

Wahlg.: Aufenthaltshd. $S(\vec{x}, t) = |\psi(\vec{x}, t)|^2$

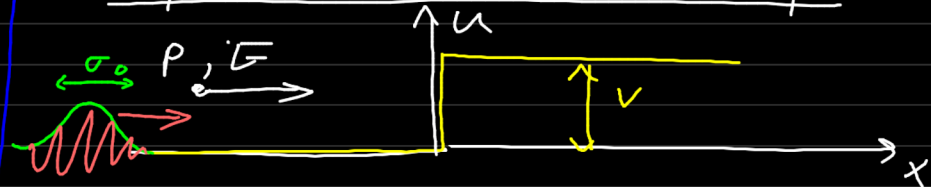
SCL, Wkterhaltung
 → Wkterstromdichte

$$\vec{j}(\vec{x}, t) = \frac{\hbar}{m} \operatorname{Im}(\psi^* \operatorname{grad} \psi)$$

→ Kont.-gl.

$$\frac{\partial S}{\partial t} + \operatorname{div} \vec{j} = 0$$

Ref. / Trans. an Pot.-stufe



Streuansatz

einfall. Teilchen

refl. Teilchen

$$\phi(x) = \begin{cases} e^{ikx} + r e^{-ikx} & : x < 0 \\ t e^{ikx} & : x \geq 0 \end{cases}$$

trans. Teilchen

+ St. S.Gl. + Stetigkeit

→ r, t

→ Reflexionswkt $R = |r|^2$.

• $E \leq V$: $R = 1$

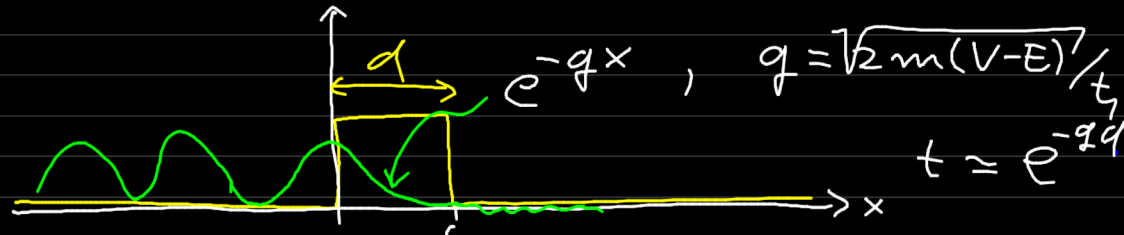
• $E > V$: $R = \left| \frac{1 - \sqrt{1 - \frac{V}{E}}}{1 + \sqrt{1 - \frac{V}{E}}} \right|^2 = \begin{cases} 1 - 4\sqrt{1 - \frac{V}{E}} & E - V \ll V \\ \left(\frac{V}{2E}\right)^2 & E \gg V \end{cases}$

auch für $E < V$ Transmission
mit Wkt. $T \approx e^{-2qd} > 0!$

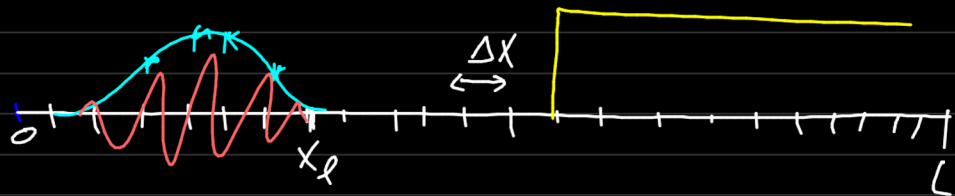
Tunneleffekt

(F. Hund, Gamov, ...)
1927 1928

Potenzialschwellen:



Voraussetzung des Streuansatzes:
Nurmerk!



