Exercise 3

We define a family of 2-qubit states $|\psi(\alpha)\rangle \in \mathcal{H}_A \otimes \mathcal{H}_B$ for $\alpha \in \mathbb{R}$ with $\frac{1}{2} \leq \alpha \leq 1$ as

$$|\psi(\alpha)\rangle = \left(\frac{\alpha}{2}\right)^{\frac{1}{2}}|00\rangle + \left(\frac{\alpha}{2}\right)^{\frac{1}{2}}|11\rangle + \left(\frac{1-\alpha}{2}\right)^{\frac{1}{2}}|01\rangle + \left(\frac{1-\alpha}{2}\right)^{\frac{1}{2}}|10\rangle \qquad (1)$$

a) You measure the first qubit. What is the probability of finding the measurement outcome that corresponds to the $|0\rangle$ state for the first qubit?

b) If you indeed find this measurement outcome, what is the post-measurement state? c) Show that $|\psi(\alpha)\rangle$ is entangled for all values of α , except $\alpha = \frac{1}{2}$. You are allowed to assume that all amplitudes are real.

Hint 1: Assume that $|\psi(\alpha)\rangle$ is a product state and work towards a contradiction for values other than $\alpha = \frac{1}{2}$.

Hint 2: Another option is to look at the Schmidt rank of the reduced density matrix $\rho_A = \text{Tr}_B(|\psi(\alpha)\rangle\langle\psi(\alpha)|).$

Hint 3: A third way is by looking at $Tr(\rho_A^2)$. If this is your preferred way, also explain why this is sufficient.

Solution

a) Let l be the measurement outcome corresponding to the $|0\rangle$ state of the first qubit.

$$P(l) = \langle \psi(\alpha) | P_0 \otimes \mathbb{I} | \psi(\alpha) \rangle = \sum_x \langle \psi(\alpha) | |0x \rangle \langle 0x | |\psi(\alpha) \rangle$$
$$= \langle \psi(\alpha) | |00 \rangle \langle 00 | |\psi(\alpha) \rangle + \langle \psi(\alpha) | |01 \rangle \langle 01 | |\psi(\alpha) \rangle$$
$$= \frac{\alpha}{2} + \frac{1 - \alpha}{2} = \frac{1}{2}$$

b) To practice the new density matrix formalism, let's do this calculation using $\rho = |\psi(\alpha)\rangle\langle\psi(\alpha)|$ and the measurement rules from exercise 0. The post-measurement state is

$$\rho' = (p(00) + p(01) \frac{(P_{00} + P_{01}) |\psi(\alpha)\rangle\langle\psi(\alpha)| (P_{00} + P_{01})}{\operatorname{Tr}((P_{00} + P_{01}) |\psi(\alpha)\rangle\langle\psi(\alpha)|)} = \alpha |00\rangle\langle00| + (1 - \alpha) |01\rangle\langle01| + \sqrt{\alpha(1 - \alpha)}(|00\rangle\langle01| + |01\rangle\langle00|).$$

c) There are three (or possibly even more) ways of solving this exercise. **Option 1:** Assume $|\psi(\alpha)\rangle$ is a product state of 2 qubits. Then it should be possible to write

$$|\psi(\alpha)\rangle = (a |0\rangle + b |1\rangle) \otimes (c |0\rangle + d |1\rangle),$$

for $|a|^2 + |b|^2 = |c|^2 + |d|^2 = 1$. Multiplying out, we get

$$|\psi(\alpha)\rangle = ac |00\rangle + ad |01\rangle + bc |10\rangle + bd |11\rangle.$$

So we have $ac = bd = \sqrt{\frac{\alpha}{2}}$ and $ad = bc = \sqrt{\frac{1-\alpha}{2}}$. Substituting a = bd/c into ad = bc we arrive at the conclusion that $c = \pm d$. Similarly, we can derive that $a = \pm b$. Using the normalisation, the only (real) solutions are $a = \pm \frac{1}{\sqrt{2}}$ and the same for b, c, d. I.e. the only valid value for α in the given range is $\alpha = \frac{1}{2}$.

Option 2: Calculate the reduced density matrix ρ_A and determine its rank. This is the Schmidt number of $\psi(\alpha)$.

$$\begin{split} \rho_A &= \operatorname{Tr}_B(|\psi(\alpha)\rangle\!\langle\psi(\alpha)|) = \sum_i (\mathbb{I}\otimes\langle i|) |\psi(\alpha)\rangle\!\langle\psi(\alpha)| \left(\mathbb{I}\otimes|i\rangle\right) \\ &= \ldots = \frac{1}{2} |0\rangle\!\langle 0| + \frac{1}{2} |1\rangle\!\langle 1| + \frac{1}{2}\sqrt{\alpha(1-\alpha)}(|0\rangle\!\langle 1| + |1\rangle\!\langle 0|) \\ &= \frac{1}{2} \left(\frac{1}{\sqrt{\alpha(1-\alpha)}} \sqrt{\alpha(1-\alpha)}\right) \end{split}$$

This matrix is a rank 1 matrix if and only if $\alpha = \frac{1}{2}$ if α is restricted to the reals. If the Schmidt rank is larger than 1, the state is entangled. **Option 3:** Again looking at ρ_A , calculate

$$\operatorname{Tr}(\rho_A^2) = \operatorname{Tr}\left(\frac{1}{4} \begin{pmatrix} 1+4\alpha(1-\alpha) & \dots \\ \dots & 1+4\alpha(1-\alpha) \end{pmatrix} \right)$$
$$= \frac{1}{4} \cdot 2 \cdot (1+4\alpha(1-\alpha)) = \frac{1}{2} + 2\alpha(1-\alpha).$$

This equals 1 if and only if $\alpha = \frac{1}{2}$. This is a sufficient condition, because it means that ρ_A is pure and pure states can't be entangled with other systems. Conversely, every mixed state is entangled with its purification.