

Wdhlg.: Zeitentwicklungsoperator

$$U(t) = e^{-iHt/\hbar}$$

$$\leadsto |\psi_0\rangle \xrightarrow{t} |\psi(t)\rangle = U(t) |\psi_0\rangle$$

- $U(t)$ unitär: $U(t)^\dagger = U^{-1}(t)$

(d.h. $U(t)$ Norm-erhaltend, invariant)

- $U(t+t') = U(t)U(t')$

Energiedarstellung:

$$H = \sum_n E_n |\varphi_n\rangle \langle \varphi_n|$$

Eigenenergien
Energiezustände



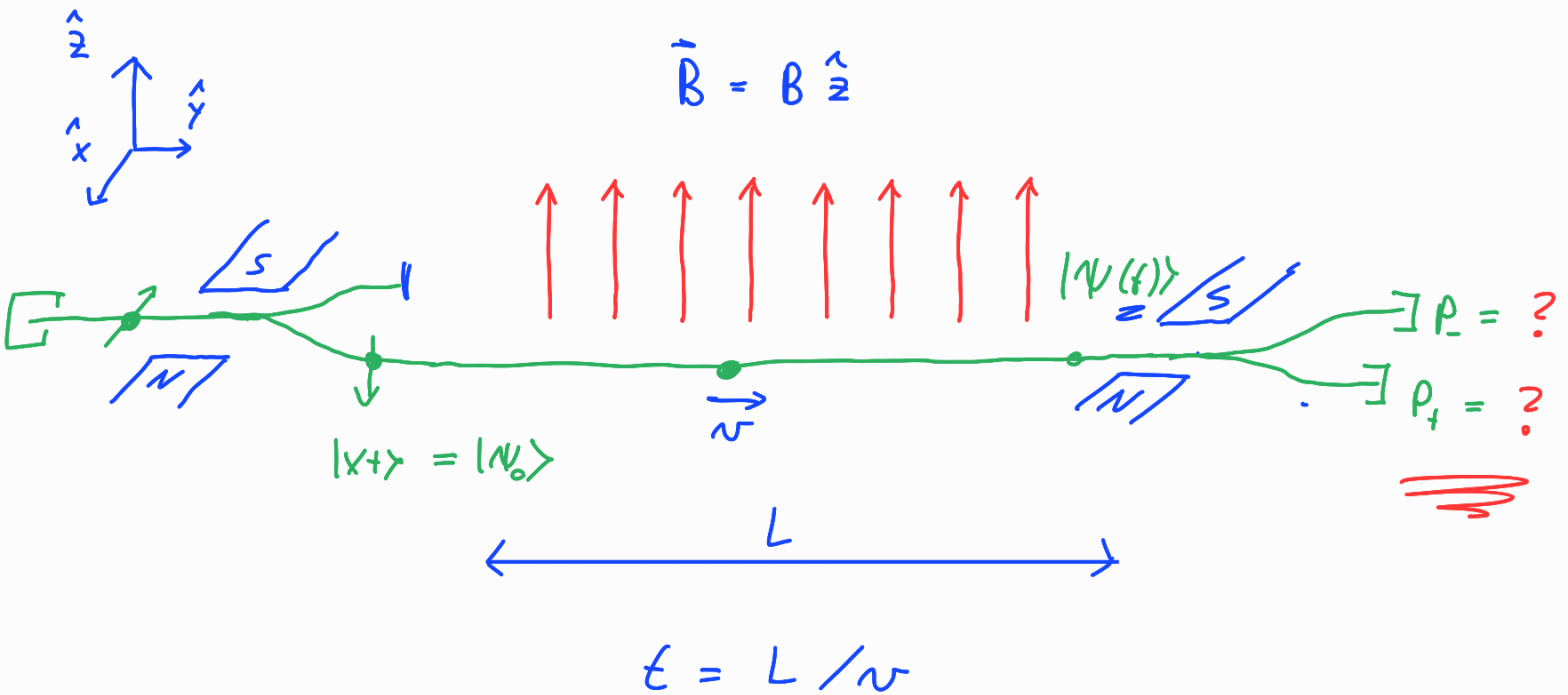
$$U(t) = \sum_n e^{-iE_n t/\hbar} |\varphi_n\rangle \langle \varphi_n| = \sum_n e^{-i\omega_n t} |\varphi_n\rangle \langle \varphi_n|$$

