## Quantum Field Theory III Exercise sheet 7

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## 7. Equation of motion for the WZW action

Consider the non-linear  $\sigma$ -model action in d = 2 dimensions

$$S_{\text{WZW}}[g] = S_2[g] + \Gamma[g] = \frac{1}{8\pi} \int_{S^2} d^2 x \operatorname{tr}(\partial_\mu g \partial_\mu g^{-1}) + \Gamma[g], \qquad (1)$$

where the field  $g \in U(N)$ , the two-dimensional space is compactified to the 2-sphere  $S^2$ , and  $\Gamma[g]$  is the Wess-Zumino-Witten term

$$\Gamma[g] = -\frac{i}{12\pi} \int_{B^3} d^3x \,\epsilon^{\mu\nu\rho} \mathrm{tr}(g^{-1}\partial_\mu g g^{-1}\partial_\nu g g^{-1}\partial_\nu g g^{-1}\partial_\nu g). \tag{2}$$

It is assumed here that the sphere  $S^2$  serves as the boundary of one half of a three-dimensional unit ball  $B^3$  (i.e.  $\partial B^3 \equiv S^2$ ) and the field  $g(x_0, x_1)$  initially defined in d = 2 is smoothly extended in the third  $(x_2)$ -direction.

- a) Check first that without the WZW term, the saddle point equation corresponding to the  $\sigma$ -model action  $S_2[g]$  has the form  $\partial_{\mu}(g^{-1}\partial_{\mu}g) = 0$ .
- b) Derive further the saddle point equation of the full action S[g] with the WZW term included. You have to arrive to the result  $\partial_z \bar{j} = 0$ , where the current  $\bar{j} \propto (\partial_{\bar{z}}g)g^{-1}$  and the complex coordinates are defined as  $z = (x_0 + ix_1)/\sqrt{2}$  and  $\bar{z} = (x_0 ix_1)/\sqrt{2}$ .

## 8. The $\sigma$ -model action of the disordered BCS superconductor

The disordered s-wave BCS superconductor (in dimension d = 2, 3) can be described by the following low-energy action

$$S[Q] = \frac{1}{2}\pi\nu \int d^d \mathbf{r} \operatorname{tr} \left[ \frac{1}{4} D(\nabla_{\mathbf{r}} Q)^2 - \epsilon_n \tau_3 Q - \Delta \tau_2 Q \right], \qquad (3)$$

where the field Q is the Matsubara Green's function at energy  $\epsilon_n$  (evaluating at coinciding spatial points) obeying the constrain  $Q^2 = 1$ , Pauli matrices  $\tau_i$  act in the Nambu space,  $\Delta$  is the BCS gap, D is the diffusion coefficient, and  $\nu$  is the density of states (DoS) in the normal state at Fermi energy.

a) By minimazing the action S[Q] derive the saddle point equation for Q,

$$D\nabla_{\mathbf{r}}(Q\nabla_{\mathbf{r}}Q) = \left[\epsilon_n \tau_3 + \Delta \tau_2, Q\right],\tag{4}$$

which is known as the Usadel equation.

**b)** Using the above result, find the stationary (position-independent) Green's function  $Q(\epsilon_n)$ . Performing the analytic continuation to the real energy, deduce the DoS in the disordered superconductor,  $\nu(\epsilon) = \nu \operatorname{Re} \operatorname{tr} \tau_3 Q(\epsilon_n) \Big|_{i\epsilon_n \to \epsilon + i0^+}$ . Compare it with the density of states known from the BCS model.

## References

[1] Altland, A. and Simons, B. Condensed Matter Field Theory, 2nd edition, Cambridge University Press (2010)