$$\psi_l = \psi(x_l)$$

$$\hookrightarrow \psi'_l = \psi'(x_l) = (\psi_{l+1} - \psi_l) / \Delta x$$

$$\psi''_l = \psi''(x_l) = (\psi_{l+1} - 2\psi_l + \psi_{l-1}) / \Delta x^2$$

$$\psi(x, t) \xrightarrow{\Delta t} \psi(x, t + \Delta t)$$

$$\{\psi_l(t)\} \xrightarrow{\Delta t} \{\psi_l(t + \Delta t)\}_l$$

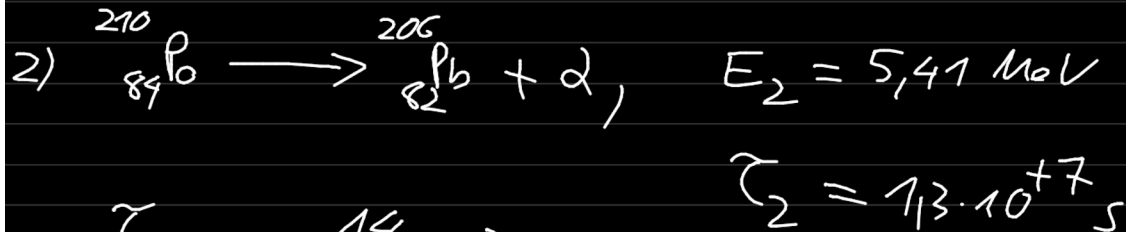
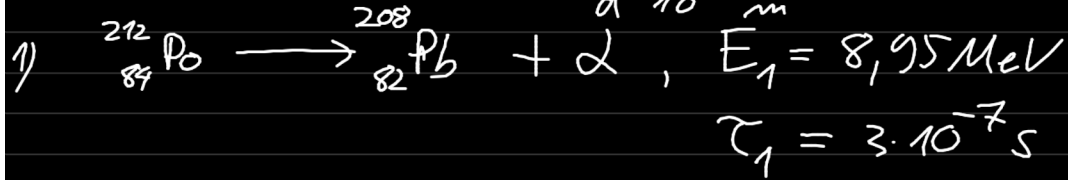
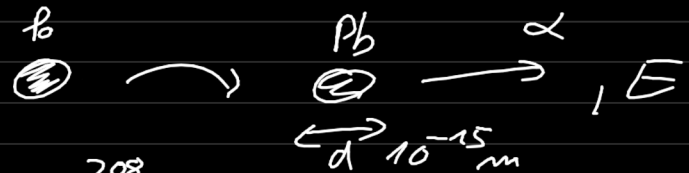
$$\psi(x_l, t + \Delta t) \approx \psi(x_l, t) + \dot{\psi}(x_l, t) \Delta t$$

$$\downarrow$$

$$\psi_l(t + \Delta t) = \psi_l(t) + \dot{\psi}_l(t) \Delta t$$

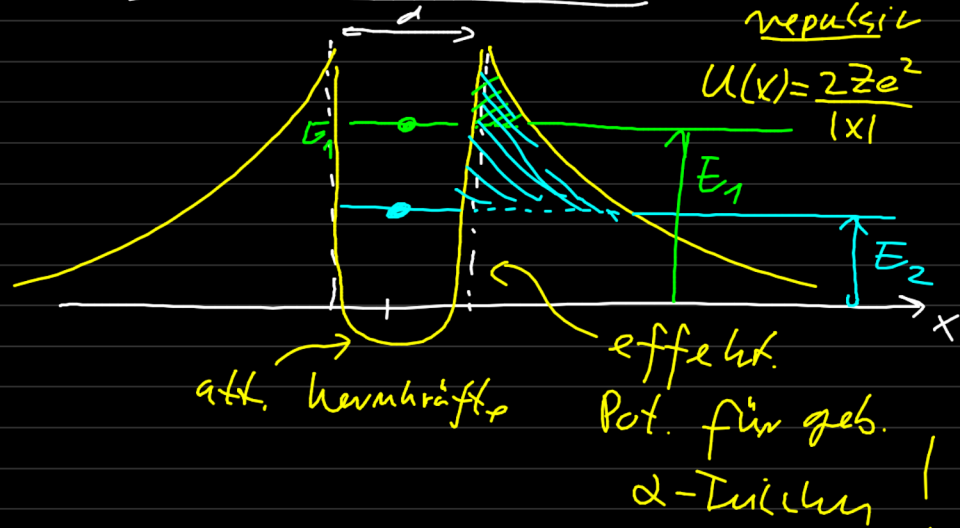
$$\dot{\psi}_l = \frac{i}{\hbar} \left(\frac{\hbar^2}{2m} \psi''_l - U_l \psi_l \right)$$