($H|\varphi_n\rangle = E_n|\varphi_n\rangle$)

$$\leadsto |\psi_0\rangle = \sum_n a_n |\varphi_n\rangle \xrightarrow{t} |\psi(t)\rangle = \sum_n e^{-i\omega_n t} a_n |\varphi_n\rangle$$

Beispiel: Dynamik eines Spins im homog. B-Feld
 $\hat{S} \hat{=} \text{mag. Moment}$

Zweistufiges Stern-Gerlach-Experiment:



Hamilton-Op.: $H = -B \hat{\mu}_z$

$$= \mu_0 B (|z-\rangle\langle z-| - |z+\rangle\langle z+|) = \mu_0 B \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$$

d.h. $|z\pm\rangle$ Energiezust. zu Eigenenergien $\mp \mu_0 B$

$\rightarrow U(t) = e^{-i\omega t} |z-\rangle\langle z-| + e^{+i\omega t} |z+\rangle\langle z+| \quad \omega = \frac{\mu_0 B}{\hbar}$

$\rightarrow |\psi_0\rangle = |x+\rangle = \frac{|z+\rangle + |z-\rangle}{\sqrt{2}} \xrightarrow{t} |\psi(t)\rangle = \frac{1}{\sqrt{2}} \left(e^{+i\omega t} |z+\rangle + e^{-i\omega t} |z-\rangle \right)$

$$|\psi(t)\rangle = \frac{1}{\sqrt{2}} (e^{+i\omega t} |z+\rangle + e^{-i\omega t} |z-\rangle)$$

$$= \cos(\omega t) |x+\rangle + i \sin(\omega t) |x-\rangle$$

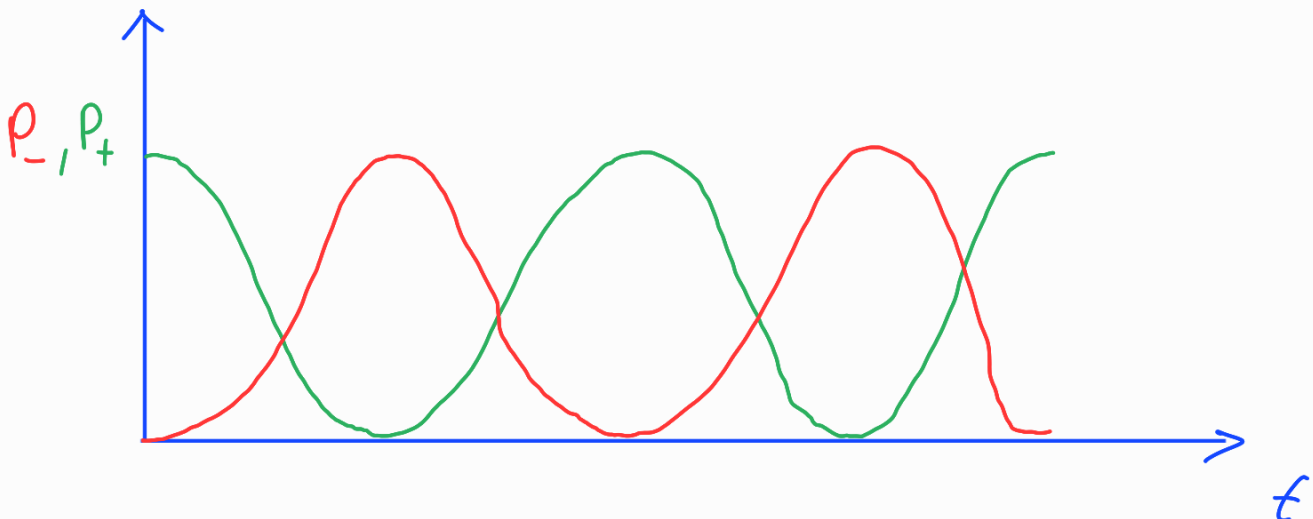
$$|z\pm\rangle = \frac{|x+\rangle \pm |x-\rangle}{\sqrt{2}}$$

