

Quantum Field Theory II

Problem Set 8

Please hand in your solution of the exercise marked with (*), for which we offer correction,
in the mail box until Monday 12am.

1. Quantum phase transition: self-consistent one-loop approximation (*)

The goal of this exercise is to draw a phase diagram as a function of temperature T and a parameter (r_0 in this exercise) under the self-consistent Hartree-Fock approximation. To this end, we consider an action for a single-component bosonic field $\phi_{\mathbf{k}}(i\omega_n)$, written as

$$S = \sum_{\mathbf{k}, \omega_n} \phi_{\mathbf{k}}^*(i\omega_n) (r_0 + \mathbf{k}^2 + |\omega_n|^{\frac{2}{z}}) \phi_{\mathbf{k}}(i\omega_n) + \frac{g}{4\beta\mathcal{V}} \sum_{\omega_{n_1}, \omega_{n_2}, \omega_{n_3}} \sum_{\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3} \phi_{\mathbf{k}_1}^*(i\omega_{n_1}) \phi_{\mathbf{k}_2}^*(i\omega_{n_2}) \phi_{\mathbf{k}_3}(i\omega_{n_3}) \phi_{\mathbf{k}_1 + \mathbf{k}_2 - \mathbf{k}_3}(i\omega_{n_1} + i\omega_{n_2} - i\omega_{n_3}). \quad (1)$$

Here z is called "dynamical critical exponent" with, e.g., $z = 1$ for an antiferromagnetic transition in an insulator while $z = 2$ in a metal due to the friction from electrons (Landau-damping). In small g limit, the first-order perturbation is taken. Taking into account the Hartree-Fock diagram, the full Green's function for the field is given by

$$G_{\mathbf{k}}(i\omega_n) \equiv \langle \phi^*(\mathbf{k}, i\omega_n) \phi(\mathbf{k}, i\omega_n) \rangle = \frac{1}{r_0 + \mathbf{k}^2 + |\omega_n|^{\frac{2}{z}} - \Sigma} = \frac{1}{r + \mathbf{k}^2 + |\omega_n|^{\frac{2}{z}}}, \quad (2)$$

where the Hartree-Fock contribution to the self energy is just a number independent of ω_n and \mathbf{k} given by

$$\Sigma = -\frac{g}{\beta\mathcal{V}} \sum_{i\omega_n, \mathbf{k}} G_{\mathbf{k}}(i\omega_n) \equiv \delta r(r, T), \quad (3)$$

and redefine a parameter $r \equiv r_0 - \delta r$. r_0^c is defined as $r_0^c \equiv \delta r(r = 0, T = 0)$. With this definition, r vanishes for $r_0 = r_0^c$ at $T = 0$.

- a) Let us consider the zero temperature case. How does r depend on $r_0 - r_0^c$? At zero temperature, the summation $\frac{1}{\beta} \sum_{i\omega_n}$ becomes the integral form of $\int (d\omega/2\pi)$. You may use the identity

$$\int \frac{d^d k}{(2\pi)^d} \int \frac{d\omega}{2\pi} \frac{1}{r + k^2 + |\omega|^{2/z}} = D \int d^{d+z} k \frac{1}{r + k^2}. \quad (4)$$

D is a numerical constant. This identity implies that a quantum theory in d dimension is related to a classical theory in $d + z$ dimension. The integral of the right hand side of Eq. (4) is approximately computed by the division of the integral

$$\int_0^{k_{\max}} dk = \int_0^{\sqrt{r}} dk + \int_{\sqrt{r}}^{k_{\max}} dk, \quad (5)$$

and setting the denominator

$$r + k^2 \simeq \begin{cases} r, & r > k^2 \\ k^2, & r < k^2 \end{cases}. \quad (6)$$

Show that it leads to the self-consistent equation

$$r = r_0 - r_0^c + Cr^{\frac{d+z-2}{2}} \quad \text{for } d + z > 2. \quad (7)$$

C is a numerical constant and you do not need to compute C . The condition $d + z > 2$ is needed to avoid the singularity at $r = 0$.

