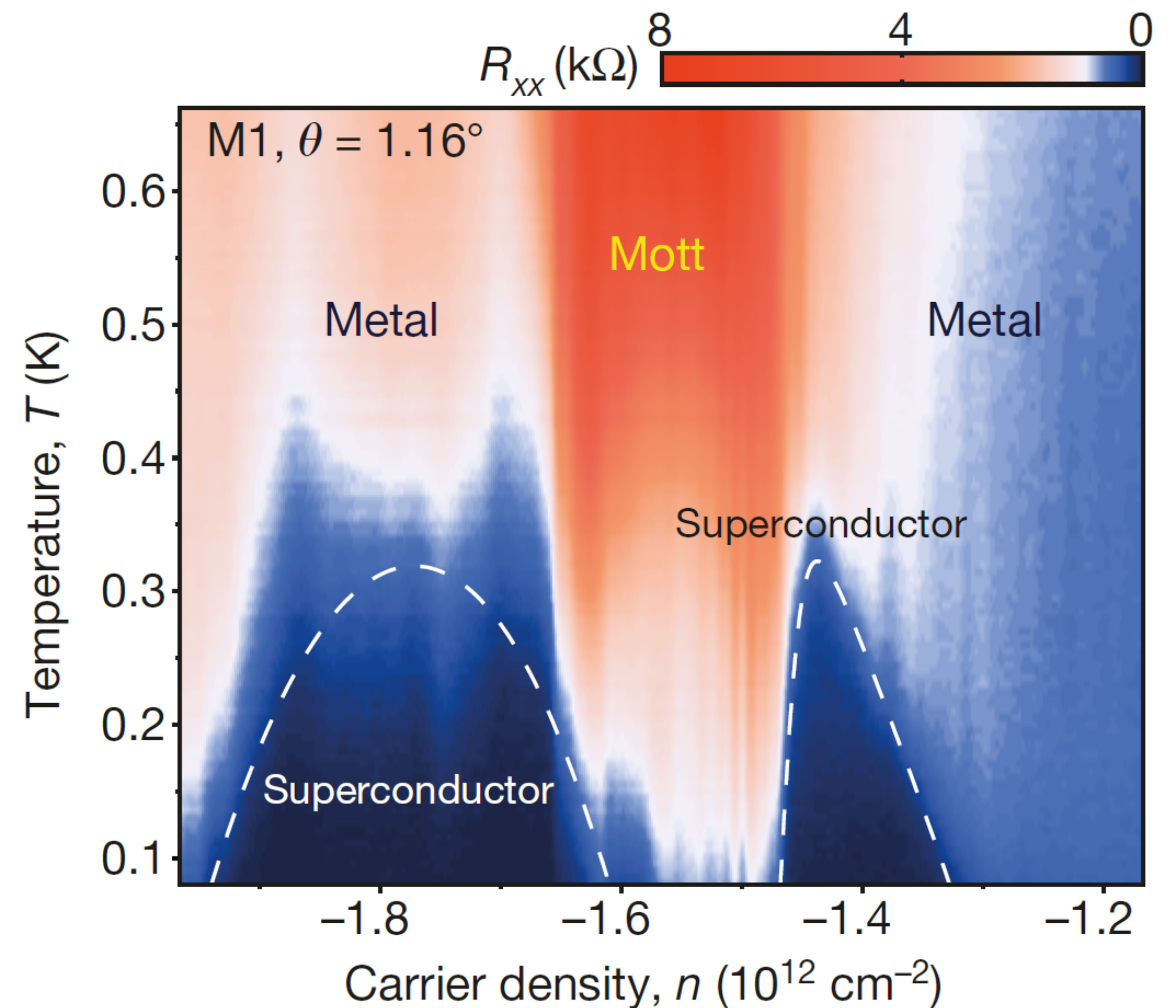
moiré materials

Why moiré systems with giant unit cells?

- easy to add 1 electron per unit cell
⇒ tunable by gate
- additional tunability from twist angle, chemical decoration, ...

This talk

- flat bands natural or “**magic**” needed ?
- fate of all **the other 10^4 bands** ?
- what is **universally** valid ?



Cao et al., Nature (2018)

correlations & topology
in a *single* highly-tunable system

meet the team



Jan Attig



Jinhong Park



Michael Scherer



Alex Altland

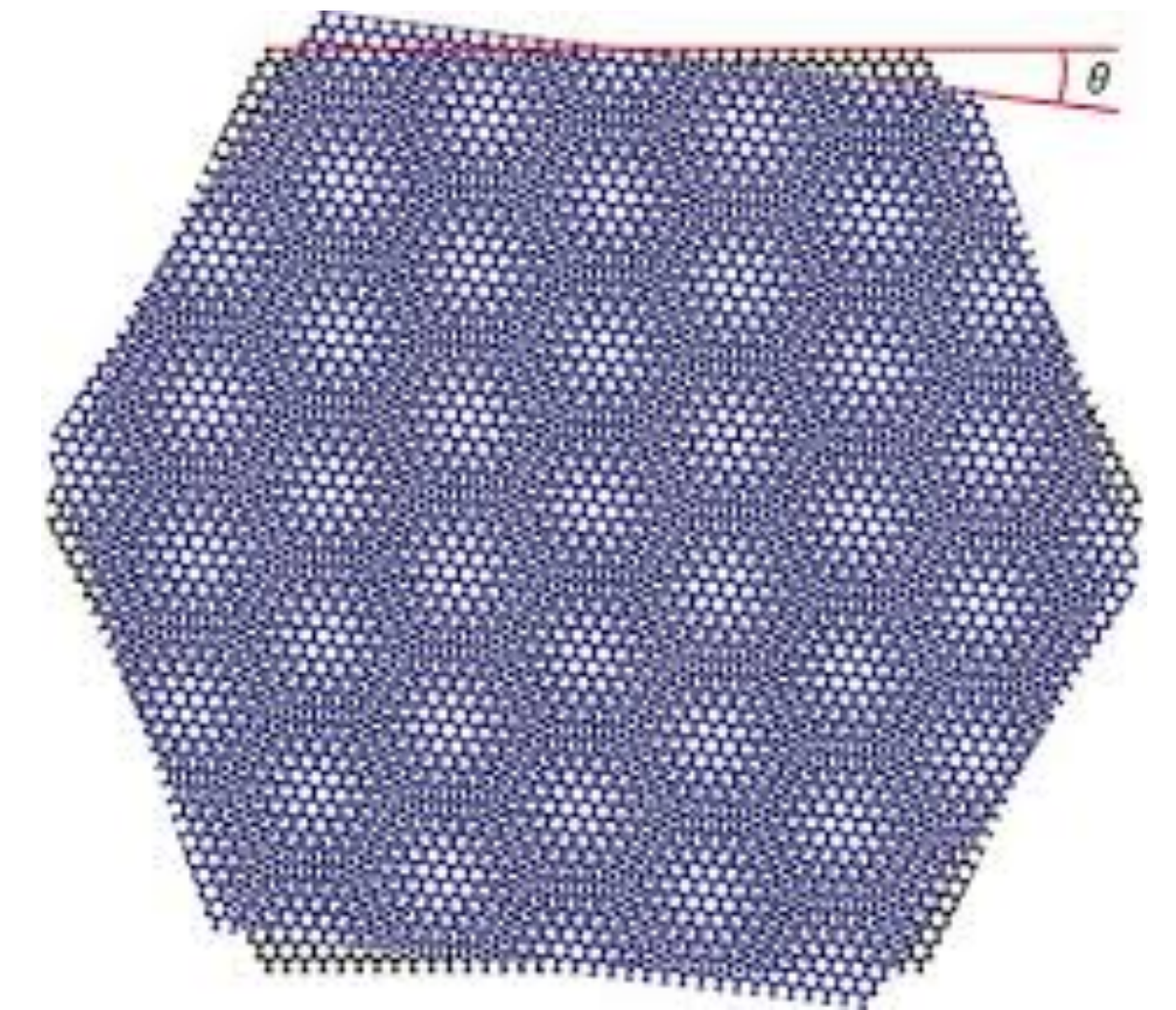


Achim Rosch

basic principles

lattice periodicity

- linear size of moiré unit cell $N \gg 1$
- reciprocal lattice vector $G_M \sim \frac{G}{N}$
- number of atoms = number of bands N^2
- moiré potential \Rightarrow effective **hopping in reciprocal space**



Anderson localization

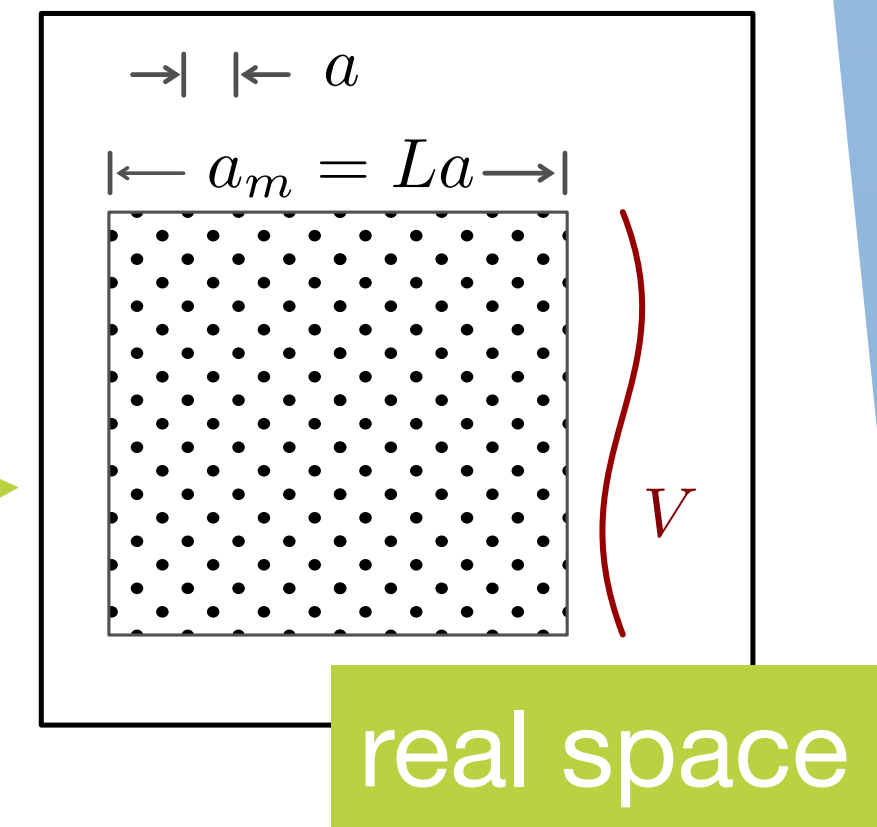
- aperiodic site-to-site variations
- quantum disorder
- dimensional reduction (1D Fermi surfaces)
- **momentum space localization**
 \Rightarrow strong band dispersions

quantum chaos

- dimensional crossover 1D \Rightarrow quasi-2D
- **momentum space delocalization**
- ergodicity hampered
by discrete lattice symmetries

momentum space phenomenology

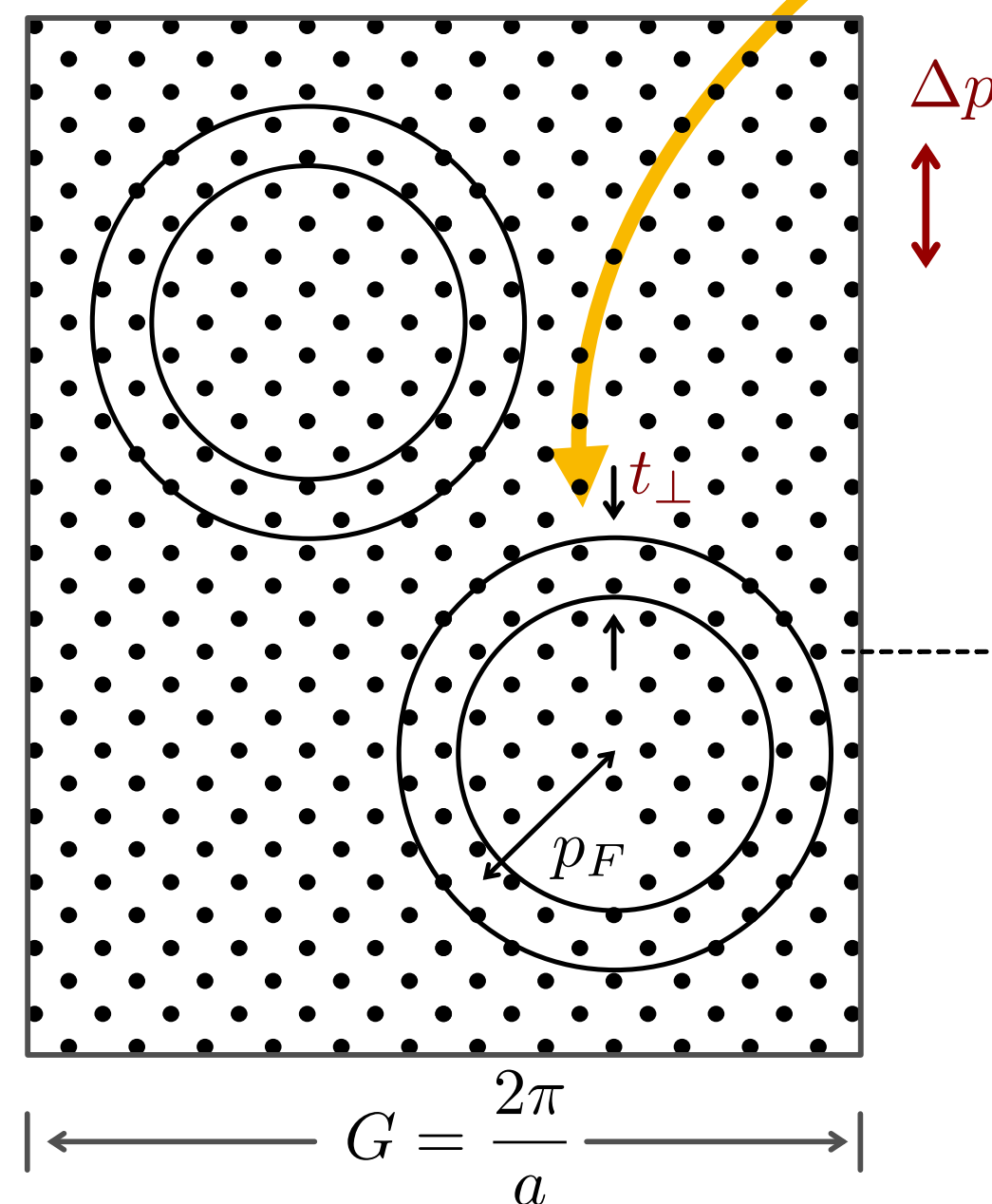
- 2D crystalline lattice subject to a **perturbation** V periodic over distances $N \gg 1$
- V defines periodic “**hopping potential**” in momentum space