α -Zerfall vom Polonium zu Blei



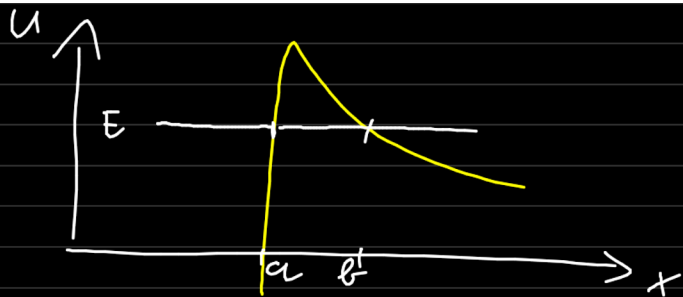
$\frac{\tau_2}{\tau_1} = 10^{14} \quad ?!$
 $\approx$

Gamov's Modell: (10)



Tunnel-Wkt für allg. Pot. schnell:

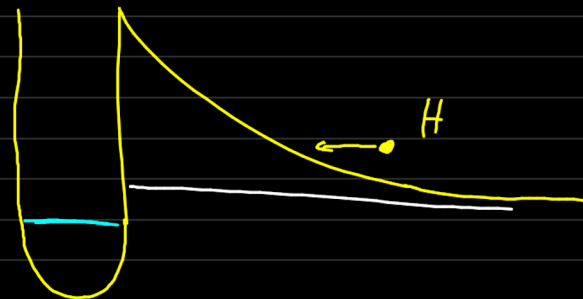
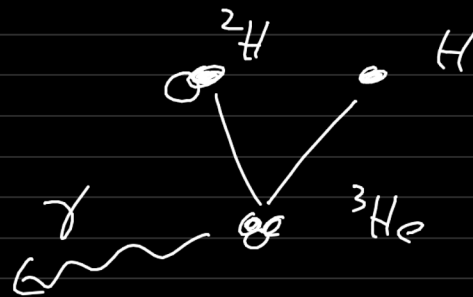
(T. Fließbach, Q11)



$$\rightarrow T = \exp\left(-\frac{1}{\hbar} \int_a^b \sqrt{8m(U(x) - E)} dx\right)$$

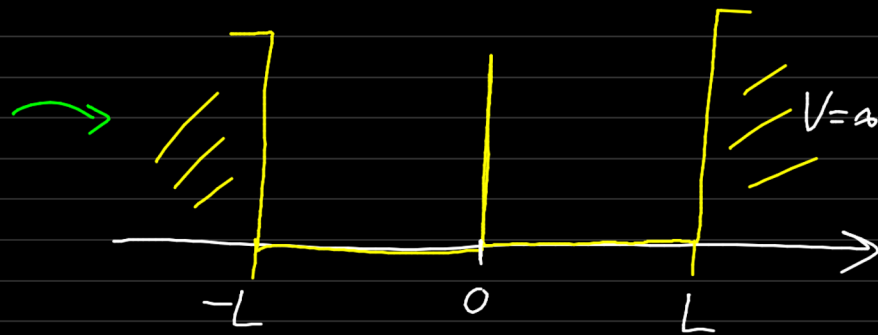
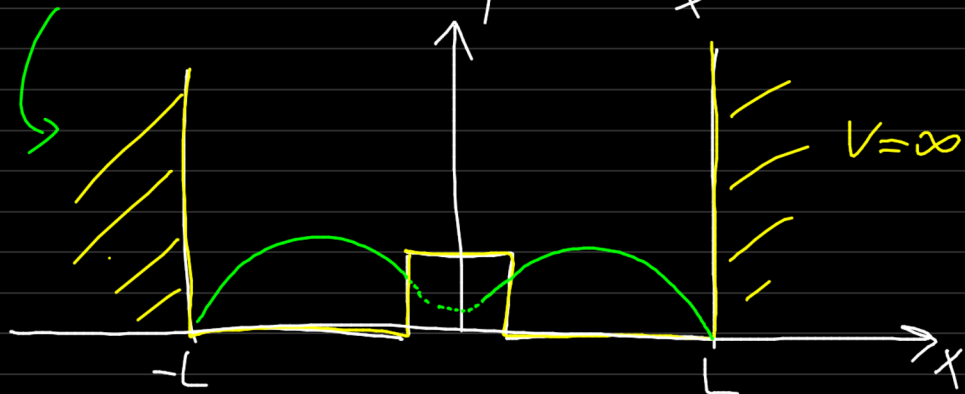
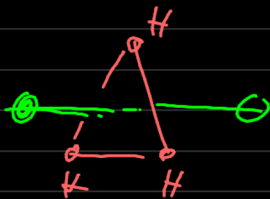
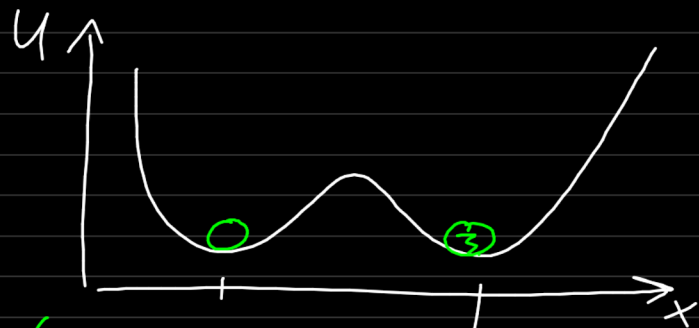
$$\rightarrow \frac{T_1}{T_2} = \frac{T_2}{T_1} = \begin{cases} 4 \cdot 10^{14} & ! \text{ (Theorie)} \\ 2 \cdot 10^{14} & \text{Exp.} \end{cases}$$

