

Quantum operation, Kraus-operator expansion

time-evolution of an open quantum system

over a time interval $[0, t]$ defines a

quantum operation

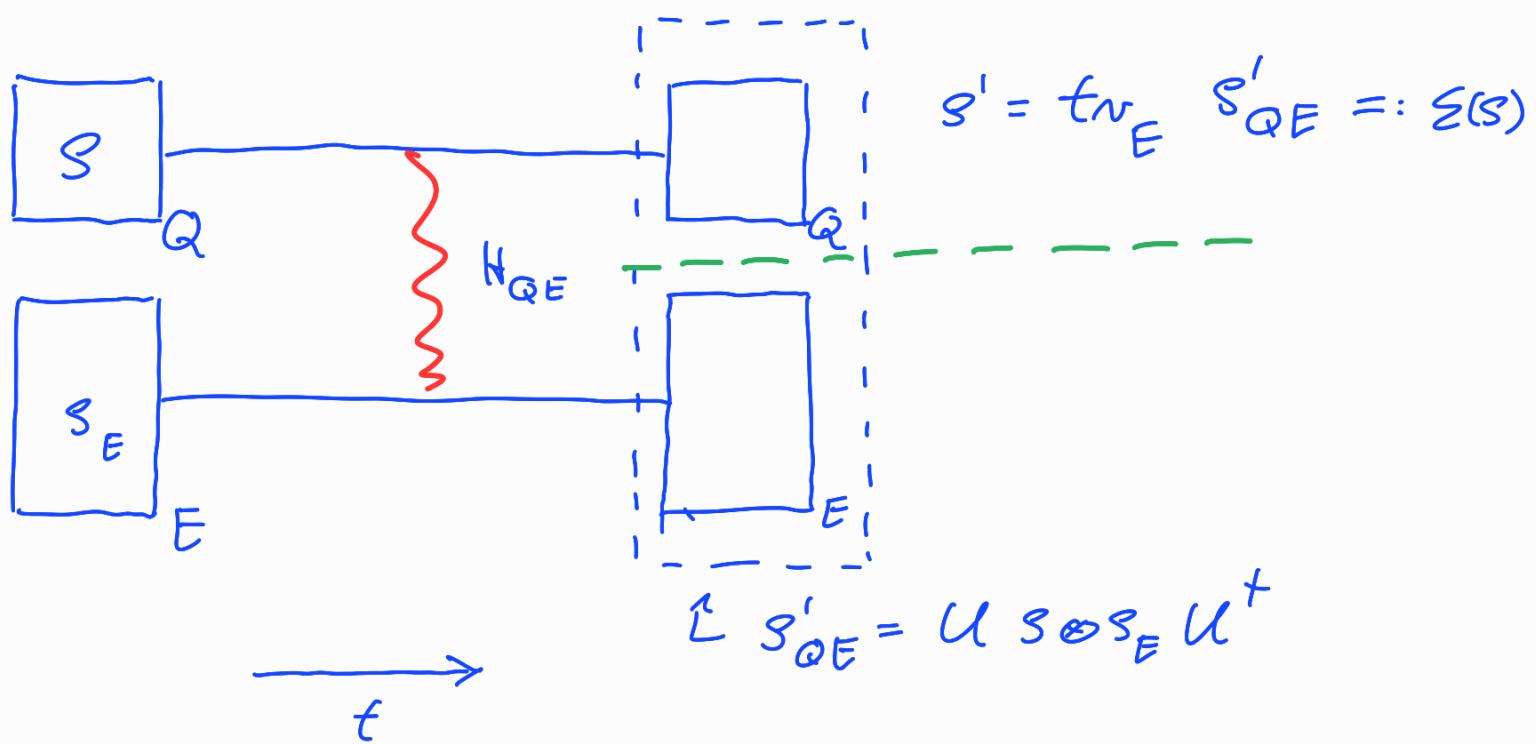
$$\Sigma: S \mapsto \Sigma(S)$$

with

$$\Sigma(S) = \text{tr}_E U S \otimes S_E U^+$$

- S_E : initial state of environment
- $U = e^{-iH_{QE}t/\hbar}$: time evolution

of joint system QE



quantum operations can always be
casted in a relatively simple

Kraus-operator expression:

Thm.:

$\Sigma(S)$ quantum operation

\Leftrightarrow

\exists Kraus-ops E_1, \dots, E_k cu \mathcal{H}_Q s.t.:

$$\Sigma(S) = \sum_{k=1}^K E_k S E_k^+$$

- $\sum_k E_k^+ E_k = \mathbb{I}_Q$
- $k \leq (\dim \mathcal{H}_Q)^2$

Examples:

- unitary time evolution of closed system:

$$S \mapsto S_t = U_t S U_t^+ , \quad U^\dagger U = \mathbb{I}$$

$\hat{=}$ Kraus-op. expans. with $K=1$

- projective measurement:

$$S \mapsto S' = \sum_{i=1}^k P_i S P_i$$

P_1, \dots, P_n orthog. projections, $\sum_{i=1}^n P_i = I$

- bif-fid: $S \mapsto \Sigma(S) = (1-p)S + p \Gamma_X S \Gamma_X^*$
 $= E_1 S E_1^* + E_2 S E_2^*$

with $E_1 = \sqrt{1-p} I$, $E_2 = \sqrt{p} \Gamma_X$.