$$e^{i\omega t} \pm e^{-i\omega t} = 2 \begin{cases} \cos(\omega t) \\ i \sin(\omega t) \end{cases}$$

$$\rightarrow P_+(t) = |\langle x+ | \psi(t) \rangle|^2 = \cos^2(\omega t)$$

$$P_-(t) = |\langle x- | \psi(t) \rangle|^2 = \sin^2(\omega t) \quad (= 1 - P_+(t))$$

$$\begin{aligned} \rightarrow \langle \mu_x \rangle_t &= \mu_0 (P_+(t) - P_-(t)) = \\ &= \mu_0 (\cos^2(\omega t) - \sin^2(\omega t)) \\ &= \mu_0 \cos(2\omega t) \end{aligned}$$



alternative Rechnungen:

$$b) U(t) = e^{i\omega t} |z+\rangle\langle z+| + e^{-i\omega t} |z-\rangle\langle z-| = \begin{pmatrix} e^{i\omega t} & 0 \\ 0 & e^{-i\omega t} \end{pmatrix}$$

$$\rightarrow P_{+}(t) = |\langle x+ | \psi(t) \rangle|^2 = |\langle x+ | U(t) | x+ \rangle|^2$$

$$= \left| \frac{1}{2} (1, 1) \begin{pmatrix} e^{i\omega t} & 0 \\ 0 & e^{-i\omega t} \end{pmatrix} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \right|^2$$

$$= \left| \frac{1}{2} (1, 1) \begin{pmatrix} e^{i\omega t} \\ e^{-i\omega t} \end{pmatrix} \right|^2 = \left| \frac{1}{2} (e^{i\omega t} + e^{-i\omega t}) \right|^2$$

$$= \cos^2(\omega t)$$

c) mittels Pauli-Operatoren:

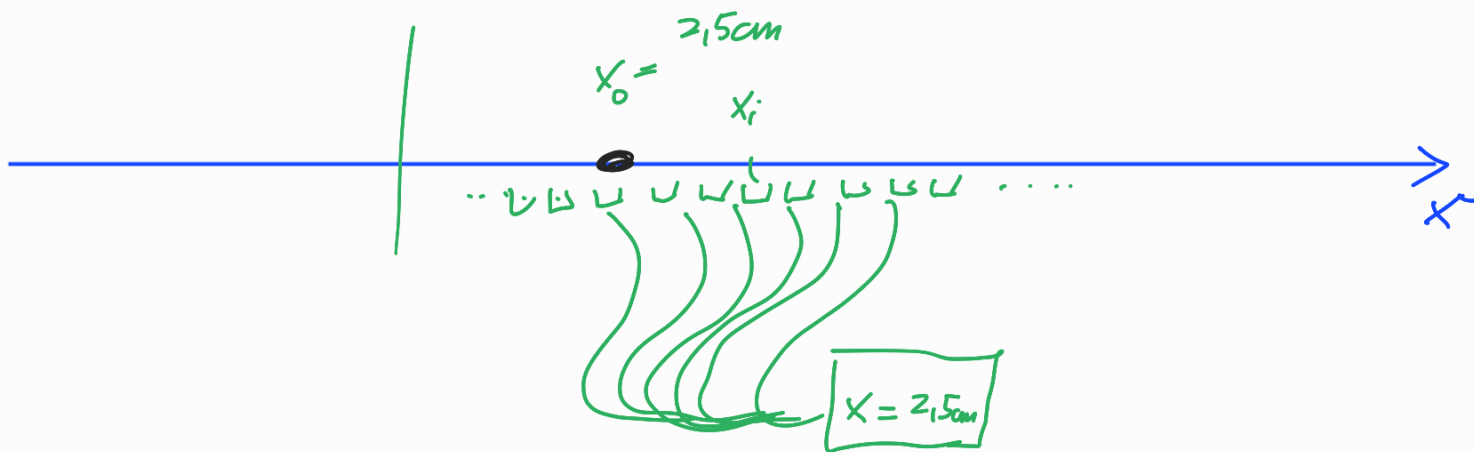
$$H = -B\hat{\mu}_z = -B\mu_0 \hat{\sigma}_z \quad \left(= -B\mu_0 \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \right)$$

$$\rightarrow U(t) = e^{i \frac{B\mu_0 t}{\hbar} \cdot \hat{\sigma}_z} = \cos(\omega t) \frac{1}{2} + i \sin(\omega t) \hat{\sigma}_z$$