- b) Now we are interested in the static susceptibility $\chi(\mathbf{k} = 0, \omega = 0) = 1/r$. Show that depending on the condition of the exponent $d + z$, $\chi(\mathbf{k} = 0, \omega = 0)$ in small r limit has different scalings in $1/(r_0^c - r_0)$. What is the scaling of $\chi(\mathbf{k} = 0, \omega = 0)$ for two different cases $d + z > 4$ and $2 < d + z < 4$?
- c) The correlation length ξ is determined by $\xi = 1/\sqrt{r}$ from the right hand side of Eq. (2). Find the scaling exponent ν of ξ in $1/(r_0^c - r_0)$ for the different cases $d + z > 4$ and $2 < d + z < 4$.
- d) Properties at $T > 0$ can be obtained from the T dependence of δr . But it is tedious to find the T dependence of δr using the Matsubara summation. Instead, we use a dimensional analysis to know the T dependence of δr . For the case of $d + z > 4$, estimate T_c assuming this T dependence is valid in $T \simeq T_c$. What is the scaling exponent ψ in the formula $T_c \sim (r_0^c - r_0)^\psi$? A schematic phase diagram (T_c as a function of r_0) is drawn in Fig. 1.

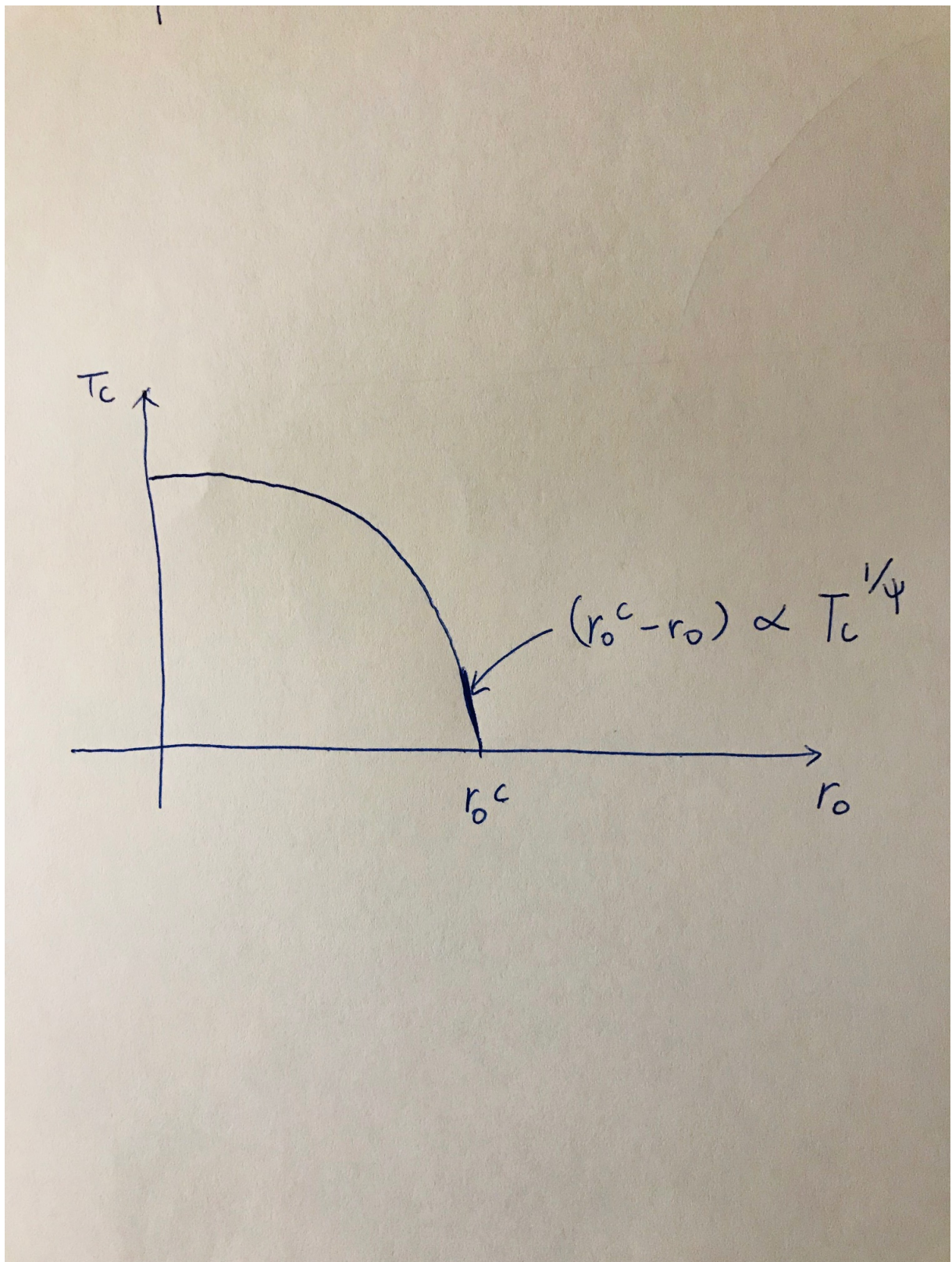


Figure 1: A schematic phase diagram of T_c as a function of r_0 .