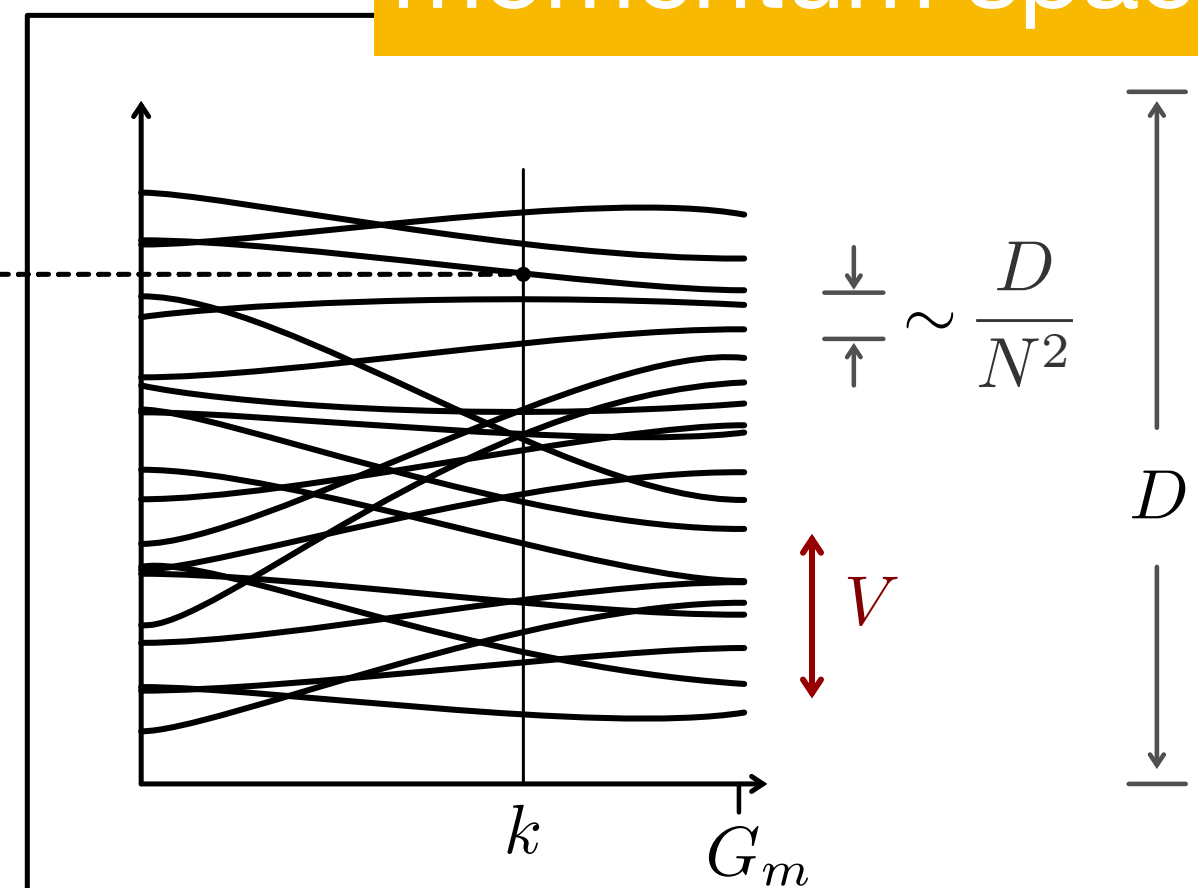
$$G_m = \frac{2\pi}{a_m}$$

→ | ←

- **quasi-1D Fermi surfaces** (of width t_{\perp})
- effective **disorder**



momentum space



Anderson localization

numerical simulations



Jan Attig

real space

twisted bilayer graphene in real space
with parameters:

- twist angle
- distance of graphene layers
- strength of corrugation

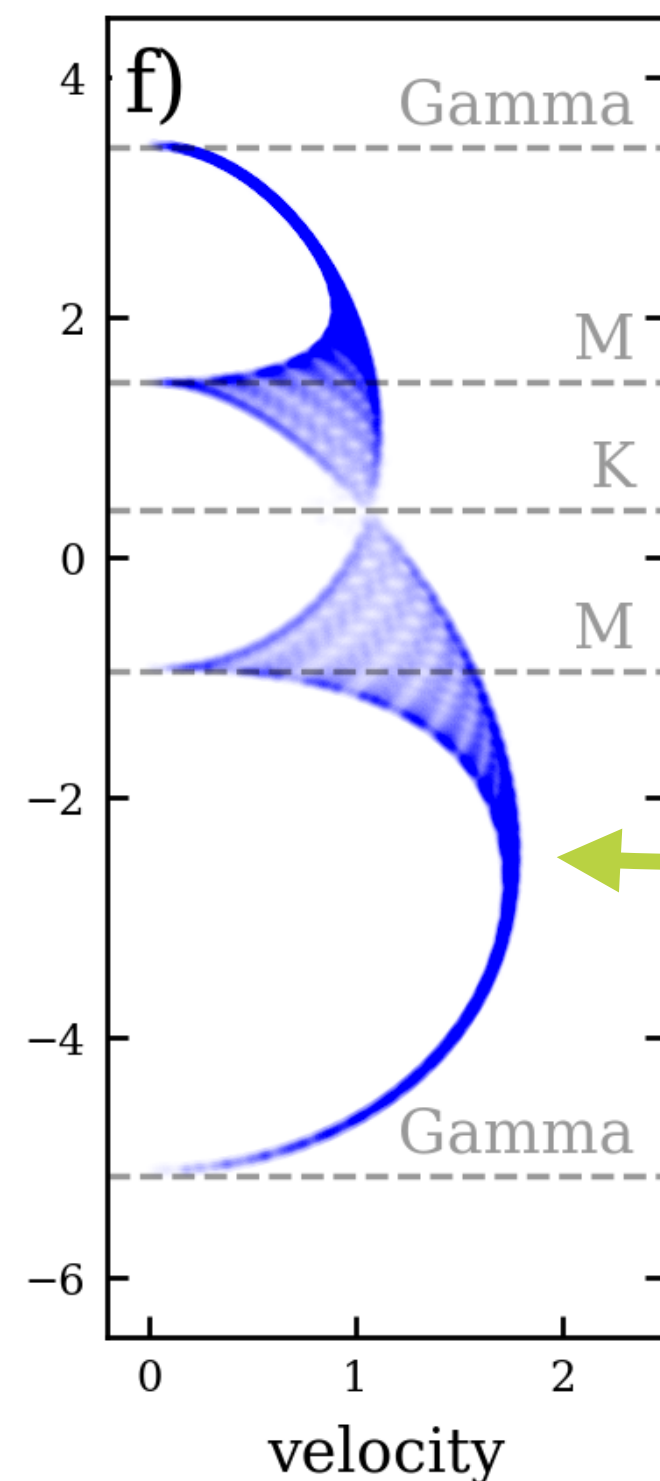
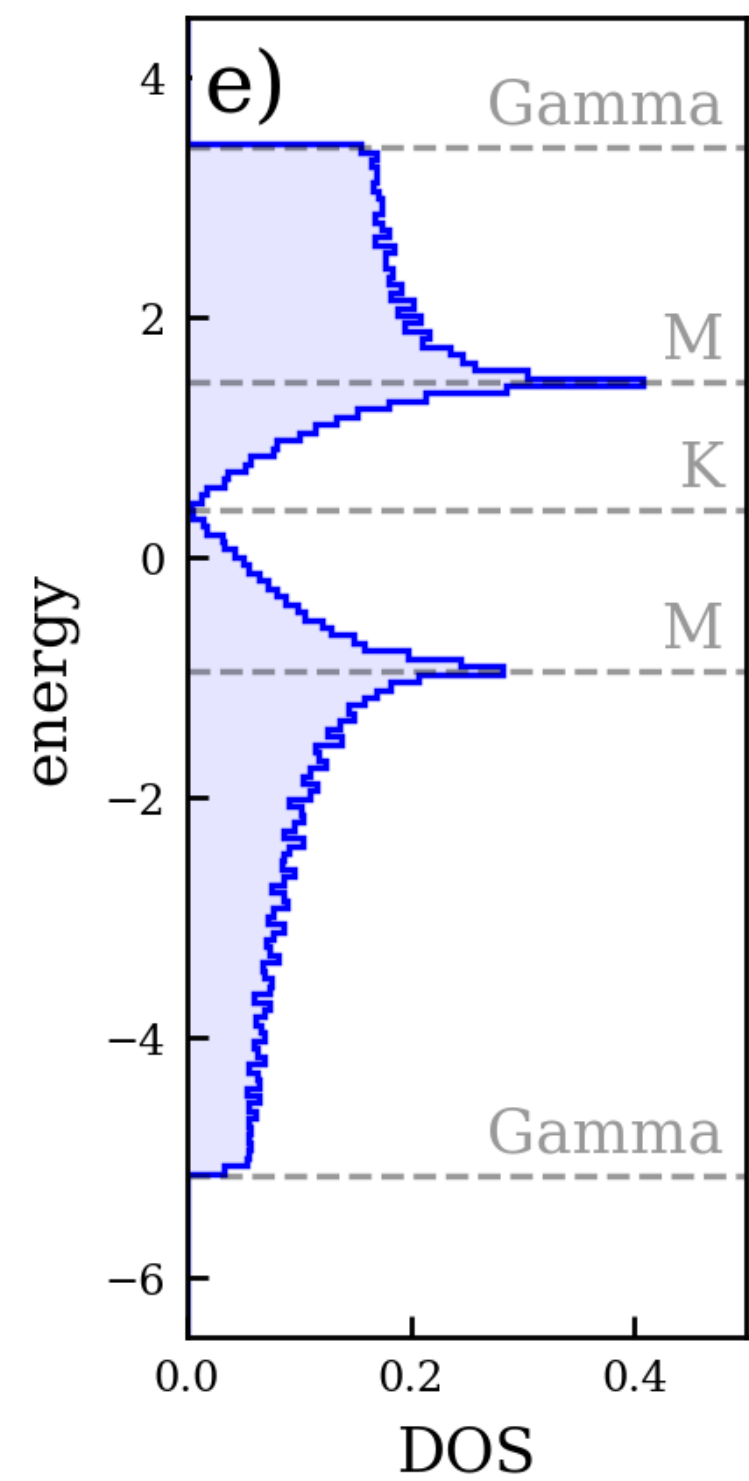
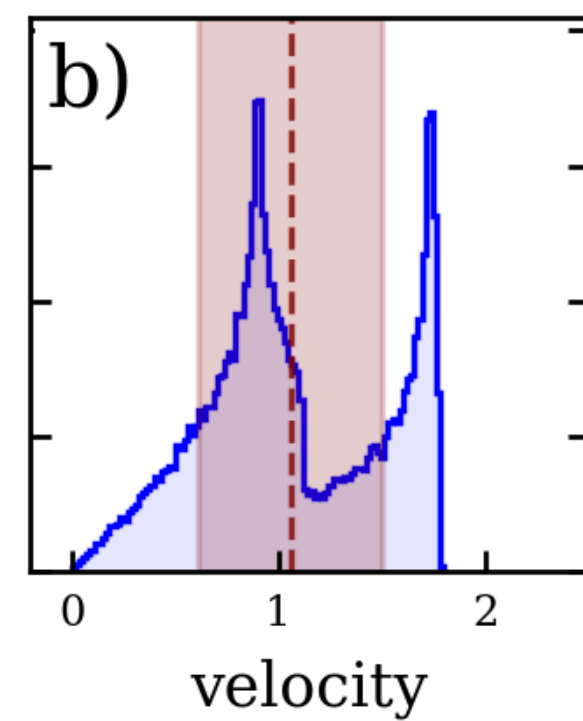
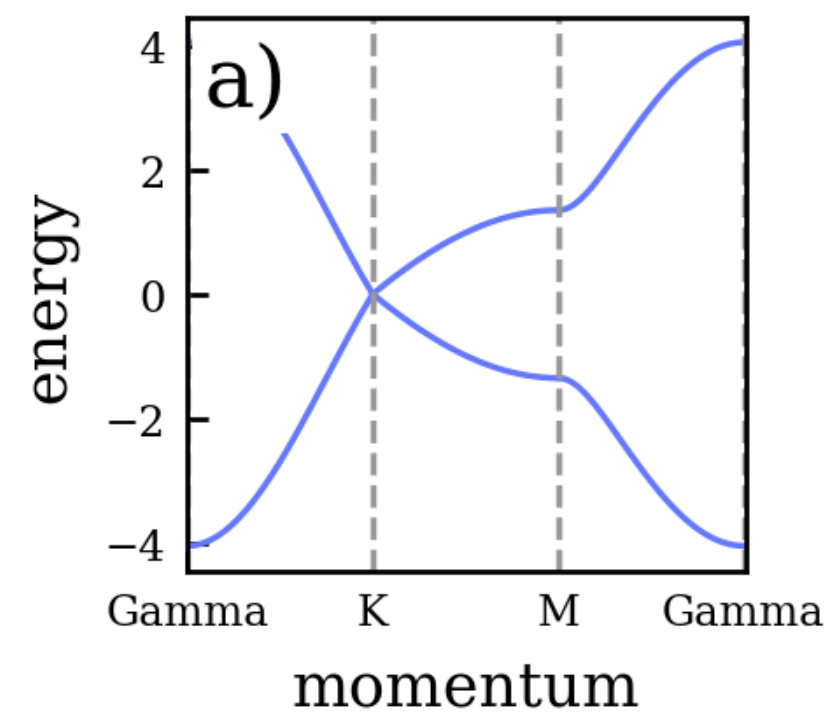
momentum space