Euler-Formel

$$= \begin{pmatrix} e^{i\omega t} & 0 \\ 0 & e^{-i\omega t} \end{pmatrix} \quad \text{dann wie in b)}$$

Quantenmechanik des Punktteilchens (1D)



Ortsmessung: Resultat x $\in \mathbb{R}$

\rightarrow Ortsobservable \hat{x} = herm. Operator

mit Kontinuum an EWen $\{x\}_{x \in \mathbb{R}}$

mit zugeh. EZen $\{|\varphi_x\rangle\}_{x \in \mathbb{R}}$

\rightarrow Impassung des Dirac-Formalismus

diskret (D)

Kontinuierlich (K)

D

Obs. $\hat{A} = \hat{A} = \sum_i a_i |\varphi_i\rangle \langle \varphi_i|$

EWe $\{a_i\}_{i=1,2,3,\dots}$

Eze $\{|\varphi_i\rangle\}_{i=1,2,3,\dots}$

↑
ONB!

$\sum_i \dots \rightarrow \int_{-\infty}^{\infty} dx \dots$ K

Obs. $X = \text{Obs. op. } \hat{X}$

$\hat{X} = \int_{-\infty}^{\infty} dx x |\varphi_x\rangle \langle \varphi_x|$

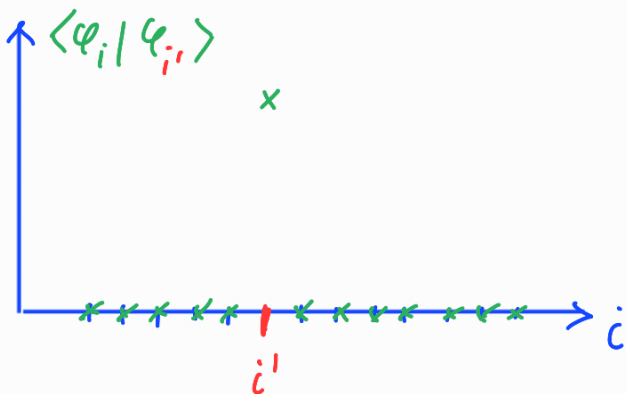
EWe $\{x\}_{x \in \mathbb{R}}$

Eze $\{|\varphi_x\rangle\}_{x \in \mathbb{R}}$

Orthogonalität

klar $\{|\varphi_i\rangle\}$:

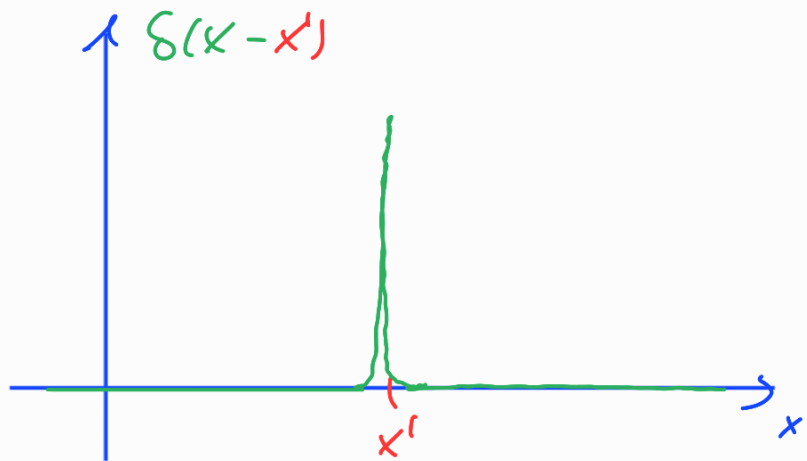
$\langle \varphi_i | \varphi_{i'} \rangle = \delta_{ii'}$

