

Lösungshinweise Blatt 5

$$23) \quad n=0: \quad \int_0^{\infty} e^{-x} dx = 1 = 0! \quad \checkmark$$

$$n \rightarrow n+1: \quad \int_0^{\infty} x^{n+1} e^{-x} dx = -x^{n+1} e^{-x} \Big|_0^{\infty} + (n+1) \int_0^{\infty} x^n e^{-x} dx \stackrel{\text{I.V.}}{=} (n+1)!$$

$$24a) \quad \| |\psi_0\rangle \|^2 = \frac{1}{\pi a_0^3} 4\pi \int_0^{\infty} dr r^2 e^{-2r/a_0} = \frac{1}{2} \int_0^{\infty} dx x^2 e^{-x} \stackrel{23)}{=} 1$$

$r = a_0 x/2$

$$b) \quad \langle \vec{x} \rangle_{\psi_0} = \vec{0} \quad \text{und} \quad \langle \vec{p} \rangle_{\psi_0} = \vec{0} \quad \text{aufgrund Symmetrie ;}$$

$$\langle |\vec{x}|^2 \rangle_{\psi_0} = \frac{1}{\pi a_0^3} 4\pi \int_0^{\infty} dr r^4 e^{-2r/a_0} \stackrel{r=a_0 x/2}{=} \frac{a_0^2}{8} \int_0^{\infty} dx x^4 e^{-x} \stackrel{23)}{=} 3 a_0^2$$

$$\langle |\vec{p}|^2 \rangle_{\psi_0} = \frac{1}{\pi a_0^3} 4\pi \int_0^{\infty} dr r^2 e^{-r/a_0} (-\hbar^2 \Delta e^{-r/a_0}) ;$$

$$\begin{aligned} \text{mit } \Delta e^{-r/a_0} &= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} e^{-r/a_0} \right) = \frac{1}{r^2} \frac{\partial}{\partial r} \left(-\frac{r^2}{a_0} e^{-r/a_0} \right) \\ &= \frac{1}{a_0 r} \left(-2r + \frac{r^2}{a_0} \right) = -\frac{2}{a_0} + \frac{1}{a_0^2} r \end{aligned}$$

$$\begin{aligned} \text{folgt: } \langle |\vec{p}|^2 \rangle_{\psi_0} &= \frac{8\hbar^2}{a_0^2} \int_0^{\infty} dr r e^{-2r/a_0} - \frac{4\hbar^2}{a_0^5} \int_0^{\infty} dr r^3 e^{-2r/a_0} \\ &\stackrel{r=a_0 x/2}{=} \frac{2\hbar^2}{a_0^2} \underbrace{\int_0^{\infty} dx x e^{-x}}_{=1} - \frac{\hbar^2}{2a_0^2} \underbrace{\int_0^{\infty} dx x^3 e^{-x}}_{=2} = \left(\frac{\hbar}{a_0} \right)^2 \end{aligned}$$

$$25a) \quad \text{Impuls Erhaltungsgröße} \rightarrow \langle \hat{p} \rangle_{\psi(t)} = \langle \hat{p} \rangle_{\psi_0} = p_0 ;$$

$$\begin{aligned} \text{nach 13a): } \frac{d}{dt} \langle \hat{x} \rangle_{\psi(t)} &= \left\langle \frac{i}{\hbar} [H, \hat{x}] \right\rangle_{\psi(t)} = \left\langle \frac{i}{\hbar} \left[\frac{\hat{p}^2}{2m}, \hat{x} \right] \right\rangle_{\psi(t)} \\ &= \left\langle \frac{i}{\hbar} \underbrace{\frac{\hat{p}}{m} [\hat{p}, \hat{x}]}_{=-i\hbar} \right\rangle_{\psi(t)} = \frac{1}{m} \langle \hat{p} \rangle_{\psi(t)} = \frac{p_0}{m} \end{aligned}$$

$$\rightarrow \langle \hat{x} \rangle_{\psi(t)} = x_0 + \frac{p_0}{m} t$$

25b) 3D: gelte $\langle \hat{\vec{x}} \rangle_{\psi_0} = \vec{x}_0$, $\langle \hat{\vec{p}} \rangle_{\psi_0} = \vec{p}_0$.