momentum space code
based on
continuum model



Jinhong Park

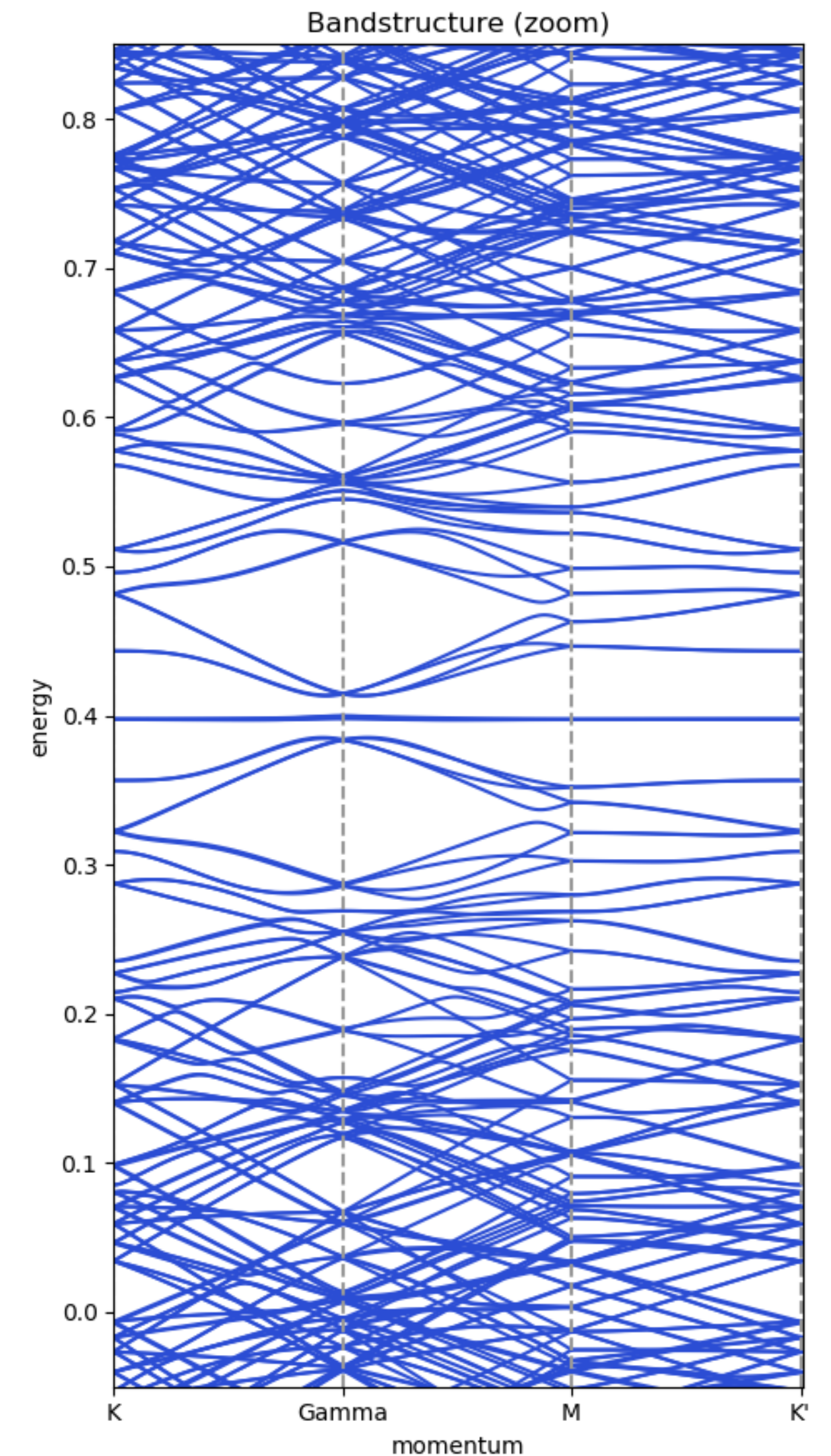
real-space numerical simulations



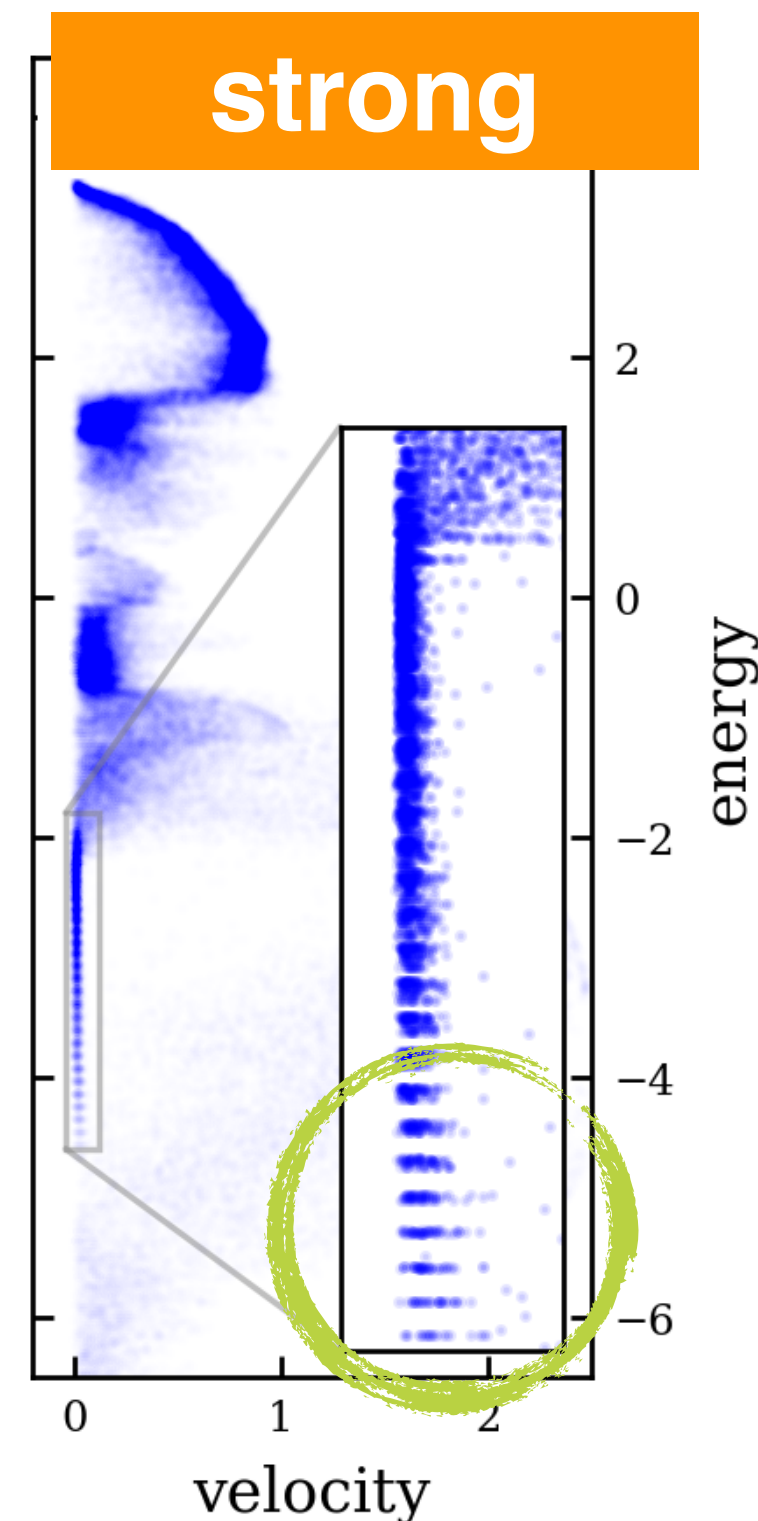
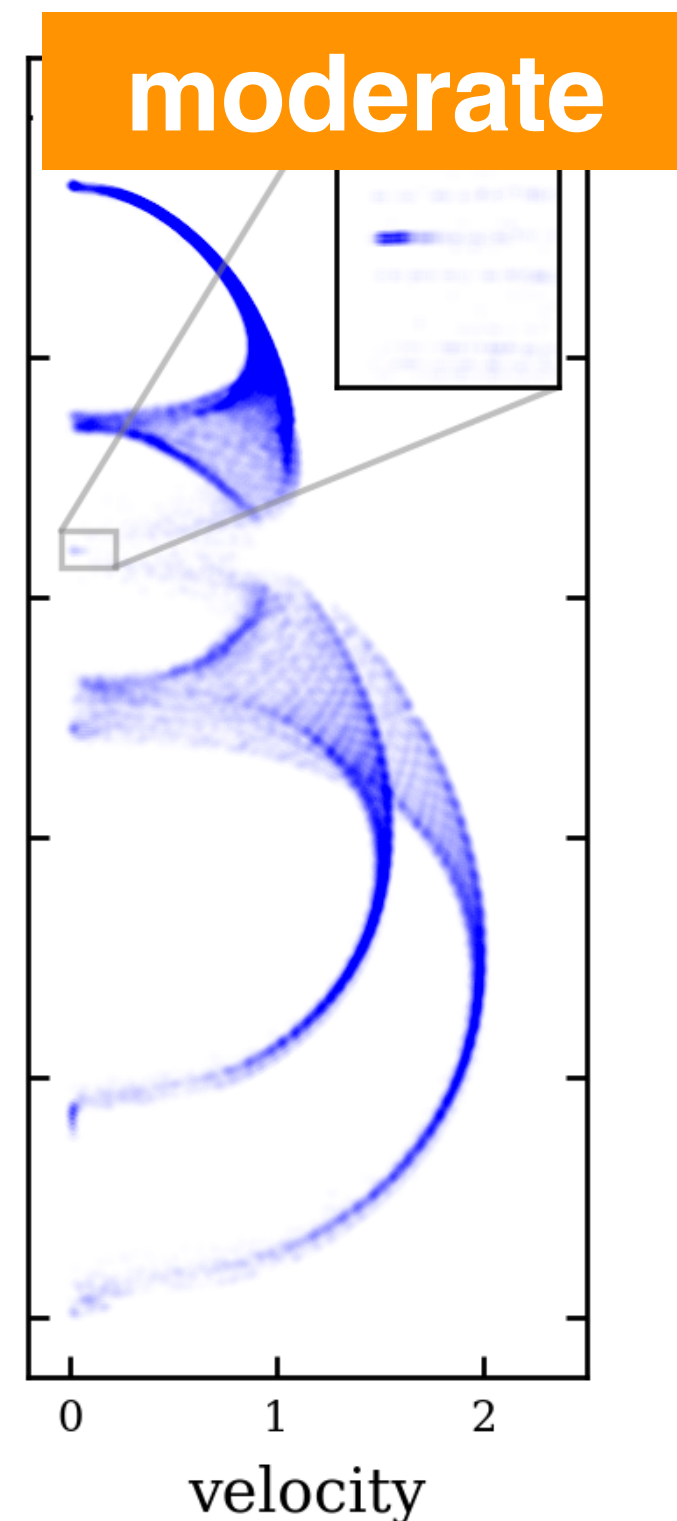
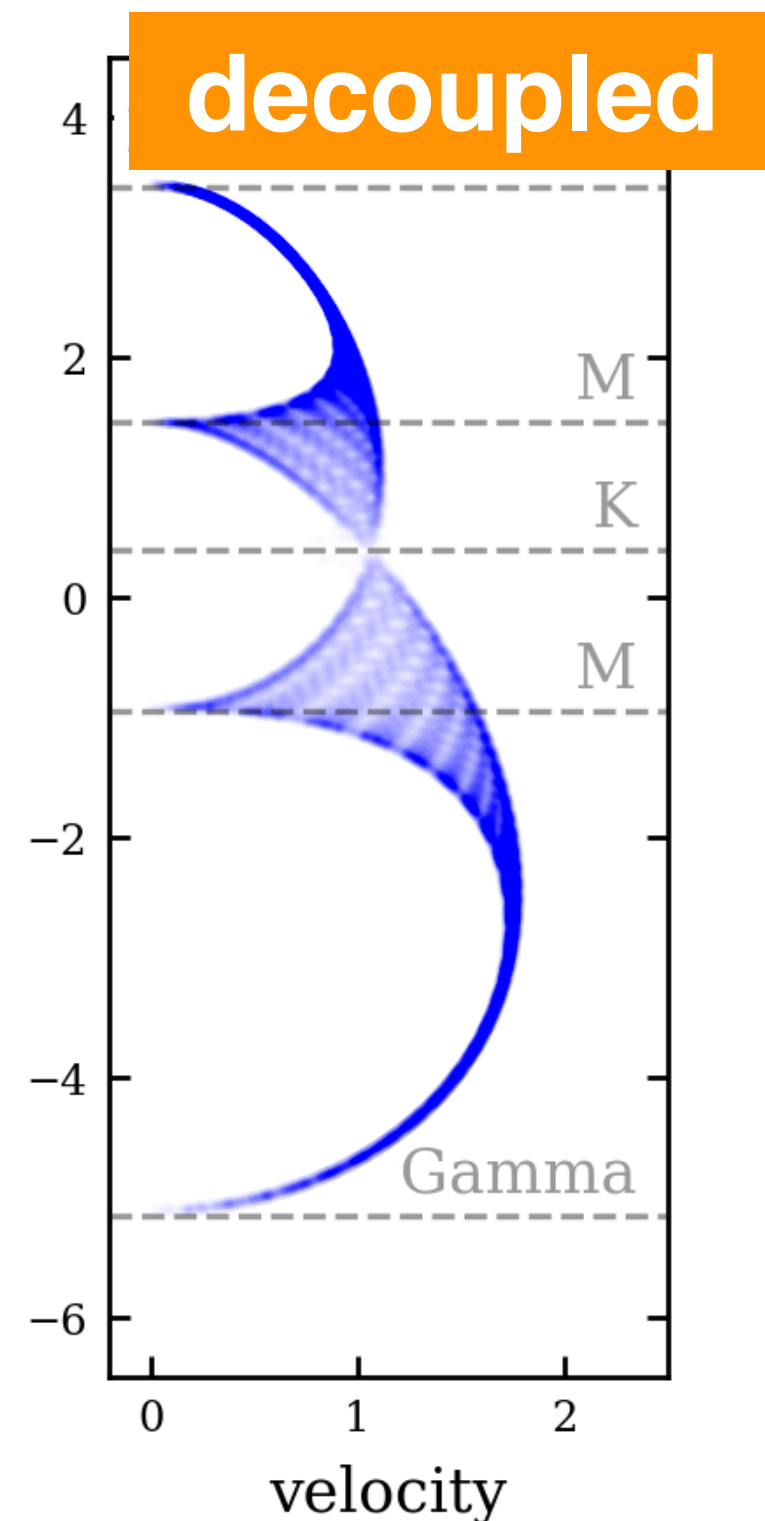
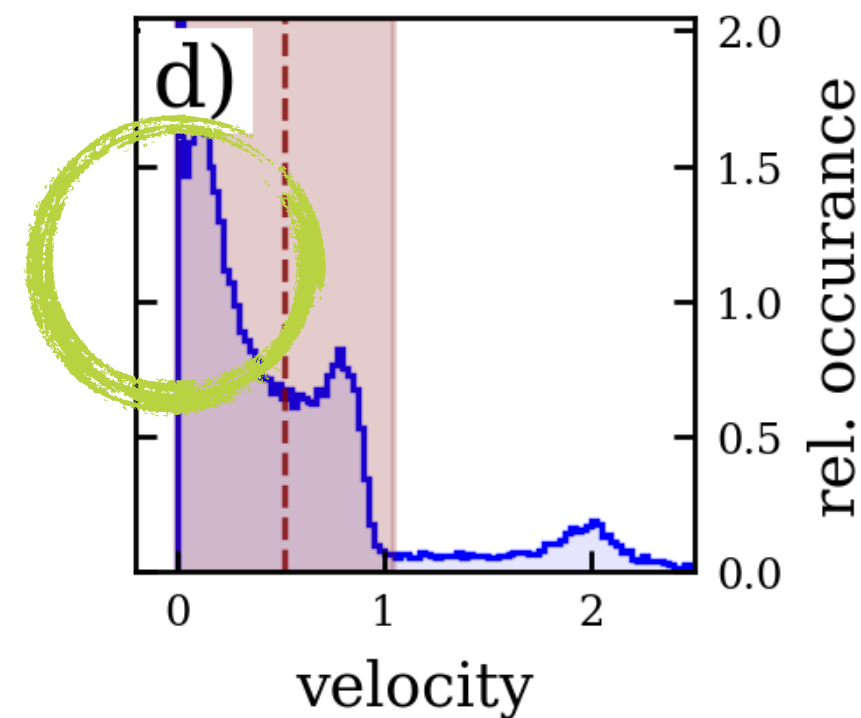
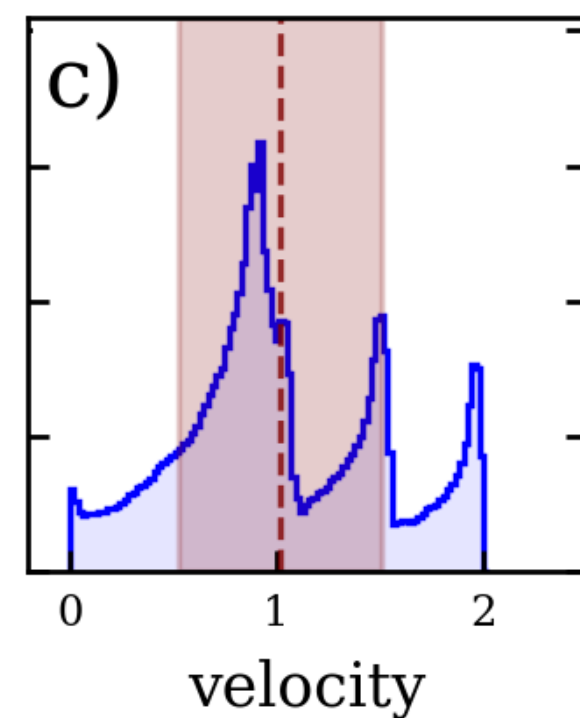
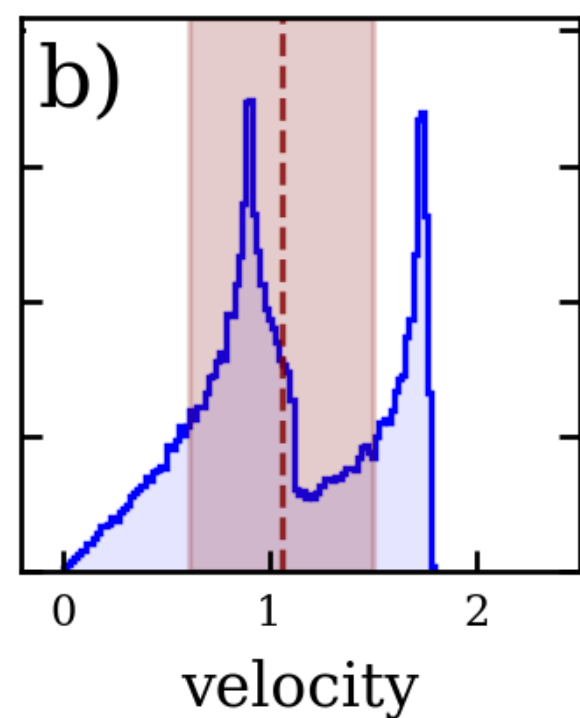
distribution of **velocities**

