

Solution 3

Tutor: Chae-Yeun Park

Nov. 5

Exercise 1

Solution to 1.1 Matrices $\{\mathbb{1}, X, Y, Z\}$ form a basis of the vectorspace of 2×2 matrices over \mathbb{C} with complex coefficients. In other word, any 2×2 matrix A can be written as $A = \sum_{\mu=0}^4 \alpha_{\mu} \sigma_{\mu}$ where $\sigma_0 = \mathbb{1}$, σ_i are Pauli matrices for $i \in \{1, 2, 3\}$ and $\alpha_{\mu} \in \mathbb{C}$. Using the orthogonality of Pauli matrices, we can also write $\alpha_{\mu} = \text{Tr}[\sigma_{\mu} A]/2$.

Using this result, we also can decompose error E_k as $E_k = \sum_{\mu} e_{\mu}^k \sigma_{\mu}$. The constraint of this channel is $\sum_k E_k^{\dagger} E_k = \mathbb{1}$ which is the trace preserving condition. Calculating this, we obtain

$$\sum_k E_k^{\dagger} E_k = (\overline{e_0^k} \mathbb{1} + \sum_i \overline{e_i^k} \sigma_i) (e_0^k \mathbb{1} + \sum_j e_j^k \sigma_j) \quad (1)$$

$$= \overline{e_0^k} e_0^k \mathbb{1} + \sum_{i,j} \overline{e_i^k} e_j^k \sigma_i \sigma_j + \sum_i \overline{e_i^k} e_0^k \sigma_i + \sum_j \overline{e_0^k} e_j^k \sigma_j \quad (2)$$

$$= \sum_{\mu} \overline{e_{\mu}^k} e_{\mu}^k \mathbb{1} + \left[\sum_{i,j} i \overline{e_i^k} e_j^k \epsilon_{ijl} + \overline{e_l^k} e_0^k + \overline{e_0^k} e_l^k \right] \sigma_l, \quad (3)$$

where we used $\sigma_i \sigma_j = \delta_{i,j} + i \epsilon_{ijk} \sigma_k$ and \bar{a} is a complex conjugate of a . Thus we require $\overline{e_{\mu}^k} e_{\mu}^k = 1$ and $\sum_{i,j} i \overline{e_i^k} e_j^k \epsilon_{ijl} + \overline{e_l^k} e_0^k + \overline{e_0^k} e_l^k = 0$.

Solution to 1.2 Recall that logical qubit of the Shor code is given as $|0_L\rangle = |+++ \rangle$ and $|1_L\rangle = |-- - \rangle$ where $|\pm\rangle = (|000\rangle \pm |111\rangle)/\sqrt{2}$. Knill-Laflamme theorem states that $\{E_k\}$ are correctable if and only if $PE_i^{\dagger} E_j P = \alpha_{i,j} P$ for some $\alpha_{i,j}$ that is a Hermitian matrix (i.e., $\alpha_{i,j} = \alpha_{j,i}^*$). Here, P is a projector on the logical qubit subspace. In our case, $P = |0_L\rangle \langle 0_L| + |1_L\rangle \langle 1_L|$. Putting this P , we obtain

$$PE_i^{\dagger} E_j P = |0_L\rangle \langle 0_L| \langle 0_L| E_i^{\dagger} E_j |0_L\rangle + |0_L\rangle \langle 1_L| \langle 0_L| E_i^{\dagger} E_j |1_L\rangle \\ + |1_L\rangle \langle 0_L| \langle 1_L| E_i^{\dagger} E_j |0_L\rangle + |1_L\rangle \langle 1_L| \langle 1_L| E_i^{\dagger} E_j |1_L\rangle. \quad (4)$$