dann $\langle \hat{\vec{x}} \rangle_{\psi(t)} = \vec{x}_0 + \frac{\vec{p}_0}{m} t$

denn: $\hat{p}_1, \hat{p}_2, \hat{p}_3$ Erhaltungsgrößen $\rightarrow \langle \hat{p}_l \rangle_{\psi(t)} = \langle \hat{p}_l \rangle_{\psi(0)} = (\vec{p}_0)_l$

$$\begin{aligned} \frac{d}{dt} \langle \hat{x}_l \rangle_{\psi(t)} &= \left\langle \frac{i}{\hbar} [H, \hat{x}_l] \right\rangle_{\psi(t)} = \left\langle \frac{i}{\hbar} \frac{1}{2m} [\hat{p}_1^2 + \hat{p}_2^2 + \hat{p}_3^2, \hat{x}_l] \right\rangle_{\psi(t)} = \\ &= \frac{1}{m} \langle \hat{p}_l \rangle_{\psi(t)} = \frac{1}{m} (\vec{p}_0)_l \end{aligned}$$

$2\hat{p}_l \cdot [\hat{p}_l, \hat{x}_l] = -2i\hbar \hat{p}_l$

$\rightarrow \langle \hat{x}_l \rangle_{\psi(t)} = (\vec{x}_0)_l + \frac{1}{m} (\vec{p}_0)_l \cdot t$ für $l=1, 2, 3$.

26a) $|\psi(t)\rangle = e^{-iHt/\hbar} |\psi(0)\rangle = \int \frac{d\hbar}{2\pi} \tilde{\psi}(\hbar) e^{-i\frac{\hbar^2}{2m} t/\hbar} |\chi_\hbar\rangle$

$= \int_{-\infty}^{\infty} \frac{d\hbar}{2\pi} \tilde{\psi}(\hbar) e^{-i\frac{\hbar^2}{2m} t} |\chi_\hbar\rangle$

$\hat{p}|\chi_\hbar\rangle = \hbar|\chi_\hbar\rangle$

b) $\psi(x, t) = \langle \varphi_x | \psi(t) \rangle = \int \frac{d\hbar}{2\pi} \tilde{\psi}(\hbar) e^{-i\frac{\hbar^2}{2m} t} \underbrace{\langle \varphi_x | \chi_\hbar \rangle}_{= e^{i\hbar x}}$

c) $\psi(x) = (4\pi\sigma_0^2)^{1/4} \int \frac{d\hbar}{2\pi} e^{-\frac{\sigma_0^2}{2}\hbar^2} e^{i\hbar x} = \frac{(4\pi\sigma_0^2)^{1/4}}{2\pi} \sqrt{\frac{2\pi}{\sigma_0^2}} e^{-\frac{x^2}{2\sigma_0^2}}$

$= \frac{1}{(\pi\sigma_0^2)^{1/4}} e^{-x^2/2\sigma_0^2}$

$\psi(x, t) \stackrel{B)}{=} (4\pi\sigma_0^2)^{1/4} \int \frac{d\hbar}{2\pi} e^{-\underbrace{\left(\frac{\sigma_0^2}{2} + i\frac{\hbar}{2m} t\right)}_{\parallel \frac{\sigma_0^2}{2} \left(1 + i\frac{\hbar}{m\sigma_0^2} t\right) = \frac{\sigma_0^2}{2} (1 + i t/\tau)}} \hbar^2 + i\hbar x$

$\parallel \frac{\sigma_0^2}{2} \left(1 + i\frac{\hbar}{m\sigma_0^2} t\right) = \frac{\sigma_0^2}{2} (1 + i t/\tau)$

$\parallel 1/\tau$

$= \frac{(4\pi\sigma_0^2)^{1/4} \sqrt{2\pi}}{2\pi \sqrt{\sigma_0^2 |1 + i t/\tau|}} e^{-\frac{x^2}{2\sigma_0^2 (1 + i t/\tau)}}$

$= \frac{1}{(\pi\sigma_0^2 |1 + i t/\tau|^2)^{1/4}} e^{+i \frac{x^2 t}{2\sigma_0^2}} e^{-\frac{x^2}{2\sigma_t^2}}$

$\uparrow$
 $\frac{1}{1 + i t/\tau} = \frac{1 - i t/\tau}{1 + t^2/\tau^2}$

$\rightarrow |\psi(x, t)|^2 = \frac{1}{\sqrt{\pi\sigma_t^2}} e^{-x^2/\sigma_t^2}$