decoupled layers

“waterfall” plots

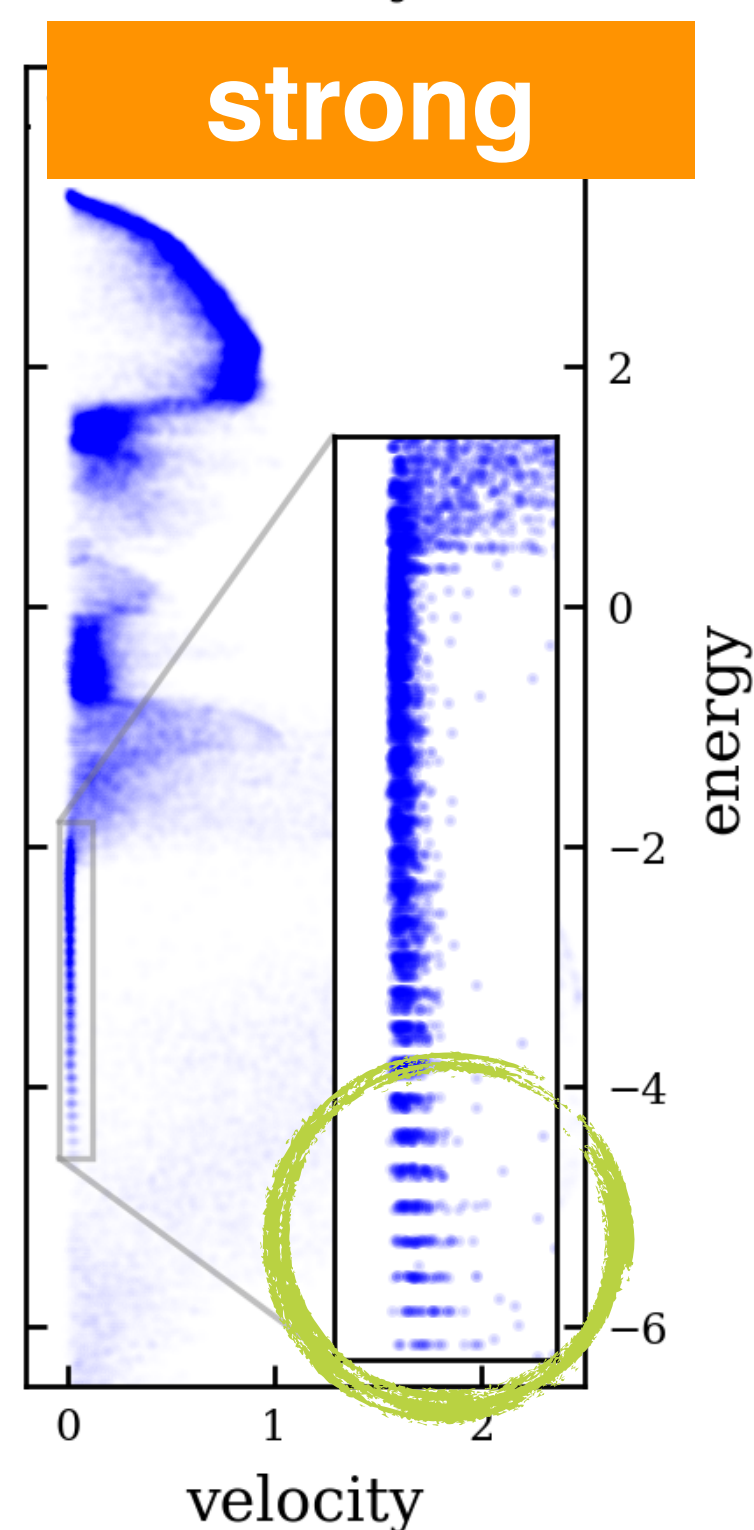
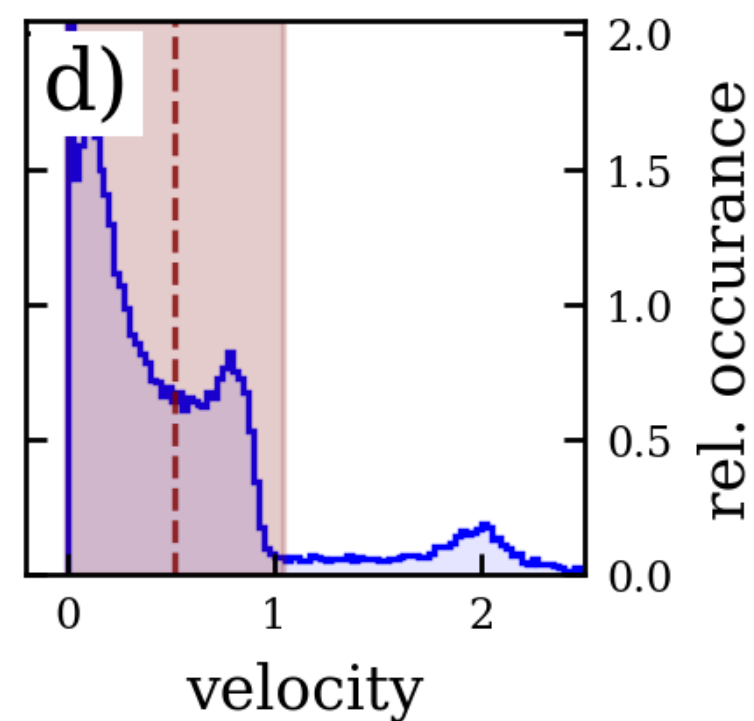


real-space numerical simulations

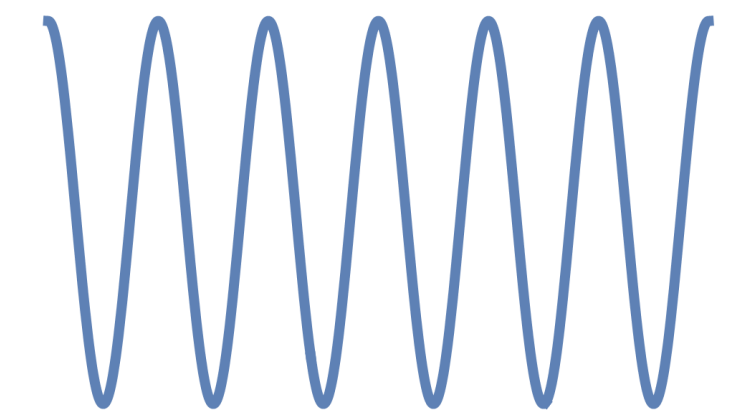


- **typical** velocities of **$O(1)$** , not $O(1/N)$
- **enhanced** probability for small v ?
- why are some regions unaffected ?
- **equally spaced**, very flat bands ?

real-space numerical simulations



- close to minimum/maximum of graphene band, map to tunnelling in a potential $t_{\perp} \cos(Q_M x)$



harmonic oscillator states

- potential is large

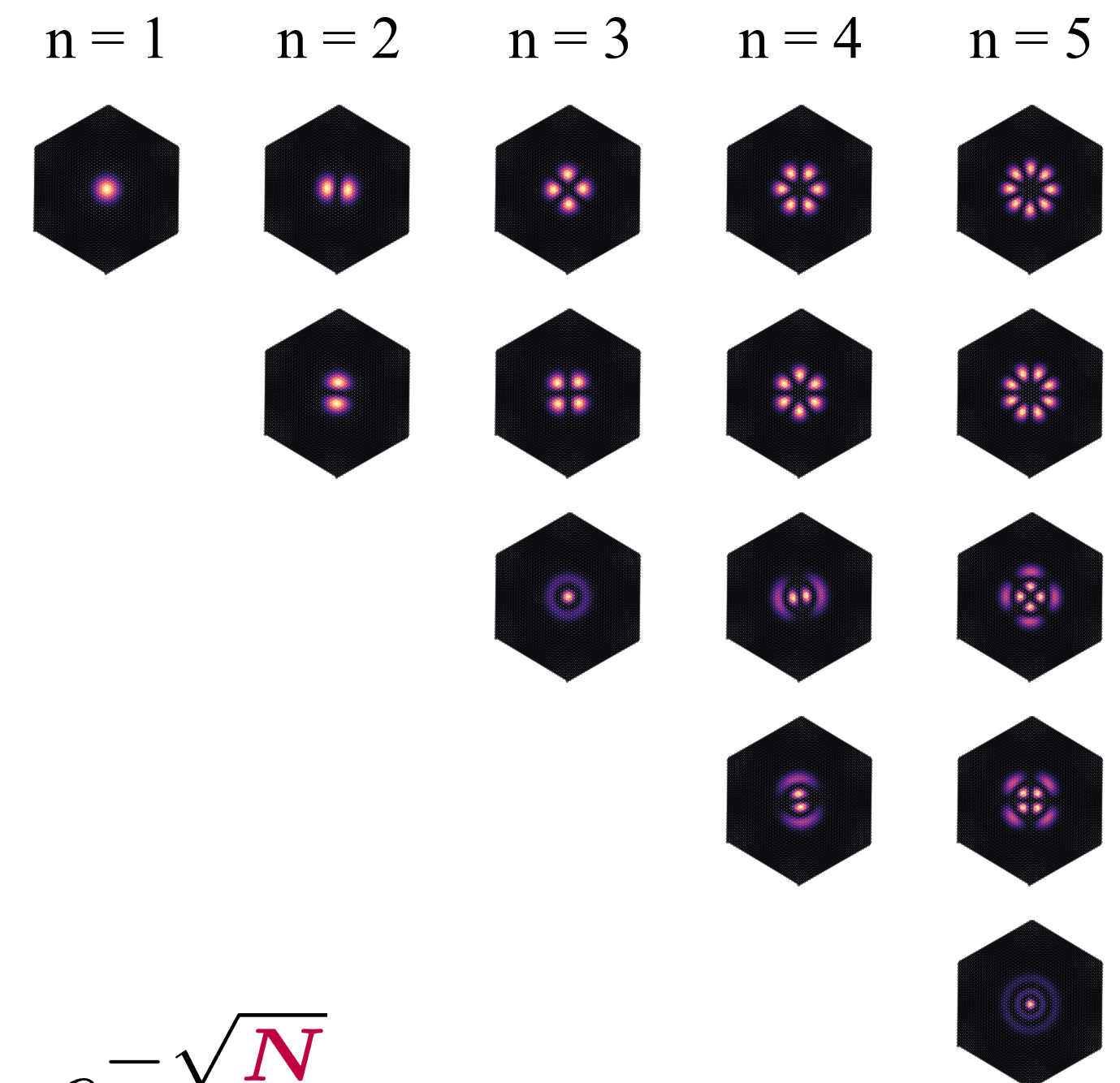
$$t_{\perp} \sim \frac{1}{N} \gg \frac{Q_M^2}{2m} \sim \frac{1}{N^2}$$