Proof:

" \Rightarrow ": w.l.o.g. S_E pure state, say

$$S_E = |\varphi_1\rangle\langle\varphi_1| \quad (\text{a mixed state}$$

S_E can always be "purified": $S_E = \text{tr}_E |\varphi_1'\rangle\langle\varphi_1'|$

let $|\varphi_1\rangle, \dots, |\varphi_K\rangle$ be ONB of \mathcal{H}_E

$$\begin{aligned} \rightarrow \Sigma(S) &= \text{tr}_E U (S \otimes |\varphi_1\rangle\langle\varphi_1|) U^* \\ &= \sum_{l=1}^K \sum_{k=1}^K \underbrace{\langle\varphi_k| U |\varphi_l\rangle\langle\varphi_k|}_{\cong S_{lk}} \underbrace{S \otimes |\varphi_1\rangle\langle\varphi_1| U^*}_{\cong S_{l1}} |\varphi_l\rangle \\ &= \sum_{l=1}^K \langle\varphi_k| U |\varphi_l\rangle S \langle\varphi_1| U^* |\varphi_l\rangle \end{aligned}$$

\rightarrow with $E_h := \langle \varphi_h | u | \varphi_i \rangle :$

$$\Sigma(S) = \sum_{h=1}^k E_h \leq E_h^+ .$$

What is meant by $\langle \varphi_h | u | \varphi_i \rangle$?

with $|1\rangle, \dots, |d\rangle$ ONB of \mathcal{H}_Q

$|\varphi_1\rangle, \dots, |\varphi_k\rangle$ ONB of \mathcal{H}_S

$$u = \sum_{ij=1}^d \sum_{l,m=1}^k c_{ij,lm} |i\rangle\langle jl| \otimes |\underline{\varphi_l}\rangle\langle \underline{\varphi_m}|$$

$$\rightarrow \underbrace{\langle \varphi_h |}_{=} \underbrace{u | \varphi_i \rangle}_{=} = \sum_{ij=1}^d c_{ij,h1} |i\rangle\langle jl|$$

operator on \mathcal{H}_Q !

]

• HS: $1 = \text{tr } \Sigma(S) = \text{tr} \left(\sum_h E_h S E_h^+ \right) = \text{tr} \left(\sum_h E_h^+ E_h \right) S$

$$\rightarrow \sum_h E_h^+ E_h = \mathbb{1}_Q ;$$

• $n \leq (\dim \mathcal{H}_Q)^2$: cf. exercise 19.

" \Leftarrow " given E_1, \dots, E_k , $\sum_s E_s^* E_s = \mathbb{1}$, find

\mathcal{H}_E , φ_E , and unitary U on $\mathcal{X}_Q \otimes \mathcal{H}_E$

s.t.

$$\sum_{s=1}^k E_s s E_s^* = U \mathcal{H}_E U^* \otimes \varphi_E U^*$$

for all s

choose $\dim \mathcal{H}_E = k$ with $|\varphi_1\rangle, \dots, |\varphi_k\rangle$

ONB of \mathcal{H}_E , $\varphi_E := |\varphi_1\rangle\langle\varphi_1|$, and

unitary U s.t.

$$U |i\rangle |\varphi_1\rangle = \underbrace{\sum_{s=1}^k (E_s |i\rangle) \otimes |\varphi_s\rangle}_{\cong |\psi_i\rangle}$$

$$\text{Then: } \langle \varphi_i | \varphi_j \rangle = \sum_{s, s' = 1}^k \langle i | E_s^* E_{s'} | j \rangle \langle \varphi_s | \varphi_{s'} \rangle$$

$$= \langle i | \underbrace{\sum_{s=1}^k E_s^* E_s}_{= \mathbb{1}} | j \rangle = \langle i | j \rangle \quad \checkmark$$

then $\text{tr}_E U S \otimes |\psi_i\rangle\langle\psi_i| U^+$

$$= \sum_{i,j} \sum_h \underbrace{\langle\psi_g|U_E|c_i\rangle\langle c_i|S|j\rangle}_{E_h} \underbrace{\langle j|\langle\psi_i|U^+|\psi_h\rangle}_{\geq E_h|c_i\rangle} \\ \geq E_h|c_i\rangle \geq \langle c_i|E_g^+$$

$$= \sum_h E_h S E_h^+.$$

Unitary freedom in the choice of basis operators :

for $S_E = |\psi_1\rangle\langle\psi_1|$ and ONB $|\psi_1\rangle, \dots, |\psi_n\rangle$

of \mathcal{H}_E we have

$$\bar{E}_g := \langle\psi_g|U|\psi_i\rangle$$

for another ONB $|X_1\rangle, \dots, |X_n\rangle$ of \mathcal{H}_E
we find alternative operators

$$\bar{F}_e = \langle X_e|U|\psi_i\rangle$$

and certainly

$$\sum_e F_e S F_e^+ = \varepsilon(S) = \sum_h E_h S E_h^+$$

$\{F_e\}$ and $\{E_h\}$ are related as follows:

$$F_e = \sum_h \underbrace{\langle \chi_e | \varphi_h \rangle}_{\cong u_{eh}} \underbrace{\langle \varphi_h | u | \varphi_i \rangle}_{\cong E_h} = \sum_h u_{eh} E_h$$

→ any unitary matrix $(u)_{eh}$ transforms

$\{E_h\}$ to an equivalent set $\{F_e\}$

with

$$F_e = \sum_h u_{eh} E_h .$$