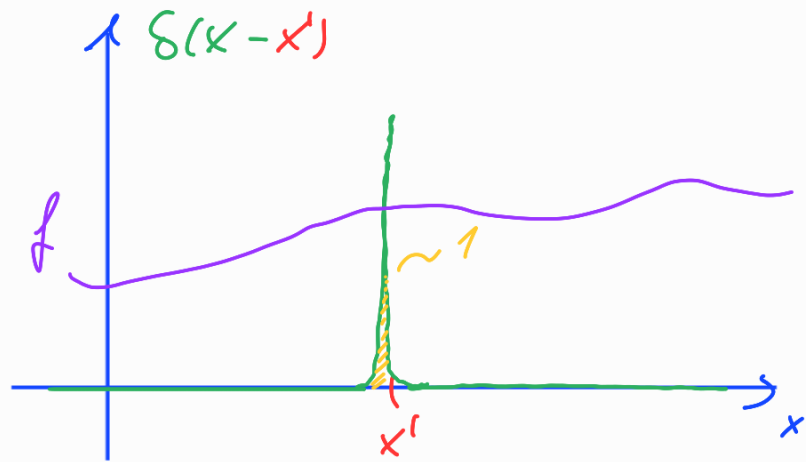
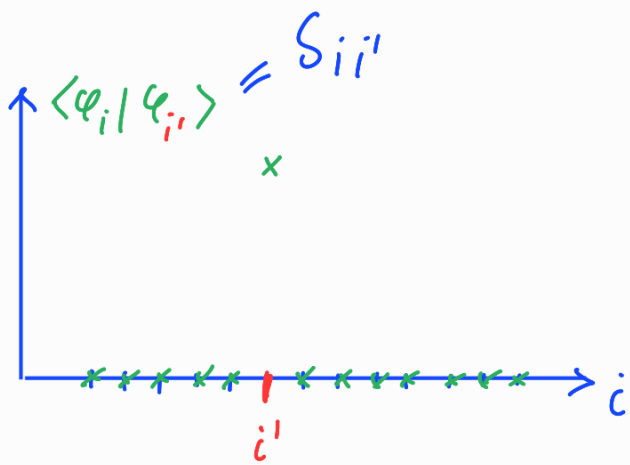


klar Obs. zu $\{|\varphi_x\rangle\}_{x \in \mathbb{R}}$

$\langle \varphi_x | \varphi_{x'} \rangle = \delta(x-x')$

↑
Dirac-Delta-Fkt.





- $i \neq i' : \delta_{ii'} = 0$

- $\sum_i \delta_{ii'} = 1$

- $\sum_i \delta_{ii'} \cdot f_i = f_{i'}$

- $x \neq x' : \delta(x-x') = 0$

- $\int_{\mathbb{R}} dx \delta(x-x') = 1$

- $\int_{\mathbb{R}} dx \delta(x-x') f(x) = f(x')$

↑ definierende Eigenschaften

• der δ -Fkt \equiv Rechenregel!

Vollständigkeit

den Zust. $\{|\varphi_i\rangle\}_{i=1,2,\dots}$

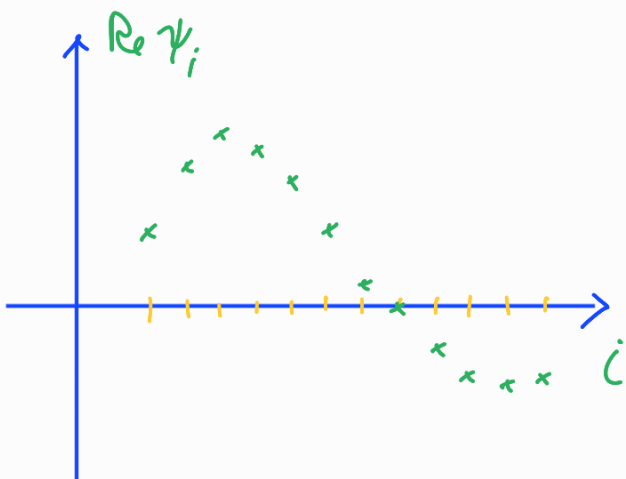
$$\mathbb{1}_{\mathcal{X}} = \sum_i |\varphi_i\rangle\langle\varphi_i|$$