ähnlich: Kernfusion



Teilchen im Doppelmuldenpot.

($\rightarrow \text{NH}_3$)



$$U(x) = \begin{cases} a & : |x| > L \\ U_S(x) & : |x| \leq L \end{cases}$$

ges.: Lsg. $\phi(x)$ der stat. S.G.
 $x \in [-L, L]$

mit Randbed. $\phi(\pm L) = 0$

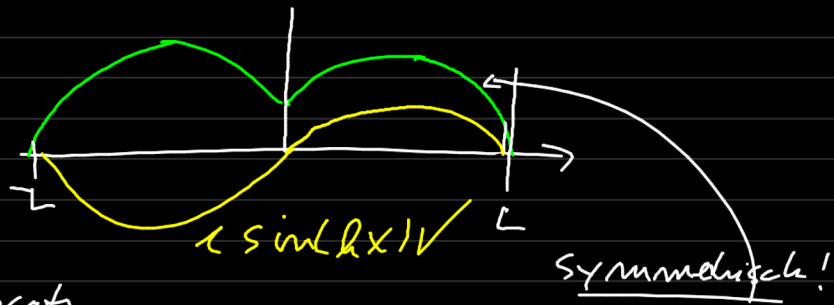
$$\phi''(x) = \frac{2m}{\hbar^2} (\mu \delta(x) - E) \phi$$

$$x \neq 0: \quad \phi''(x) = -k^2 \phi(x), \quad k = \sqrt{2mE}/\hbar$$

$$\hookrightarrow \phi(x): e^{\pm i k x}, \cos(kx), \sin(kx), \dots, \sin(kx + \varphi)$$

$x=0$: Unstetigkeit in ϕ' bei $x=0$!

$$\boxed{\Delta \phi' = \phi'(\varepsilon) - \phi'(-\varepsilon) = \int_{-\varepsilon}^{\varepsilon} \phi''(x) dx = \int_{-\varepsilon}^{\varepsilon} \left(\frac{2m}{\hbar^2} \mu \delta(x) - \frac{2mE}{\hbar^2} \right) \phi(x) dx = \frac{2m}{\hbar^2} \mu \phi(0)}$$



Ansatz $\phi(x) = \underline{1} \sin(\underline{k|x|} + \underline{\varphi})$

$$\hookrightarrow \phi(0) = 1 \sin(k \cdot 0)$$

$$\phi'(0) = 1 k \cos(k \cdot 0)$$

$$\rightarrow S\phi' = 2 \cdot h \cos(hb) \stackrel{!}{=} \frac{2m\alpha}{\hbar^2} \cdot \sin(hb) \rightarrow h_n^{\text{odd}} = \frac{\pi}{L+b} \cdot n \quad !$$

$$\tan(hb) \stackrel{!}{=} \frac{\hbar^2}{m\alpha}$$

$$\rightarrow \boxed{b = \frac{\hbar^2}{m\alpha}}$$

falls n groß
 $\leadsto hb$ klein
 $\leadsto \tan(hb) \approx hb$

R.B.: $0 \stackrel{!}{=} \sin(\underbrace{h(L+b)}_{\stackrel{!}{=} \pi \cdot n}) = \phi^{\text{even}}(L)$

$\rightarrow$ Eigenenergien symmetrisch E. fließen:

$$E_n^{\text{even}} = \frac{\hbar^2 h_n^2}{2m} = \frac{\hbar^2 \pi^2}{2m} \cdot \frac{n^2}{(L+b)^2}$$