- harmonic oscillator spacing

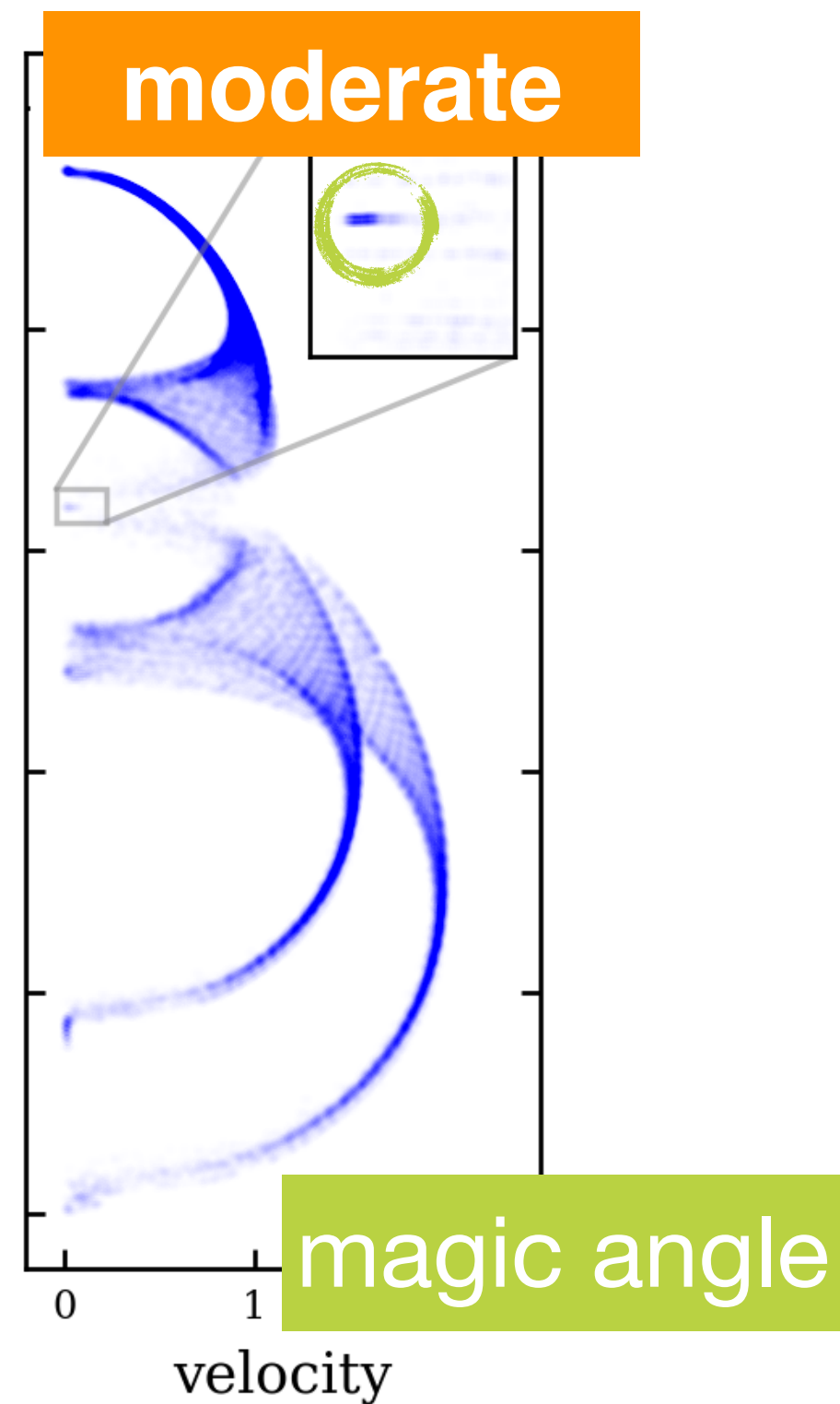
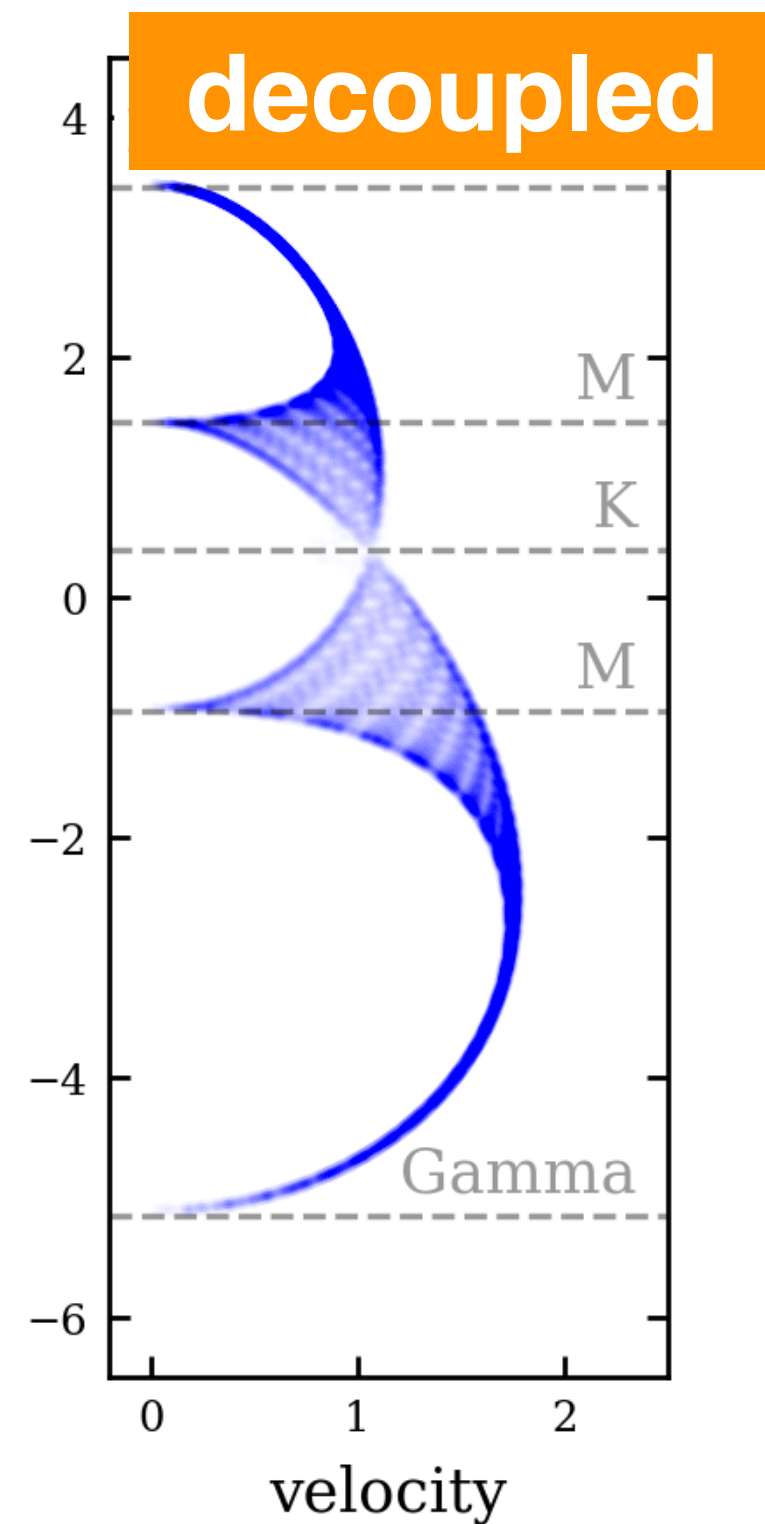
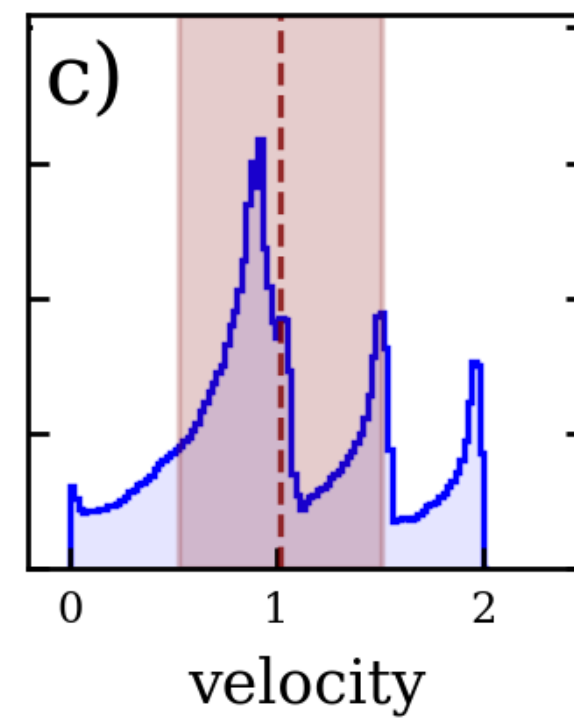
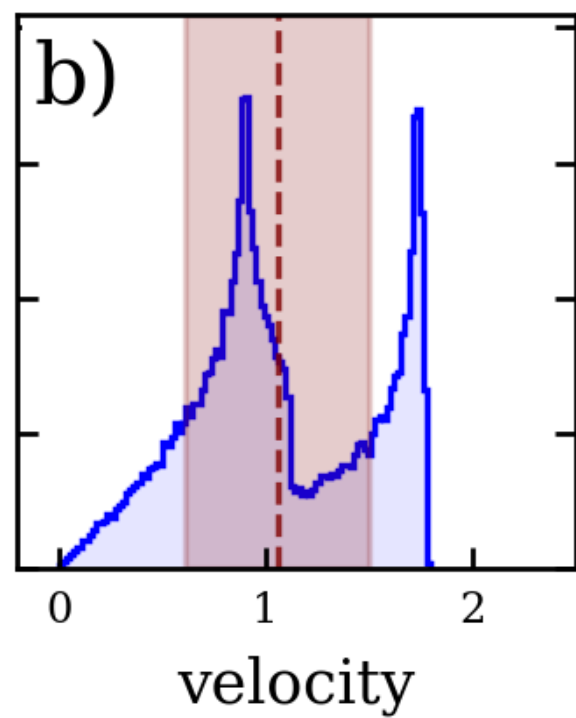
$$\sqrt{\frac{t_{\perp} Q_M^2}{m}} \sim \frac{1}{N^{3/2}} \gg \frac{1}{N^2}$$

- exponentially small bandwidth

$$e^{-\sqrt{mt_{\perp}} Q_M} \sim e^{-\sqrt{N}}$$



real-space numerical simulations



- Why is there **almost no effect** of tunnelling for most energies ?



**momentum space
localization**

momentum-space dynamics

- **dynamics in momentum space:** lattice points spanned by reciprocal moiré lattice
- **tunnelling** between graphene layers or **scattering** from moiré potential

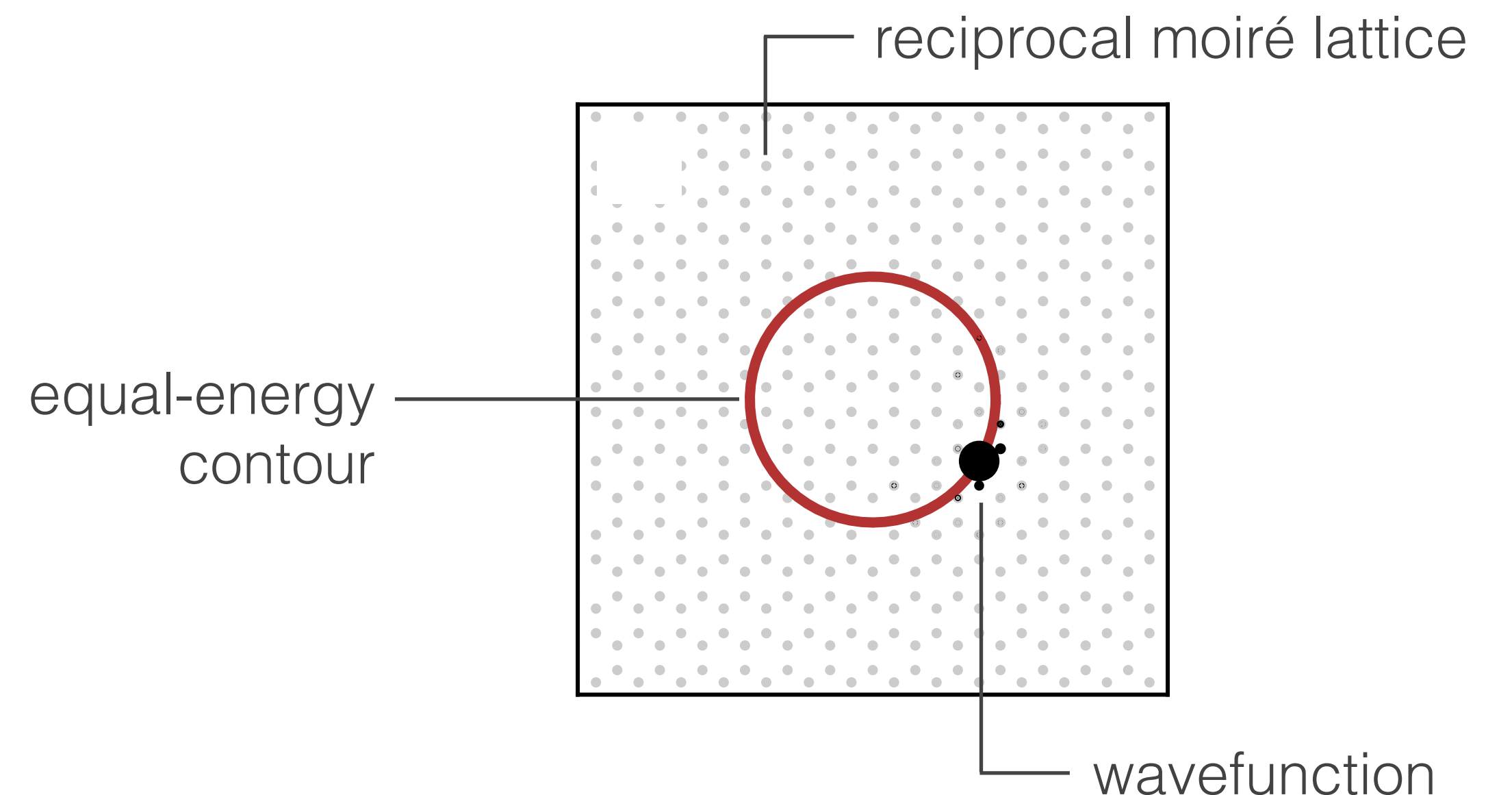
hopping in reciprocal space

- graphene **band structure**

potential term in k-space

- tunnelling along **equal-energy contours** (circles)

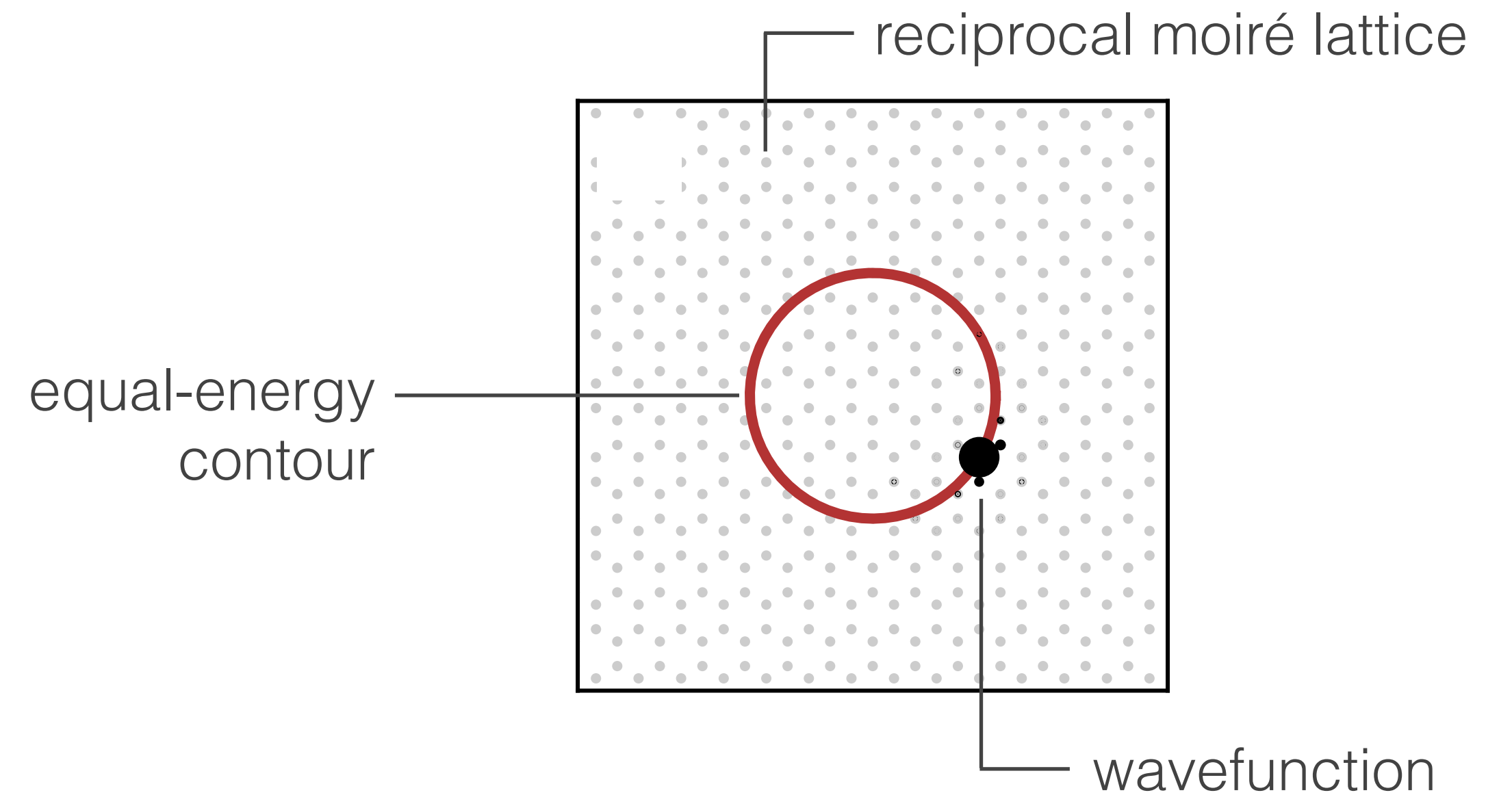
$$v_F |k| = E \pm t_{\perp}$$



momentum-space localization

Anderson localization

- **localization** by “effective disorder”
- localization **in 1D** highly efficient
 - └ why 1D?
hopping along Fermi surface
(equal-energy contour)
- localization **length**
 - = mean-free path times # of channels
- **velocity**
 - = weighted average of underlying Fermi velocities



