

gestern:

3. Postulat: Dynamik

$$\frac{d}{dt} |\psi(t)\rangle = -\frac{i}{\hbar} H |\psi(t)\rangle$$

$$i\hbar \frac{d}{dt} |\psi(t)\rangle = H |\psi(t)\rangle$$

Schrödinger-Gleichung

- H : Hamilton-Operator $\equiv$ Energie-Op.
($H = H^\dagger$)
- $\hbar = 1,054 \dots \cdot 10^{-34} \text{ Js}$

$$\rightarrow \frac{d}{dt} \langle A \rangle_{\psi(t)} = \left\langle \frac{i}{\hbar} [H, A] \right\rangle_{\psi(t)}$$

Kommutator: $[A, B] := AB - BA$

$$\rightarrow A \text{ Erhaltungsgröße} \Leftrightarrow [H, A] = 0$$

Lösung der Schrödinger-Gleichung zum
Anfangszustand $|\psi(0)\rangle = |\psi_0\rangle$:

$$|\psi(t)\rangle = e^{-\frac{i}{\hbar} H t} |\psi_0\rangle$$

$$\text{[energy } \dot{y} = -i\omega y$$

$$\rightarrow y(t) = e^{-i\omega t} y_0]$$

explizit in "Energie darstellung":

d.h.:

$$H = \sum_{n=0}^N E_n |\varphi_n\rangle \langle \varphi_n|$$

Eigenenergien $E_0 \leq E_1 \leq E_2 \dots$ (Eigenwerte)

Energieeigenzustände $|\varphi_0\rangle, |\varphi_1\rangle, |\varphi_2\rangle \dots$ (Eigenvektoren)

von H

$$H |\varphi_n\rangle = E_n |\varphi_n\rangle$$

"stationäre Schrödingergl."

(Eigenwertgleichung)

$$\rightarrow e^{-\frac{i}{\hbar} H t} = \exp\left(-\frac{i}{\hbar} H t\right) = \sum_{n=0}^N e^{-i\omega_n t} |\varphi_n\rangle \langle \varphi_n|$$

$$\omega_n = E_n / \hbar$$

$$\rightarrow e^{-\frac{i}{\hbar} H t} \equiv \exp\left(-\frac{i}{\hbar} H t\right) = \sum_{n=0}^{\infty} e^{-i \omega_n t} |\varphi_n\rangle \langle \varphi_n| \quad (*)$$

$$\omega_n = E_n / \hbar$$

• $|\psi_0\rangle = |\varphi_h\rangle \xrightarrow{t} |\psi(t)\rangle = e^{-\frac{i}{\hbar} H t} |\varphi_h\rangle$
 $= e^{-i \omega_h t} |\varphi_h\rangle$

