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 Problems in Advanced Statistical Physics
 

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**Problem 7: Equal probability of microstates and Bernoulli measure**

For the one-dimensional asymmetric exclusion process with  $N$  particles on a ring of  $L$  sites it has been shown that all  $\binom{L}{N}$  microstates are equally likely in the nonequilibrium stationary state. Prove that this implies *Bernoulli measure* in the limit  $L, N \rightarrow \infty$  at fixed density  $\rho = N/L$ , which means that in the infinite system each site is independently occupied or vacant with probability  $\rho$  and  $1 - \rho$ , respectively. To this end, compute correlation functions  $\langle \eta_i \eta_j \rangle$ ,  $\langle \eta_i \eta_j \eta_k \rangle$  etc. for the finite system, and take the limit  $L, N \rightarrow \infty$ . Also compute the leading finite size correction to the stationary particle current.

**Problem 8: Single file diffusion**

Here we want to establish the anomalous diffusion law<sup>1</sup>  $\langle X_0(t)^2 \rangle \sim t^{1/2}$  for a tagged particle in the one-dimensional Langmuir lattice gas, where  $X_0(t)$  denotes the position of a particle that resides at the origin at time  $t = 0$ . We work in the continuum setting. The key relation is

$$X_0(t) \approx \frac{1}{\bar{\rho}} \int_0^t ds j(0, s) \quad (1)$$

for long times, where  $\bar{\rho}$  is the mean density in the system and  $j(0, t)$  is the particle current through the origin at time  $t$ .

We decompose the density  $\rho(x, t) = \bar{\rho} + \phi(x, t)$  into the mean density and a fluctuating part  $\phi(x, t)$ , which satisfies the diffusion equation

$$\frac{\partial \phi}{\partial t} = D \frac{\partial^2 \phi}{\partial x^2}. \quad (2)$$

Use this to derive the time evolution of the intermediate structure function

$$S(k, t) = \langle \hat{\phi}(k, t) \hat{\phi}(-k, 0) \rangle \quad (3)$$

where brackets refer to the thermal equilibrium state and

$$\hat{\phi}(k, t) = \int dx e^{ikx} \phi(x, t). \quad (4)$$

You should find

$$S(k, t) = e^{-|k|^2 D |t|} S(k, 0) = L \bar{\rho} (1 - \bar{\rho}) e^{-|k|^2 D |t|}. \quad (5)$$

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<sup>1</sup>S. Alexander, P. Pincus, Phys. Rev. B **18**, 2011 (1978)

In the continuum limit the equal time density correlation function is  $\langle \phi(x, t)\phi(x', t) \rangle = \bar{\rho}(1 - \bar{\rho})\delta(x - x')$ , and hence  $S(k, 0) = L\bar{\rho}(1 - \bar{\rho})$ , where  $L$  is the system size.

To compute the mean square displacement  $\langle X_0^2 \rangle$ , use (1) to relate it to the correlation function of the current  $j(0, t)$ . The latter can then be obtained from the density correlation function (5) using the fact that  $j = -D\partial\phi/\partial x$ . The calculation is simplified by introducing the *particle counting function*

$$N(x, t) = \int_0^x dy \phi(y, t), \quad (6)$$

which satisfies the relations  $j(0, t) = \partial N/\partial t$  and  $\phi(x, t) = \partial N/\partial x$ . Finally, examine how  $\langle X_0(t)^2 \rangle$  behaves in the low density limit  $\bar{\rho} \rightarrow 0$ , and provide an interpretation.

### Problem 9: The zero range process

In the zero range process (ZRP)<sup>2</sup> an unlimited number  $n_i = 0, 1, 2, \dots$  of particles can occupy each site  $i = 1, \dots, L$  of the one-dimensional lattice with periodic boundary conditions (a ring). A particle at site  $i$  jumps to the right ( $i \rightarrow i + 1$ ) with probability  $p$  and to the left ( $i \rightarrow i - 1$ ) with probability  $1 - p$  at a rate which is a function  $\gamma(n_i)$  of the number of particles at the site of origin with  $\gamma(0) = 0$ . There is no dependence on the occupancy of the target site (= *zero range interaction*). The ZRP has the remarkable property that the stationary distribution is a *product measure* for a broad class of functions  $\gamma(n)$ , i.e. the stationary weight of a configuration  $\{n_1, \dots, n_L\}$  is of the form

$$\text{Prob}[n_1, \dots, n_L] \sim \prod_{i=1}^L f(n_i). \quad (7)$$

- a.) In the symmetric case  $p = 1/2$  use the condition of detailed balance to show that

$$f(n) \sim \prod_{k=1}^n \gamma(k)^{-1} \quad (8)$$

You may ignore the constraint of constant total particle number  $N = \sum_{i=1}^L n_i$ , i.e. work in a “grand-canonical” setting.

- b.) Write down the master equation for the asymmetric case and show that the product measure (7, 8) remains stationary also when  $p \neq 1/2$ .
- c.) Under what conditions on the rate function  $\gamma(n)$  is the single-site probability distribution (8) normalizable?

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<sup>2</sup>F. Spitzer, Adv. Math. **5**, 246 (1970); M.R. Evans, T. Hanney, J. Phys. A **38**, R195 (2005).