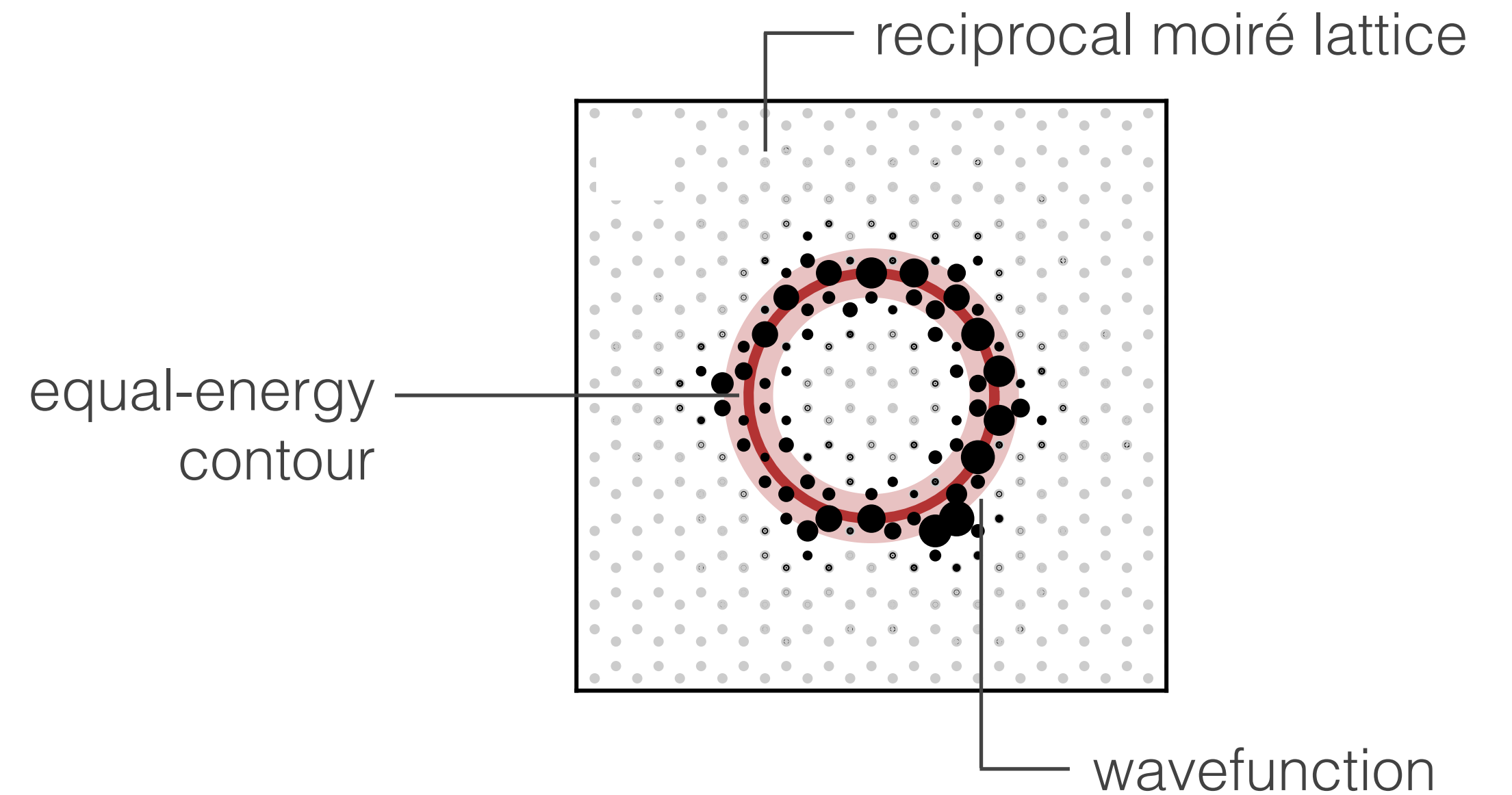
localization in k-space

⇒ no level repulsion
high velocity

momentum-space localization

Anderson localization

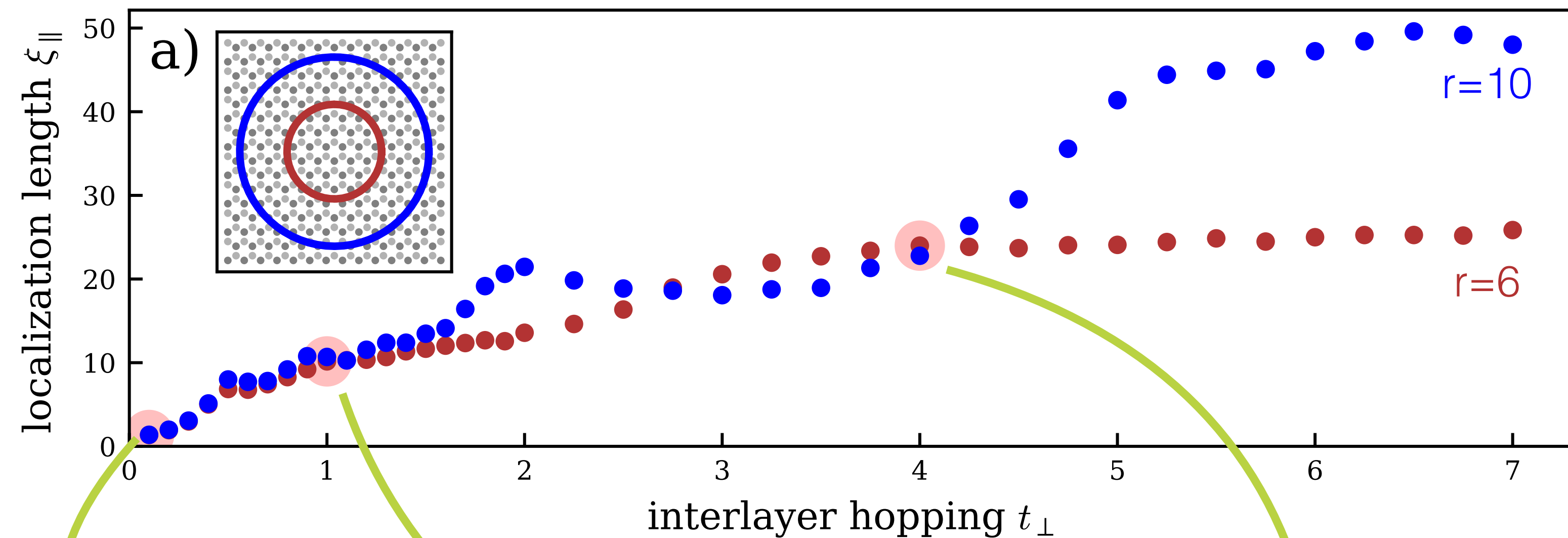
- **localization** by “effective disorder”
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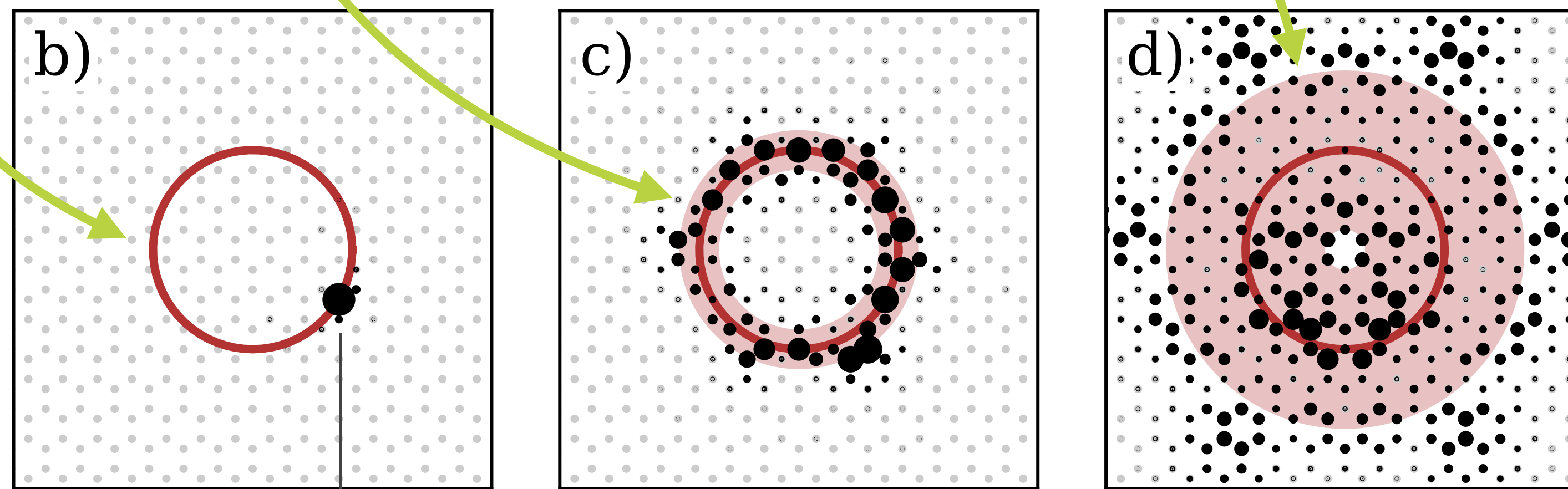
partial localization in k-space