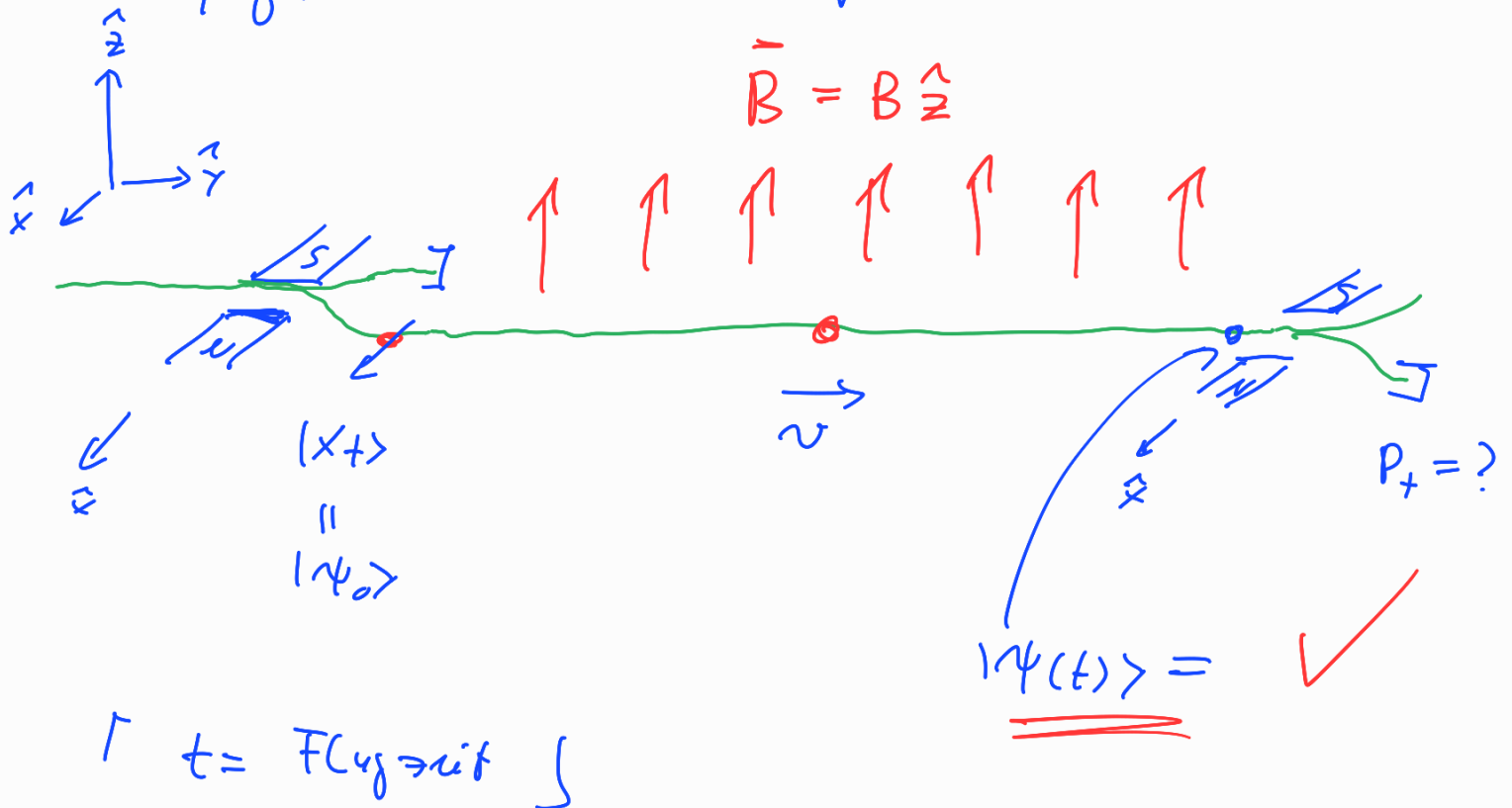
$\uparrow$
Energieeigen-
zustand

$$\left[e^{-\frac{i}{\hbar} H t} |\varphi_h\rangle = \sum_{n=0}^{\infty} \overbrace{e^{-i \hbar t / \hbar}}^{= 1} e^{-i \omega_n t} |\varphi_n\rangle \underbrace{\langle \varphi_n | \varphi_h \rangle}_{= \delta_{nh}} = e^{-i \omega_h t} |\varphi_h\rangle \right]$$

• $|\psi_0\rangle = \sum_h a_h |\varphi_h\rangle \xrightarrow{t} |\psi(t)\rangle = \sum_h a_h e^{-\frac{i}{\hbar} H t} |\varphi_h\rangle$
 $= \sum_h a_h e^{-i \omega_h t} |\varphi_h\rangle$

Dynamik eines Spins im Magnetfeld

2-stufiges Stern-Gerlach-Experiment:



Hamilton-Funktion

$$H = -\vec{B} \cdot \vec{\mu}$$

$$H = -B \mu_z$$

$(\vec{B} = B \hat{z})$

Hamilton-Operatoren:

$$\hat{H} = -B \hat{\mu}_z$$

$$\hat{\mu}_z = \frac{\mu_0}{2} |z+\rangle \langle z+| - \frac{\mu_0}{2} |z-\rangle \langle z-| \leftarrow \begin{cases} +\mu_0/2, & -\mu_0/2 \\ |z+\rangle & |z-\rangle \end{cases}$$

$$\rightarrow H = -\frac{B\mu_0}{2} |z+\rangle \langle z+| + \frac{B\mu_0}{2} |z-\rangle \langle z-|$$