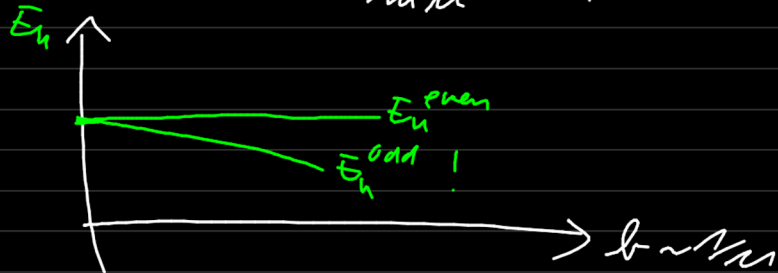
$$\rightarrow 0 \stackrel{!}{=} \sin(\underbrace{hL}_{2=\pi \cdot n}) = \phi^{\text{odd}}(L)$$

$$\rightarrow h_n^{\text{even}} = \frac{\pi}{L} \cdot n$$

$$E_n^{\text{odd}} = \frac{\hbar^2 (k_n^{\text{odd}})^2}{2m} = \frac{\hbar^2 \pi^2}{2m} \frac{n^2}{L^2}$$

for $\mu \rightarrow \infty$: $E_n^{\text{odd}} = E_n^{\text{even}}$

$\mu < \infty$: $\rightarrow b = \frac{\hbar^2}{m\mu} > 0!$



$\rightarrow$ Ground state energy:

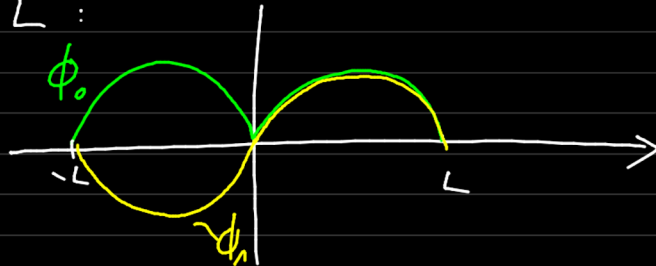
$$E_0 = E_1^{\text{even}} = \frac{\hbar^2 \pi^2}{2m} \frac{1}{(L+b)^2}$$

$$E_1 = E_1^{\text{odd}} = \frac{\hbar^2 \pi^2}{2m} \frac{1}{L^2}$$

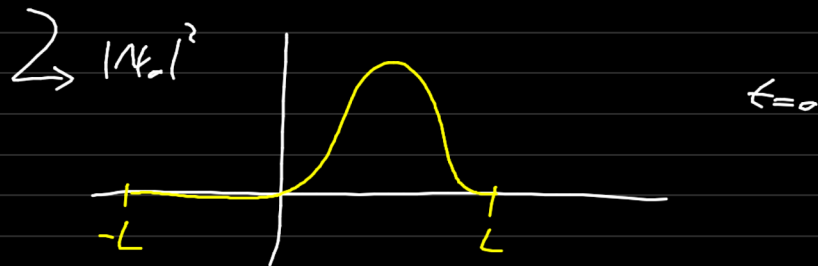
$$\phi_0(x) = \sqrt{2} \sin(k_1^{\text{ev}}(x+b))$$

$$\phi_1(x) = \sqrt{2} \sin(k_1^{\text{od}} x)$$

$0 < b \ll L$:



Dynamik: $|\psi\rangle = \frac{1}{\sqrt{2}}(|\phi_0\rangle + |\phi_1\rangle) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$



$$B = \langle |\phi_0\rangle, |\phi_1\rangle \rangle$$

$$E_0 = \frac{\hbar^2 \gamma^2}{2m} \frac{1}{(L+b)^2}$$

$$E_1 = \frac{\hbar^2 \pi^2}{2m} \frac{1}{L^2} \quad b \ll L$$

$$\hookrightarrow |\psi(t)\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} e^{-iE_0 t/\hbar} \\ e^{iE_1 t/\hbar} \end{pmatrix} = \frac{e^{-i\Omega t}}{\sqrt{2}} \begin{pmatrix} e^{+i\frac{\omega t}{2}} \\ e^{-i\frac{\omega t}{2}} \end{pmatrix}$$

$$\Delta E = E_1 - E_0 = E_1 \left(1 - \frac{1}{(1+b^2/L^2)} \right) \ll E_1$$

$$\Omega = \frac{E_0 + E_1}{2\hbar}$$

$$\omega = \frac{\Delta E}{\hbar} \ll \Omega !$$

$$\sum |\psi(t)|^2 = \dots$$