⇒ level repulsion
reduced velocity

momentum-space localization

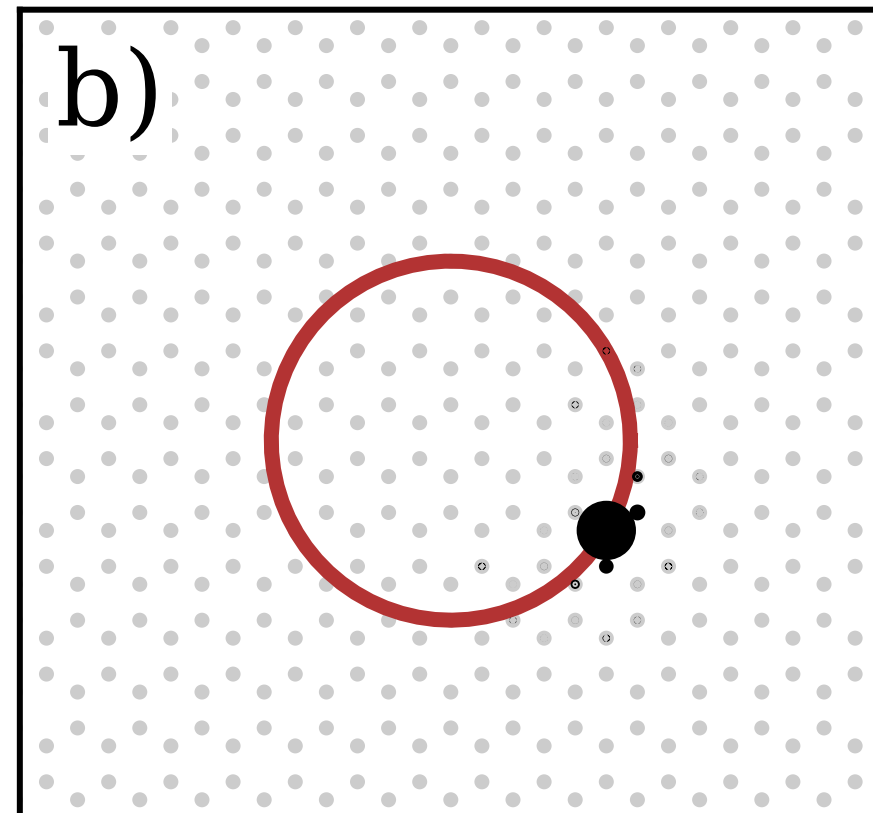


localization driven
by **interlayer hopping**

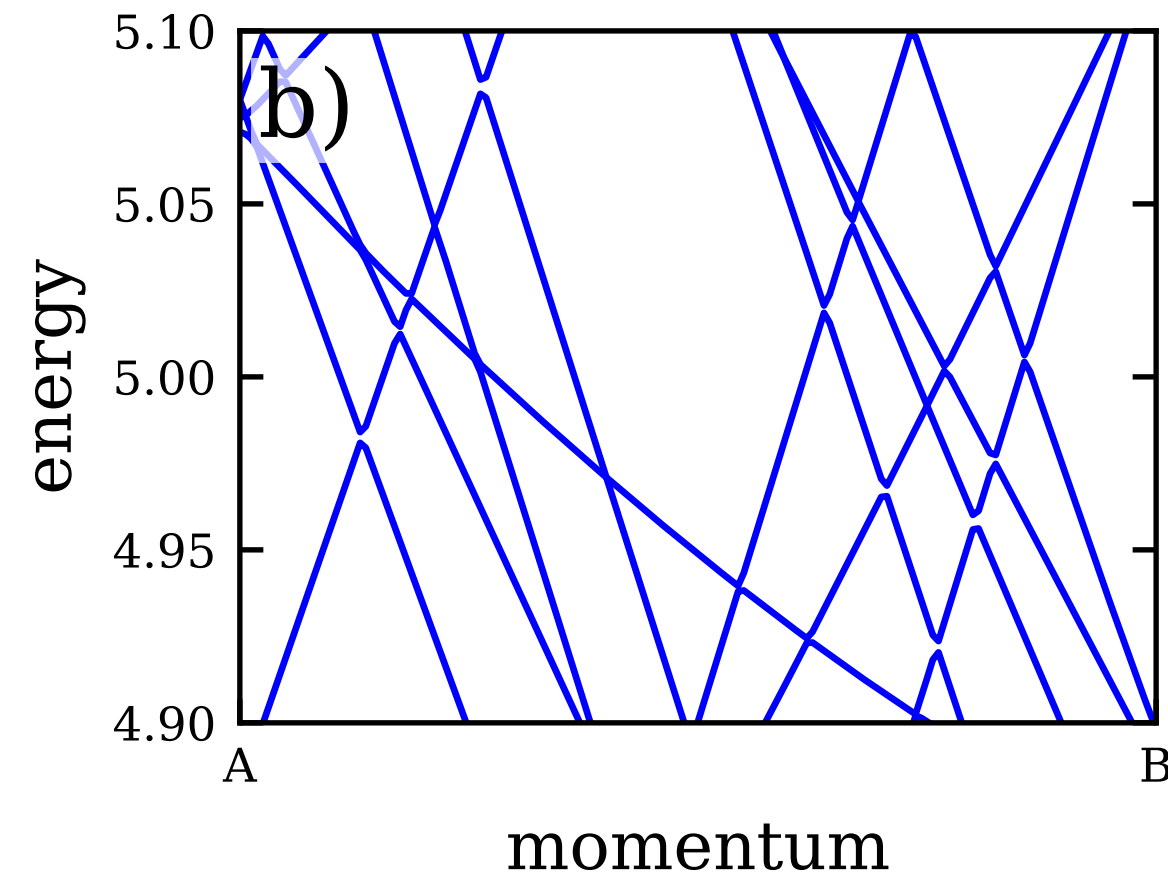


wavefunction

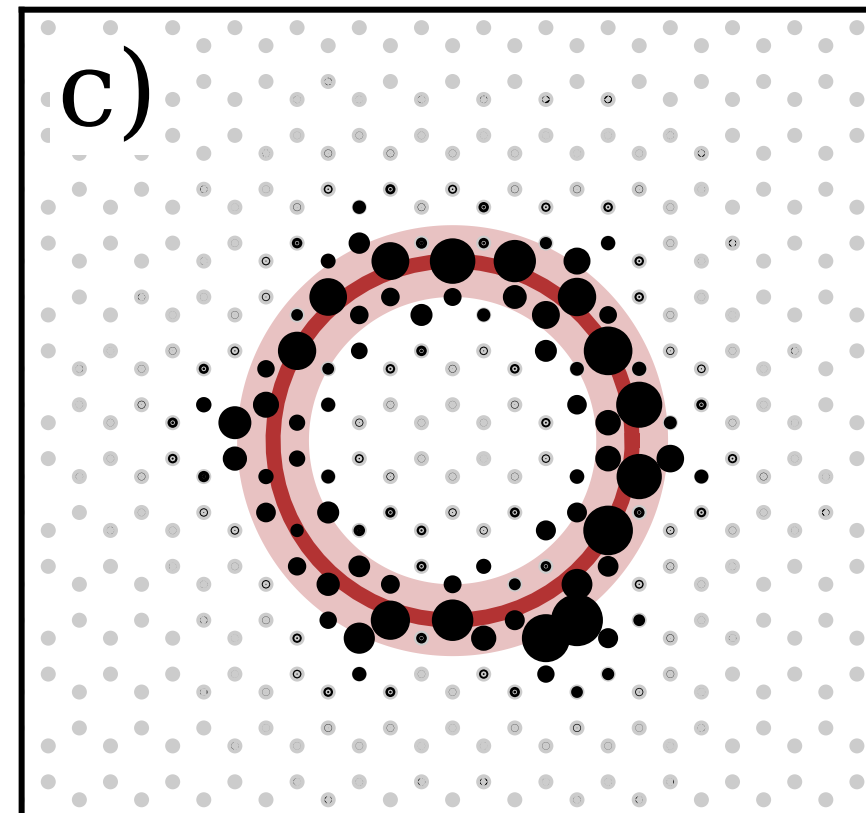
three localization regimes



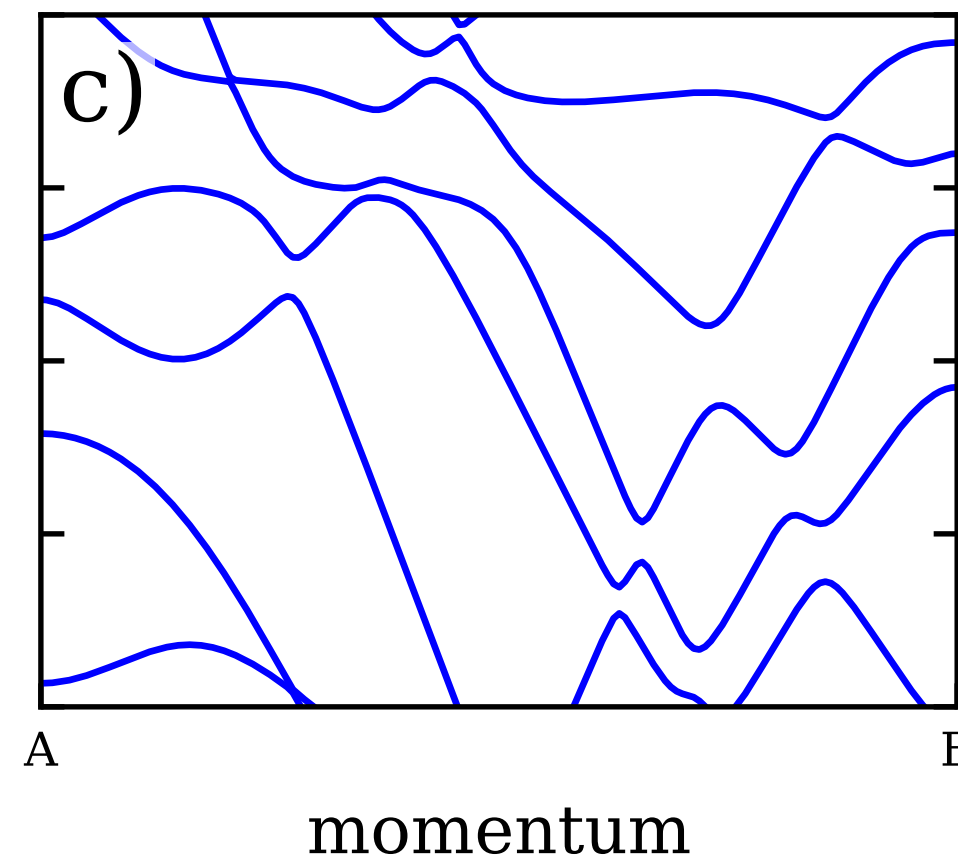
deep localization



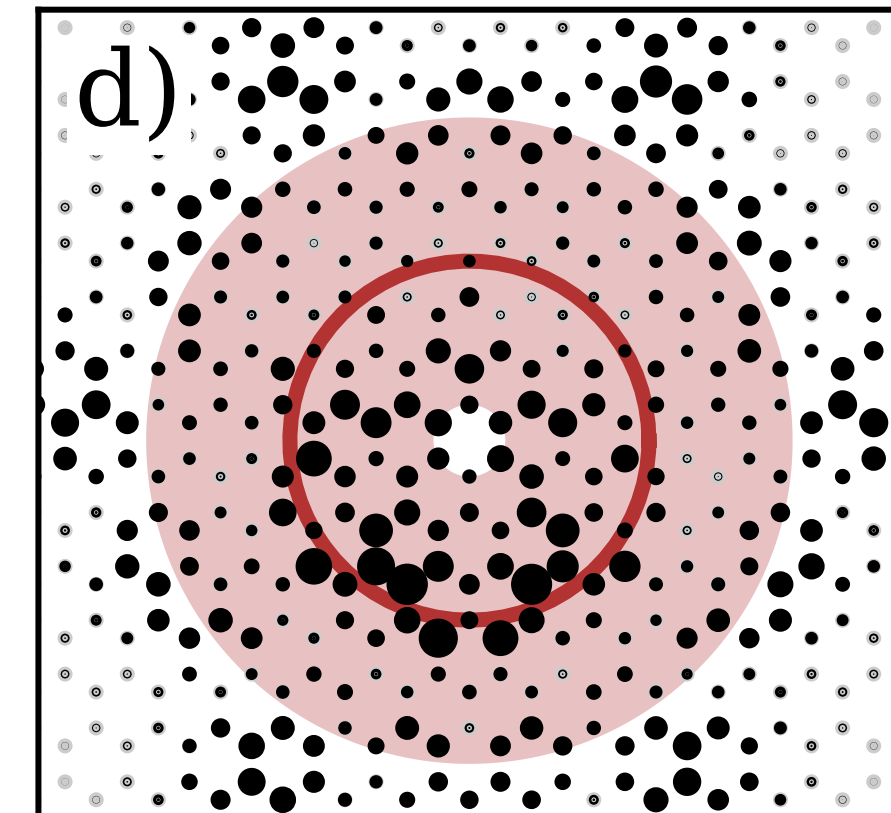
- **localization** in momentum space for incommensurate angles
- level repulsion exponentially small



1D delocalization

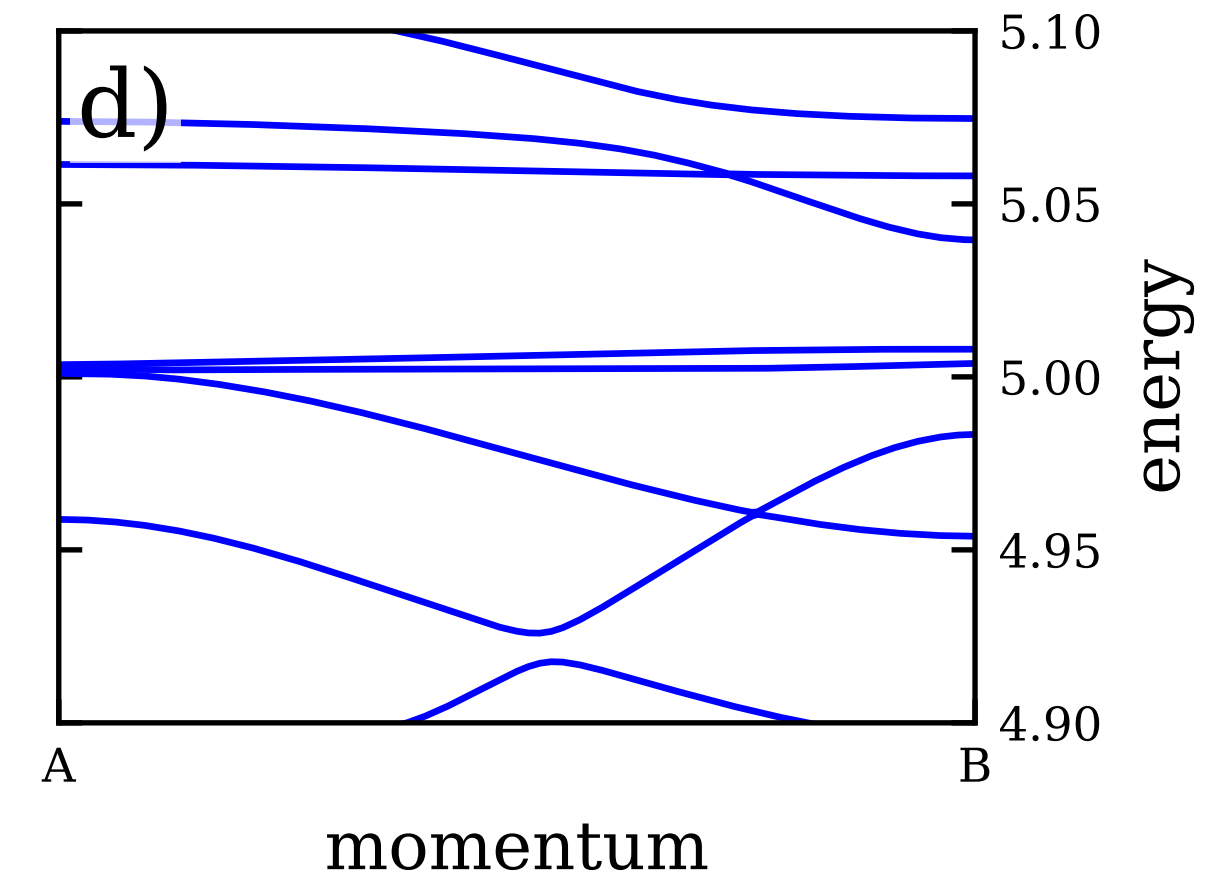


- **delocalized** along circumference
- level repulsion substantial
- regions of high velocity prevails



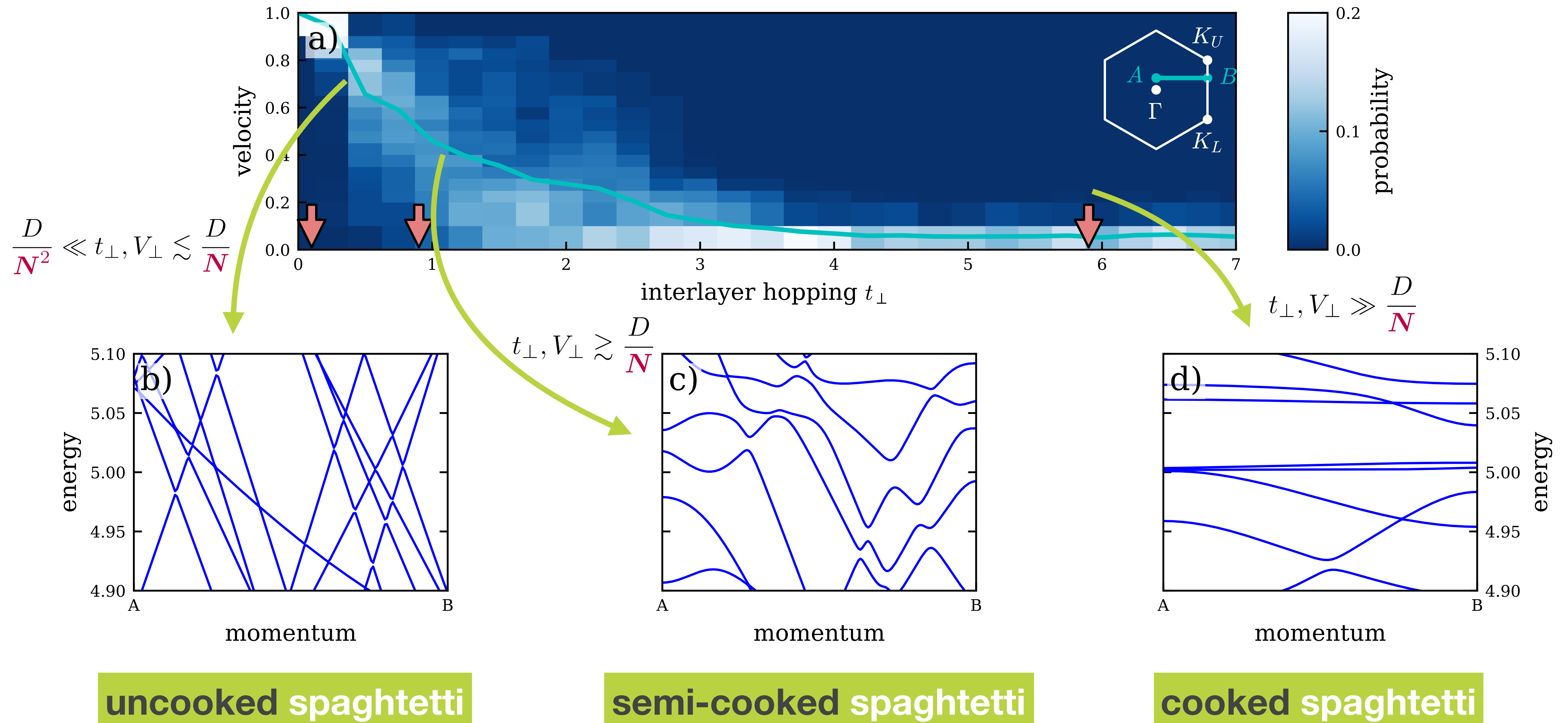
strong coupling

not ergodic!

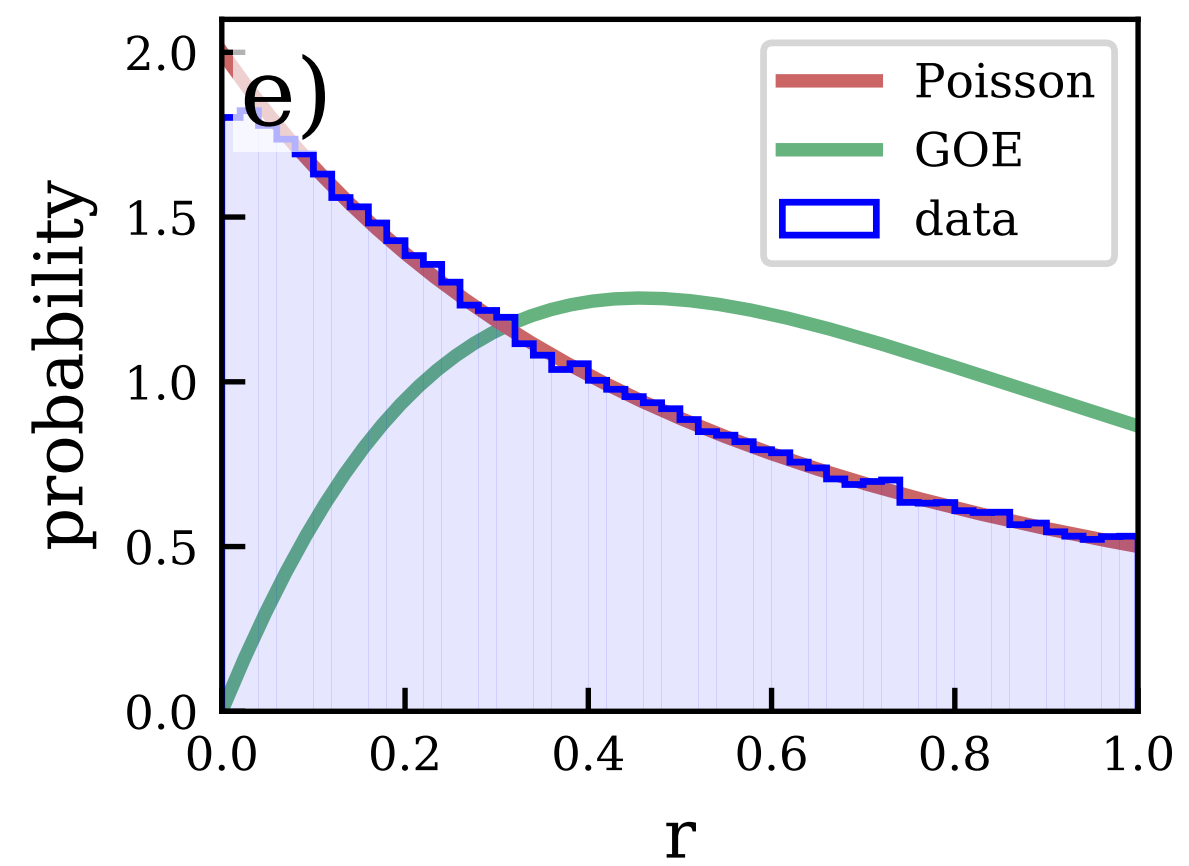
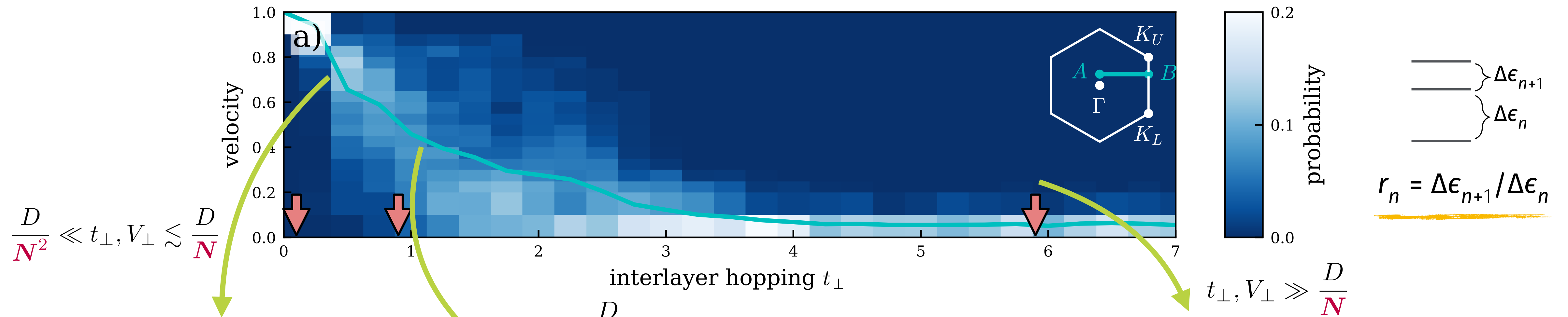