↳ Komponentenentwicklung
des Zustands $|\psi\rangle$:

$$\begin{aligned} |\psi\rangle &= \mathbb{1}_{\mathcal{X}} |\psi\rangle \\ &= \left(\sum_i |\varphi_i\rangle\langle\varphi_i| \right) |\psi\rangle \\ &= \sum_i \underbrace{\langle\varphi_i|\psi\rangle}_{\psi_i} |\varphi_i\rangle \end{aligned}$$

$$|\psi\rangle = \sum_i \psi_i |\varphi_i\rangle$$

\nearrow
 i -te Komponente von $|\psi\rangle$



den Zust. $\{|\varphi_x\rangle\}_{x \in \mathbb{R}}$

$$\mathbb{1}_{\mathcal{X}} = \int_{\mathbb{R}} dx |\varphi_x\rangle\langle\varphi_x|$$

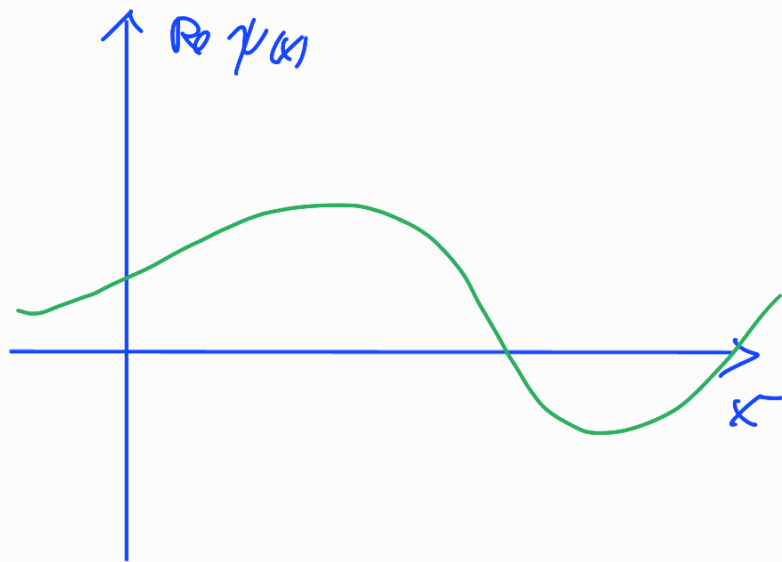
↳ Ortsdarstellung des
Zustands $|\psi\rangle$

$$\begin{aligned} |\psi\rangle &= \mathbb{1}_{\mathcal{X}} |\psi\rangle \\ &= \left(\int_{\mathbb{R}} dx |\varphi_x\rangle\langle\varphi_x| \right) |\psi\rangle \\ &= \int_{\mathbb{R}} dx |\varphi_x\rangle \underbrace{\langle\varphi_x|\psi\rangle}_{\psi(x)} \end{aligned}$$

mit Wellenfunktion des
Zustands $|\psi\rangle$ (am Ort x):

$$\psi(x) := \langle\varphi_x|\psi\rangle :$$

$$|\psi\rangle = \int_{\mathbb{R}} dx \psi(x) |\varphi_x\rangle$$



Skalarprodukt

$$\begin{aligned}
 \langle \psi | \chi \rangle &= \left(\sum_i \psi_i^* \langle \phi_i | \right) \left(\sum_j \chi_j | \phi_j \rangle \right) \\
 &= \sum_{ij} \psi_i^* \chi_j \underbrace{\langle \phi_i | \phi_j \rangle}_{\delta_{ij}} \\
 &= \sum_i \psi_i^* \chi_i
 \end{aligned}$$

$$\langle \psi | \chi \rangle = \sum_i \psi_i^* \chi_i$$

$$\begin{aligned}
 \langle \psi | \chi \rangle &= \left(\int dx \psi^*(x) \langle \phi_x | \right) \cdot \left(\int dx' \chi(x') | \phi_{x'} \rangle \right) \\
 &= \int dx \int dx' \psi^*(x) \chi(x') \underbrace{\langle \phi_x | \phi_{x'} \rangle}_{\delta(x-x')} \\
 &= \int dx \psi^*(x) \chi(x)
 \end{aligned}$$

$$\langle \psi | \chi \rangle = \int dx \psi^*(x) \chi(x)$$