$$|\psi_0\rangle = |x+\rangle \xrightarrow{t} |\psi(t)\rangle = e^{-\frac{i}{\hbar} H t} |x+\rangle$$

$$\rightarrow H = - \frac{B\mu_0}{2} |z+\rangle\langle z+| + \frac{B\mu_0}{2} |z-\rangle\langle z-|$$

$$|\psi_0\rangle = |x+\rangle \xrightarrow{t} |\psi(t)\rangle = e^{-\frac{i}{\hbar} H t} |x+\rangle$$

$$e^{-\frac{i H t}{\hbar}} = \exp\left(i \frac{B\mu_0}{2\hbar} t |z+\rangle\langle z+| - i \frac{B\mu_0}{2\hbar} t |z-\rangle\langle z-| \right)$$

$$= \exp\left(i \frac{\omega}{2} t |z+\rangle\langle z+| - i \frac{\omega}{2} t |z-\rangle\langle z-| \right)$$

mit $\omega = \frac{B\mu_0}{\hbar}$ (Larmorfrequenz)

$$\rightarrow e^{-\frac{i H t}{\hbar}} = e^{i \frac{\omega}{2} t} |z+\rangle\langle z+| + e^{-i \frac{\omega}{2} t} |z-\rangle\langle z-|$$

$$\rightarrow |\psi(t)\rangle = e^{-\frac{i H t}{\hbar}} |x+\rangle$$

$$|x+\rangle = (|z+\rangle + |z-\rangle) / \sqrt{2} \quad : \quad |z+\rangle = (|x+\rangle + |x-\rangle) / \sqrt{2}$$

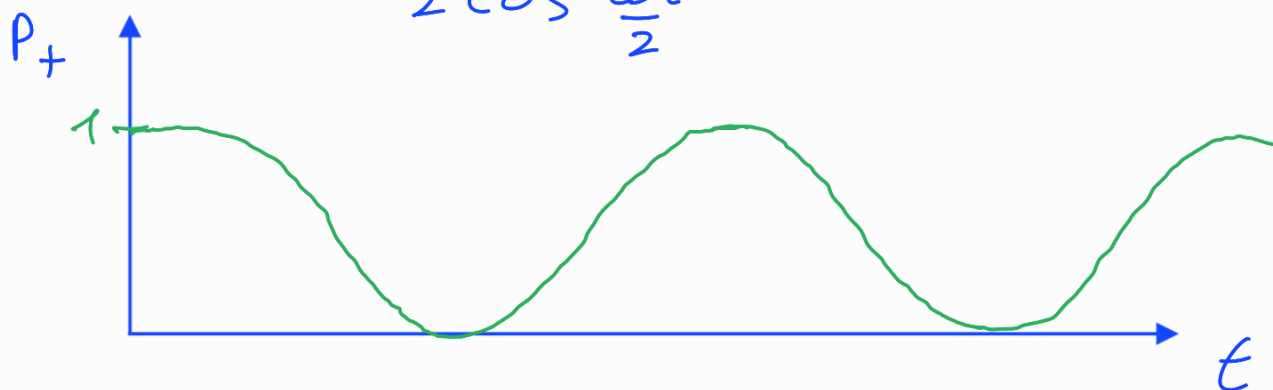
$$|x-\rangle = (|z+\rangle - |z-\rangle) / \sqrt{2} \quad : \quad |z-\rangle = (|x+\rangle - |x-\rangle) / \sqrt{2}$$

$$|\psi(t)\rangle = e^{i \frac{\omega t}{2}} \underbrace{|z+\rangle\langle z+|}_{1/\sqrt{2}} |x+\rangle + e^{-i \frac{\omega t}{2}} \underbrace{|z-\rangle\langle z-|}_{1/\sqrt{2}} |x+\rangle$$

$$|\psi(t)\rangle = e^{i\frac{\omega t}{2}} \underbrace{|z+\rangle}_{1/\sqrt{2}} \langle z+|x+\rangle + e^{-i\frac{\omega t}{2}} \underbrace{|z-\rangle}_{1/\sqrt{2}} \langle z-|x+\rangle$$