- **dimensional crossover**
- discrete symmetries become relevant
- small velocities prevail

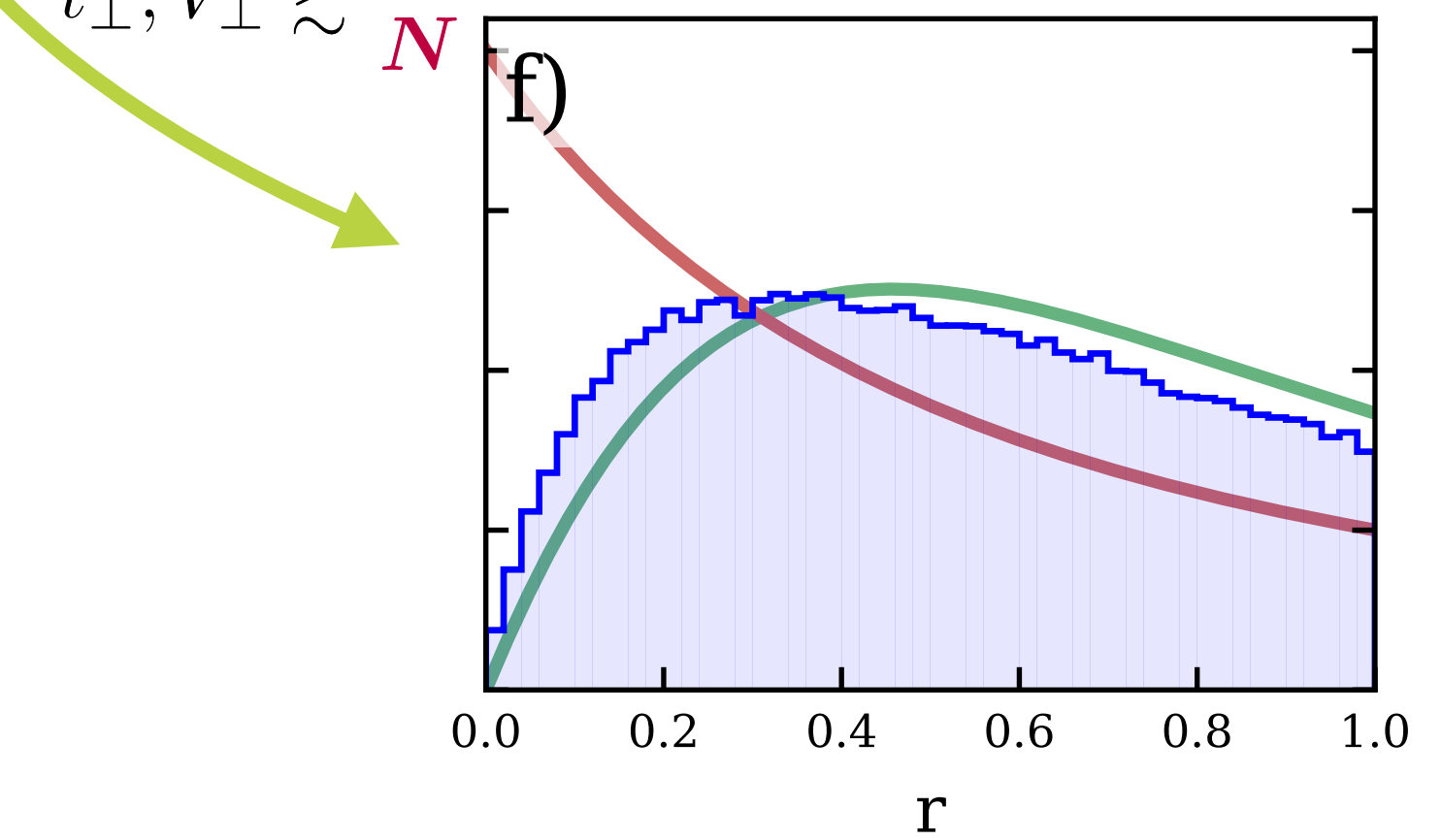
spectral statistics



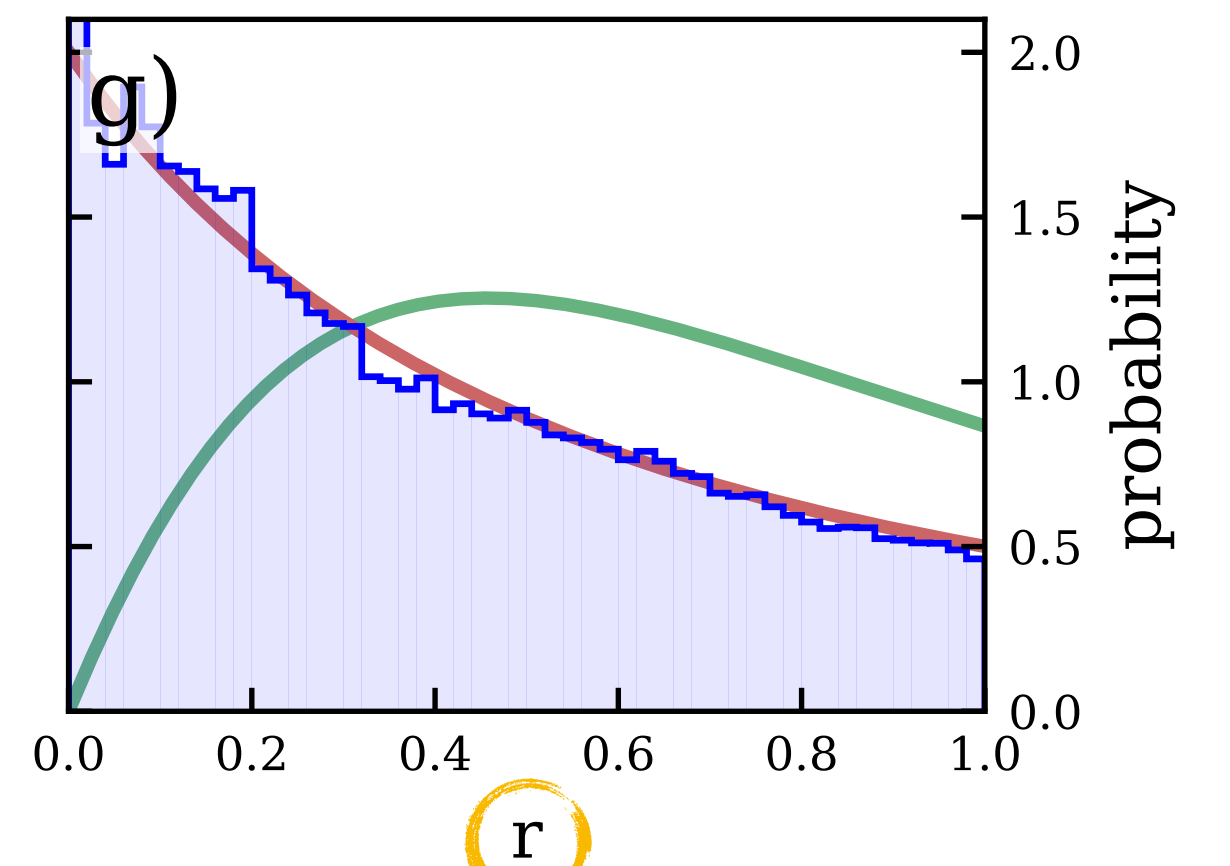
spectral statistics



uncooked spaghetti

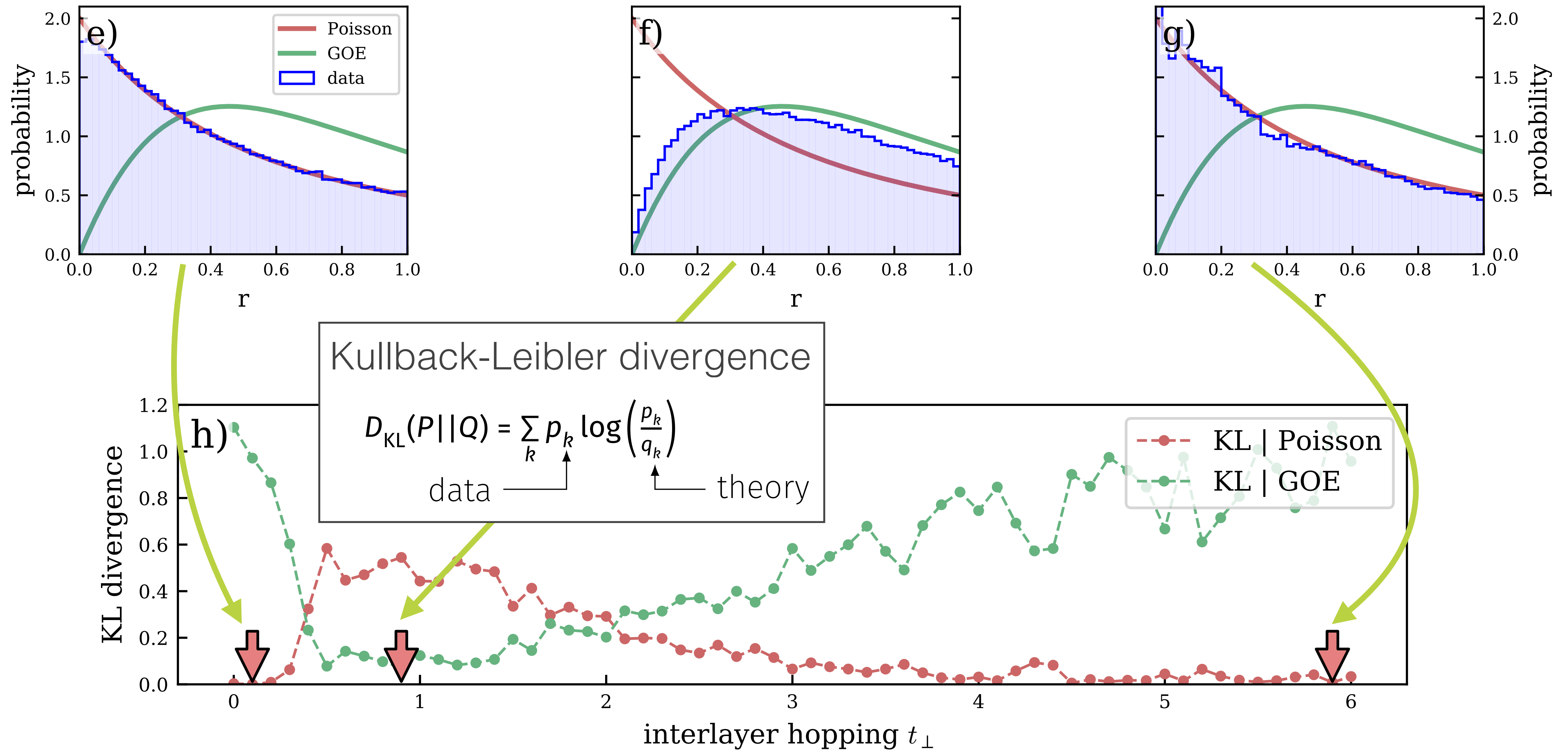



semi-cooked spaghetti



cooked spaghetti

spectral statistics





**Where to go
from here?**

summary

2D Materials **8**, 044007 (2021)



Take-away messages

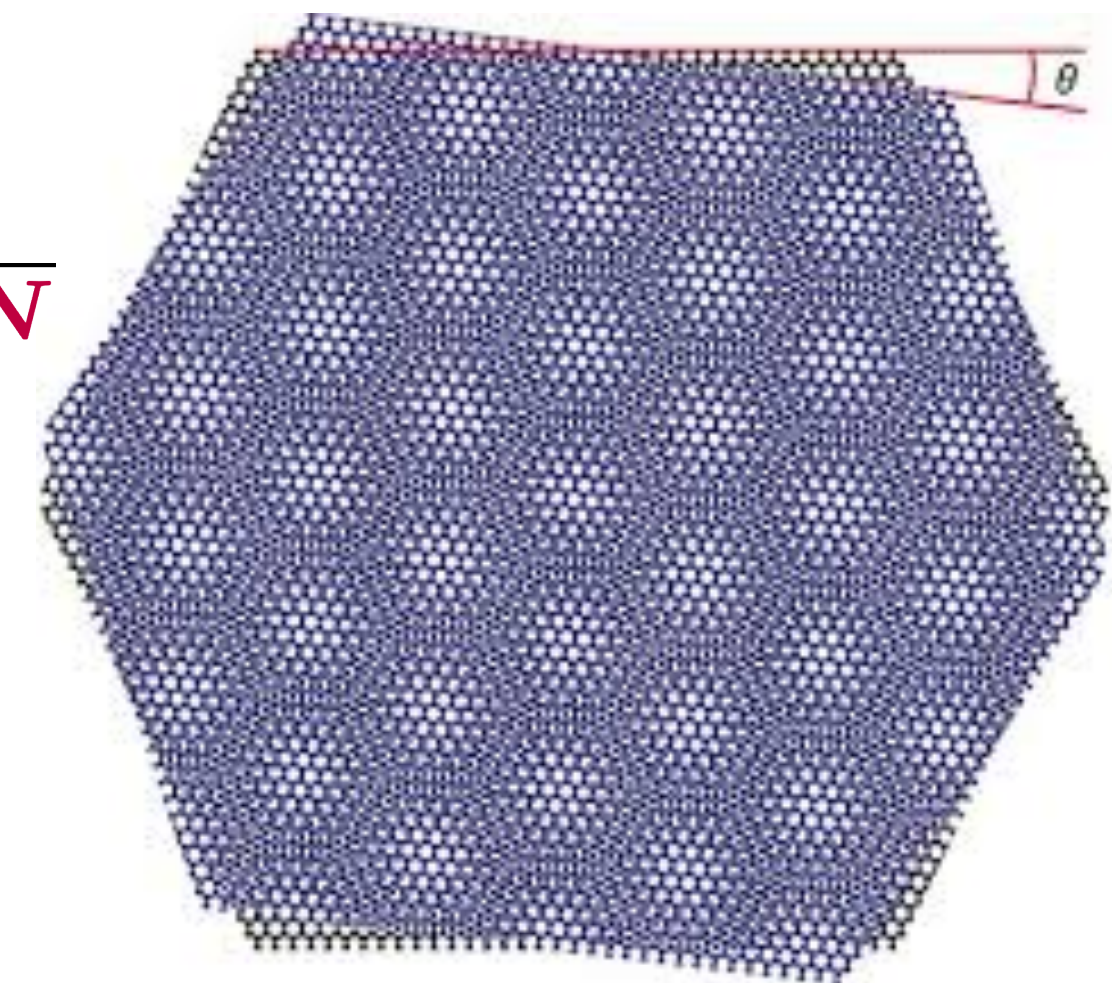
- **typical band** in generic moiré system: **not flat**, but velocities of $O(1)$
- reason – **localization in momentum space** along **1D** Fermi surface

→ **three localization regimes:**

deep localization, 1D delocalization, strong coupling

- **exceptions** to the rule:

- close to **minima/maxima** of unperturbed bands expect bandwidth $e^{-\sqrt{N}}$
- at **magic points** (derived from Dirac points of graphene)





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