For single qubit errors $\{E_k\}$, it is easy to show that $\langle 0_L| E_i^{\dagger} E_j |1_L\rangle = \langle 1_L| E_i^{\dagger} E_j |0_L\rangle = 0$. Likewise, we also obtain $\langle 0_L| X_a |0_L\rangle = \langle 0_L| Y_a |0_L\rangle = \langle 0_L| Z_a |0_L\rangle = 0$ and $\langle 1_L| X_a |1_L\rangle = \langle 1_L| Y_a |1_L\rangle = \langle 1_L| Z_a |1_L\rangle = 0$ where a is the location of the error. As $E_i^{\dagger} E_j = (\overline{e_0^i} e_0^j + \overline{e_1^i} e_1^j + \overline{e_2^i} e_2^j + \overline{e_3^i} e_3^j) \mathbb{1} + (\dots) X + (\dots) Y + (\dots) Z$, $\langle 0_L| E_i^{\dagger} E_j |0_L\rangle = \langle 1_L| E_i^{\dagger} E_j |1_L\rangle = \overline{e_0^i} e_0^j + \overline{e_1^i} e_1^j + \overline{e_2^i} e_2^j + \overline{e_3^i} e_3^j = \sum_{\mu=0}^3 \overline{e_{\mu}^i} e_{\mu}^j$. To sum up, we obtain $\alpha_{i,j} = \sum_{\mu=0}^3 \overline{e_{\mu}^i} e_{\mu}^j$ so $\alpha_{j,i} = \sum_{\mu=0}^3 \overline{e_{\mu}^j} e_{\mu}^i = \alpha_{i,j}^*$.

Exercise 2

Solution to 2.1 Depolarizing channel is given as $\mathcal{E}(\rho) = (1-p)\rho + p/3(X\rho X + Y\rho Y + Z\rho Z)$. Thus $\mathcal{E}^{\otimes 9}(\rho) = \sum_k (1-p)^{9-k} (p/3)^k \sum_{a=1}^{3^k} E_a^k \rho (E_a^k)^{\dagger}$ where $\{E_a^k\}$ are all possible configuration of errors with k

Pauli operators. For example, $E_1^0 = \mathbb{1}^{\otimes 9}$ and E_a^1 for $a \in [1, \dots, 27]$ are all possible single qubit errors. As there are three different errors on each qubit, we have 27 configurations. In addition, we already know that the Shor_9 code can correct any arbitrary single qubit error. So these 27 configurations are all correctable. Then let us consider $k = 2$ case. In this case, there are some correctable and uncorrectable errors. For example, let us consider there are two X errors in the same cluster of qubits. Then the error correcting procedure for $X_1 X_2 |\pm\rangle$ makes a logical error. To be precise, we measure two syndromes $Z_1 Z_2$ and $Z_2 Z_3$. The first gives 1 and the second gives -1 . Then the majority vote suggests correcting it by applying X_3 . Then the resulting state is $X_1 X_2 X_3 |\pm\rangle = \pm |\pm\rangle$ that means logical phase flip error in the encoded logical qubit. On the other hand, when there are two X errors in two different clusters then it is correctable. Likewise, we also see that two Z errors in the same cluster does not change the state but in different clusters makes a logical error. From this working principle, we can find the correctable error configurations as follows:

XX	XY	XZ	YY	YZ	ZZ
27	54	72	0	18	9

Here, Y error is just considered as XZ . Thus among total $\binom{9}{2}3^2 = 324$ error configurations, 180 error configurations are correctable. Likewise, we can extend this to larger k . As it is complicated to calculate all correctable configurations by hand, we instead calculate it by making a simple program. Using that, we obtain the number of correctable configurations as $\{1, 27, 180, 804, 2502, 6858, 14580, 20916, 15633, 4035\}$ from $k = 0$ to 9. It is also interesting as there are correctable error configurations even for $k = 9$. This is because two Z errors in the same cluster and X errors on two whole clusters do not change the state. Example correctable errors in this case are $Z_1 Z_2 X_3 Z_4 Z_5 X_6 Z_7 Z_8 X_9$ and $X_1 X_2 X_3 X_5 X_6 X_7 Z_1 Z_2 Z_3$.