$$\text{d.l. } |\psi(t)\rangle = \frac{1}{\sqrt{2}} \left(e^{i\frac{\omega t}{2}} |z+\rangle + e^{-i\frac{\omega t}{2}} |z-\rangle \right)$$

$$\begin{aligned} \rightarrow P_+ &= |\langle x+ | \psi(t) \rangle|^2 = \frac{1}{2} \left| \frac{1}{\sqrt{2}} e^{i\frac{\omega t}{2}} + \frac{1}{\sqrt{2}} e^{-i\frac{\omega t}{2}} \right|^2 \\ &= \frac{1}{4} \left| \underbrace{e^{i\frac{\omega t}{2}} + e^{-i\frac{\omega t}{2}}}_{2 \cos \frac{\omega t}{2}} \right|^2 = \cos^2 \frac{\omega t}{2} \end{aligned}$$



$$\langle \mu_x \rangle_t = \langle \psi(t) | \hat{\mu}_x | \psi(t) \rangle$$

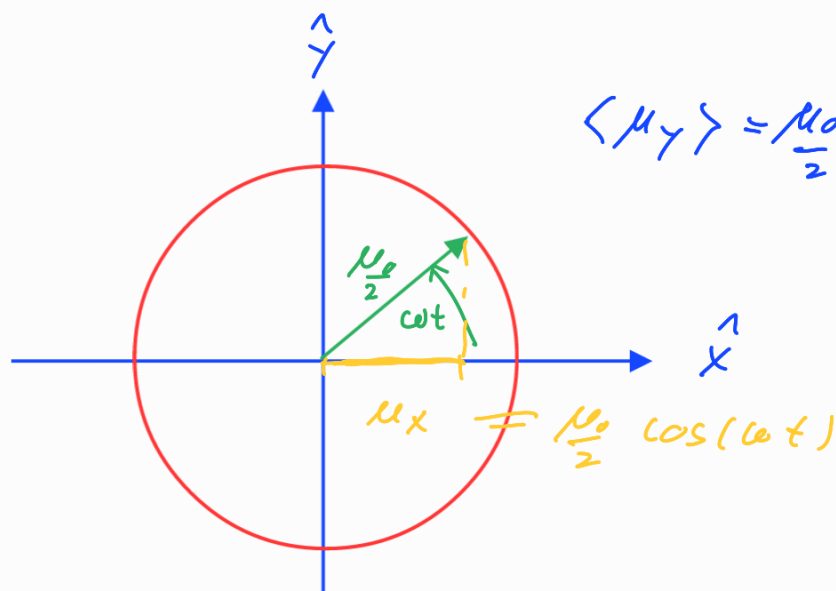
$$= P_+ \cdot (+\mu_0) + \overset{1-P_+}{P_-} \cdot \left(-\mu_0\right)$$

$$= P_+ \cdot \mu_0 - \mu_0 = \frac{\mu_0}{2} \underbrace{\left(2 \cos^2 \left(\frac{\omega t}{2} \right) - 1 \right)}_{\cos \omega t}$$

$$\begin{aligned}
 \Gamma_{NR}: \quad & 2 \cos^2\left(\frac{d}{2}\right) - 1 \quad (*) \\
 &= \cos^2\left(\frac{d}{2}\right) + \cos^2\left(\frac{d}{2}\right) - 1 \\
 &= \cos^2\left(\frac{d}{2}\right) + \underbrace{1 - \sin^2\left(\frac{d}{2}\right)} - 1 \\
 &= \cos^2 \frac{d}{2} - \sin^2 \frac{d}{2} \\
 &= \cos(d)
 \end{aligned}$$

$$\cos(a+b) = \cos a \cos b - \sin a \sin b$$

$$\langle \mu_x \rangle_t = \frac{\mu_0}{2} \cos \omega t$$



$$\langle \mu_z \rangle_t \stackrel{?}{=} \langle \mu_z \rangle_0$$

$$\hat{H} = -B \hat{\mu}_z \rightarrow [\hat{H}, \mu_z]$$

$$= [-B \mu_z, \mu_z]$$

$$= -B [\mu_z, \mu_z] = 0!$$

$$\vec{B} \parallel \hat{z} \rightarrow \hat{\mu}_z \text{ Erhaltungsgröße!}$$

